

# The role of nitrogen doping in ALD Ta<sub>2</sub>O<sub>5</sub> and its influence on multilevel cell switching in RRAM

N. Sedghi<sup>1, a)</sup>, H. Li<sup>2</sup>, I. F. Brunell<sup>3</sup>, K. Dawson<sup>3</sup>, R. J. Potter<sup>3</sup>, Y. Guo<sup>2, 4</sup>, J. T. Gibbon<sup>5</sup>, V. R. Dhanak<sup>5</sup>, W. D. Zhang<sup>6</sup>, J. F. Zhang<sup>6</sup>, J. Robertson<sup>2</sup>, S. Hall<sup>1</sup>, and P. R. Chalker<sup>3</sup>

<sup>1</sup>*Department of Electrical Engineering and Electronics, University of Liverpool, Liverpool, L69 3GJ, United Kingdom*

<sup>2</sup>*Department of Engineering, University of Cambridge, Cambridge, CB2 1TN, United Kingdom*

<sup>3</sup>*School of Engineering, University of Liverpool, Liverpool, L69 3GH, United Kingdom*

<sup>4</sup>*College of Engineering, Swansea University, Swansea SA1 8EN, United Kingdom*

<sup>5</sup>*Department of Physics, University of Liverpool, Liverpool, L69 7ZE, United Kingdom*

<sup>6</sup>*Department of Electronics and Electrical Engineering, Liverpool John Moores University, Liverpool, L3 3AF, United Kingdom*

The role of nitrogen doping on the stability and memory window of resistive state switching in N-doped Ta<sub>2</sub>O<sub>5</sub> deposited by atomic layer deposition is elucidated. Nitrogen incorporation increases the stability of resistive memory states which is attributed to neutralization of electronic defect levels associated with oxygen vacancies. Density functional simulations with screened exchange hybrid functional approximation show that the incorporation of nitrogen dopant atoms in the oxide network removes the O vacancy midgap defect states, thus nullifying excess defects and eliminating alternative conductive paths. By effectively reducing the density of vacancy-induced defect states through N doping, 3-bit multilevel cell switching is demonstrated, consisting of eight distinctive resistive memory states achieved by either controlling the set current compliance or the maximum voltage during reset. Nitrogen doping has a threefold effect; widening the switching memory window to accommodate more intermediate states, improving the stability of states, and providing gradual reset for multilevel cell switching during reset. The N-doped Ta<sub>2</sub>O<sub>5</sub> devices have relatively small set and reset voltages (< 1 V) with reduced variability due to doping.

Resistive-switching random access memory (RRAM) is the focus of development as a substitute for existing non-volatile memory (NVM) devices. In metal oxide based RRAM devices, switching between low resistive and high resistive states (LRS and HRS, respectively) occurs by field-assisted diffusion of defects such as oxygen vacancies (O<sub>vac</sub>), to form a conductive ‘filament’ (CF)<sup>1</sup>. By applying appropriate voltages across the oxide, the CF can be formed, ruptured, and restored to switch between the ‘on’ and ‘off’ memory states. It is crucial to engineer the oxygen vacancy profile in the oxide film to achieve stable memory states and a large switching memory window for immunity to noise<sup>2</sup>. Extending this requirement further leads to the

1. Corresponding author electronic mail: nsed@liverpool.ac.uk.

prevention of multiple potential CFs with different LRS resistances, which is the main cause of the variability in the multilevel cell (MLC) switching states<sup>3</sup>. The MLC switching enables multi-bit storage capacity and potentially increases the storage capacity for ultra-high density memory applications. It has been reported that the variability of resistance associated with a CF depends on the conductive filament size and also the oxygen vacancy concentration<sup>4,6</sup>. The multilevel switching is commonly achieved by adjusting the LRS resistance by setting the current compliance during the set cycle<sup>3, 4, 7, 8</sup>. Less common is to change the HRS current by changing the maximum voltage during the reset ( $V_{\max}$ )<sup>8,9</sup>. The latter suffers from larger HRS resistance variability since that depends mostly on the length of the ruptured part of the filament which varies from cycle to cycle. In the former method, the LRS resistance depends on the radius and conductivity of the filament, resulting in a far smaller variability. Therefore, MLC switching by set current is mostly used and 3-bit cells using this method have been reported<sup>3,7</sup>. However, the low resistance states at small set currents, within the lower region of memory window, have poor stability and the device often fails to be set or reset. Here, we have combined the two MLC switching methods using N-doped atomic layer deposited (ALD) Ta<sub>2</sub>O<sub>5</sub> based RRAM devices. This approach facilitates the achievement of more than 3-bit MLC by varying the low resistance states in the upper half of the memory window by set current compliance and high resistance states in the lower part by  $V_{\max}$ . The role of N doping is to control the oxygen density in the film, firstly to widen the memory window to accommodate more intermediate states, secondly, to improve the stability of states, and thirdly, to extend the region of gradual reset to lower currents, which is essential for MLC switching of HRS resistance.

The selection of an appropriate switching layer material is central to the RRAM operation<sup>10</sup> and tantalum oxide is of particular interest because of its high endurance of switching cycles<sup>11</sup>. The purpose of N doping is to reduce electronic defect states associated with the  $O_{\text{vac}}$ <sup>12</sup> and enabling CFs to be formed with more stable and discrete resistances suitable for MLC switching. The improved state stability is believed to be due to the elimination of excess conducting paths in the oxide by passivation of oxygen vacancies by nitrogen atoms<sup>3</sup>. Metal-insulator-metal (MIM) structures were fabricated on Corning glass or silicon wafer substrates. The bottom electrode is 50 nm Pt, with a 10 nm Cr adhesion layer, deposited by dc magnetron sputtering. The oxide layer, 15 nm thick Ta<sub>2</sub>O<sub>5</sub> or N-doped Ta<sub>2</sub>O<sub>5</sub>, was deposited by conventional ALD, using the precursor Ta(OC<sub>2</sub>H<sub>5</sub>)<sub>5</sub> and H<sub>2</sub>O as an oxidant. N doping was achieved by substituting a proportion of the water cycles with aqueous NH<sub>4</sub>OH solution, in the increasing ratios of 1:99, and 1:32, 1:11. The top electrode of 30 nm Ti, with a 60 nm Pt capping layer, was deposited by rf sputtering. Conventional photolithography and metal lift-off were used to define the top and bottom contacts and device active area of 512 overlapping and 512 cross-line square devices, with dimensions of 2 to 150  $\mu\text{m}$ . The samples were annealed in forming gas ambient at 350 °C for one hour. This facilitated the electroforming and

device switching, presumably by increasing the O vacancies reservoir at top metal-oxide interface. A bright field transmission electron microscopy (BF-TEM) micrograph of the MIM structure are shown in Fig. 1. The inset diffraction patterns show that the tantalum oxide film remains amorphous even after forming gas annealing. The dc set-reset cycles were performed by return sweeps of a dc voltage from 0 V to a positive value, applied to the Ti contact, for set and to a negative value for reset. The measurements were taken using either HP 4155A or Agilent B1500A Semiconductor Parameter Analyzers. The LRS resistance was programmed by setting the current compliance during the set cycle and the HRS resistance was programmed by setting  $V_{\max}$ . The resistance at each state was calculated at a 'read' voltage of  $-0.1$  V. The statistical variability of each programmed state was analyzed after performing 30 alternate set and reset cycles and calculating the cumulative distribution function (CDF) of resistance at each state. The same method was used to analyze the device to device variability on 25 devices in different regions of the sample.

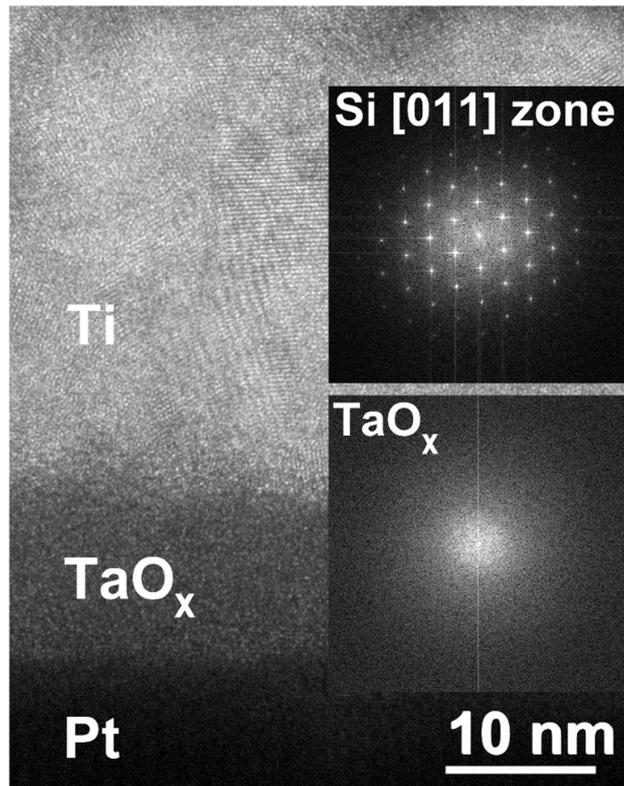


FIG. 1. BF-TEM image of device cross section and diffraction patterns from the TaO<sub>x</sub> film and silicon substrate.

Figure 2 (a) shows set and reset cycles of the devices made of undoped and N-doped Ta<sub>2</sub>O<sub>5</sub> with various doping cycles, at a set current compliance of 1 mA. Electroforming was performed prior to the first reset cycle by a positive sweep up to 10 V to the Ti contact and current compliance of 10  $\mu$ A, as shown in the inset of Fig. 2 (a). The forming voltage was typically 6-8 V and no specific difference in forming voltage or its current-voltage shape was observed between doped and undoped

devices. The switching characteristics are independent of device area ( $4\text{-}22500\ \mu\text{m}^2$ ), which is in agreement with previous reports<sup>11</sup> and indicates the presence of only one dominant conductive filament at a time in a given device. The memory window is significantly increased by the level of N doping, from 17 to 141 at  $-0.1\ \text{V}$  for the 1:11 N doping cycles, due to reduction in HRS current in the N-doped device, which is attributed to a more resistive or longer ruptured region of conductive filament. For all of the doping cycle treatments used, the nitrogen incorporation in the films is at the detection limit of XPS measurements (not shown), implying that the N content is less than 2 %. The sample with 1:12 N cycles has been used as an optimum choice for this study. Excess amount of N doping increases the resistance nonlinearity and variability<sup>3</sup>. The variation in LRS current with N doping is insignificant. The HRS current, however, is controlled by the ruptured region in the conductive filament. The decreased HRS current in N-doped devices can be attributed to a longer ruptured region of the filament and larger hopping distance due to a lower concentration of O vacancies.

The effect of set current compliance and  $V_{\text{max}}$  on 1:12 N-doped devices is shown in Fig. 2 (b) and (c), respectively. Varying the former, changes the LRS current and varying the latter changes the HRS current, both resulting in intermediate states within the upper and lower parts of the memory window, respectively, applicable for MLC switching. The increase in LRS current with increase in set current compliance is related to a larger radius of the conductive filament and also the higher concentration of O vacancies within the filament<sup>4</sup>. The reduction in HRS current with  $V_{\text{max}}$  is due to a larger emigration rate of oxygen vacancies from conductive filament, resulting in a longer ruptured region.

For MLC switching with a large number of intermediate states (3-bit or more), the device not only requires a large memory window to accommodate the states, but also the states need to have small variability to prevent intermixing. It has been shown in the literature<sup>4, 7</sup> and in this work, that atomic doping can improve both requirements. Two types of variability in RRAM devices can be discussed which, although related, are of a different nature. One is variability in resistance of a state during successive set and reset cycles, which depends on the rupture and restoration of a particular filament in the device. The other is device to device variability which depends on the uniformity of O vacancies within the film across the sample. Figure. 3 (a) shows the cumulative distribution function (CDF) of the resistance at 8 various states (3 bits) of the N-doped device programmed by set current compliance, read at voltage of  $-0.1\ \text{V}$ . The devices show good stability in LRS resistance with a large gap between highest LRS and lowest HRS resistances. The same graph for an undoped device is shown in Fig. 3 (b) for comparison. Both N-doped and undoped devices have good stability at LRS, however, the undoped device shows more variability at low set currents, has significantly greater variability in HRS with a much smaller memory window. The device to device variability of the sample is also shown in the CDF plots of Fig. 3 (c) for 25 devices with differing areas and

locations across the sample. These data demonstrate the excellent stability at LRS, but the HRS resistance has a fairly large variability.

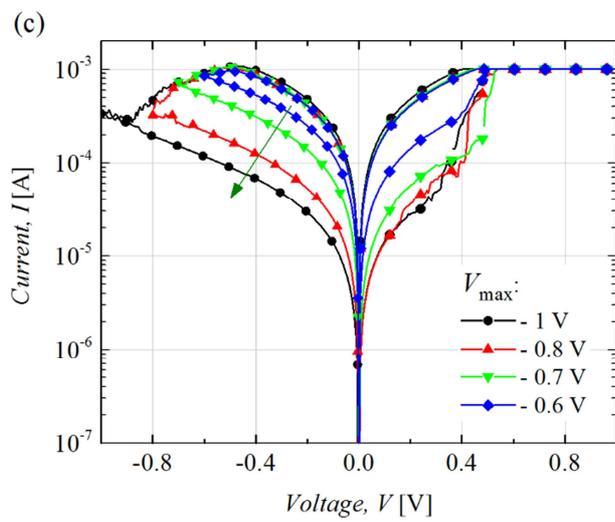
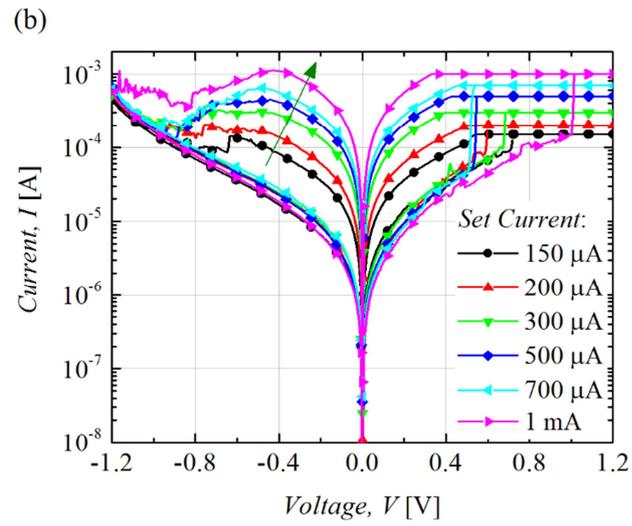
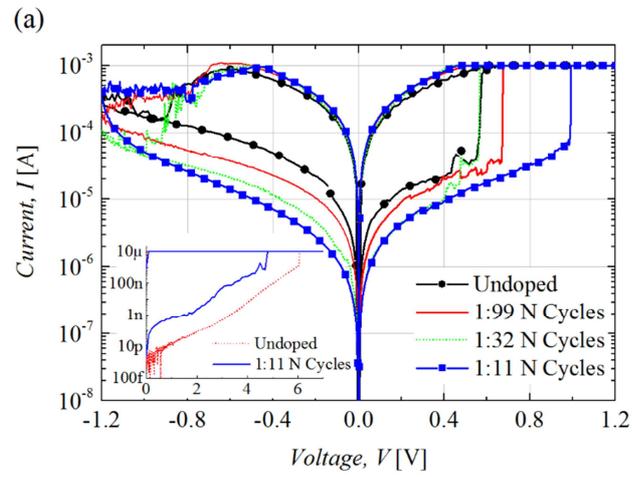


FIG. 2. (a) The switching characteristics of the N-doped and undoped RRAM devices at set current compliance of 1 mA. The forming IV characteristics are shown as an inset. Variation of LRS current with set current compliance (b) and HRS with  $V_{\max}$  (c). The arrows show the direction of change in current with setting parameters.

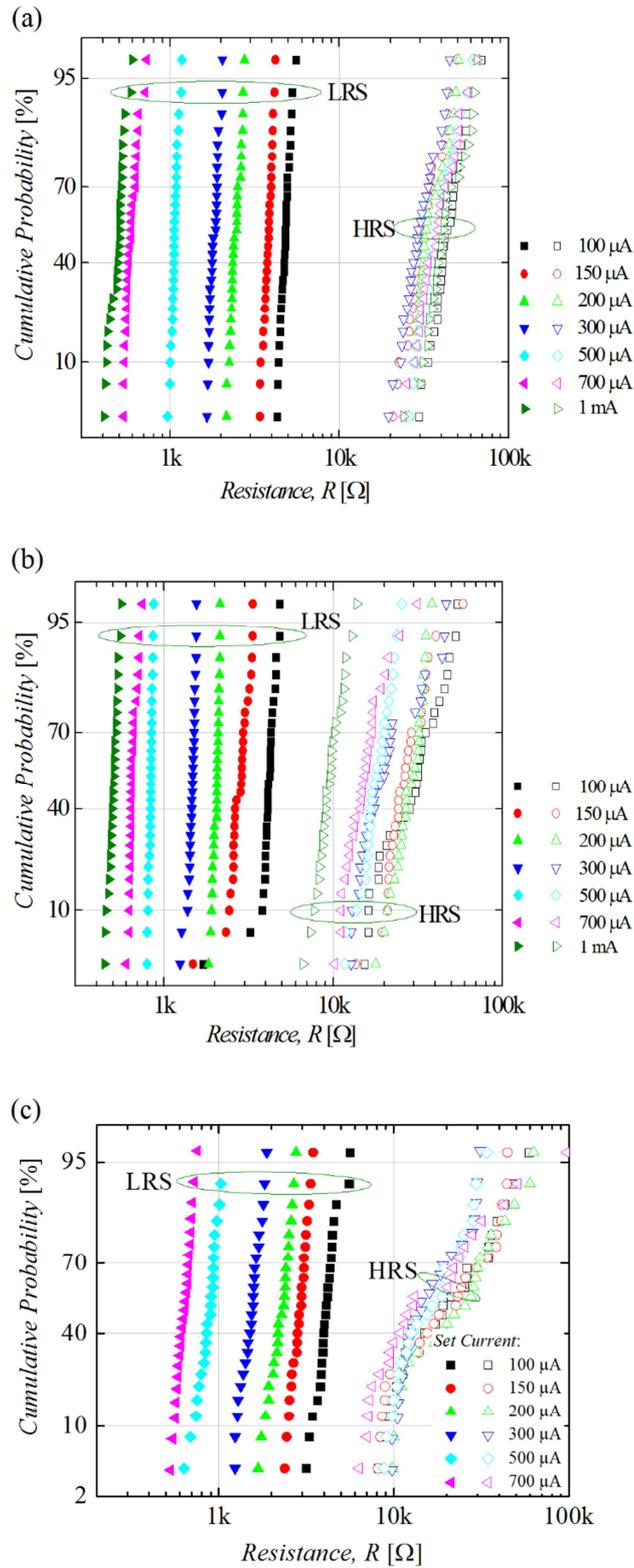


FIG. 3. The cumulative distribution of resistance states for N-doped (a) and undoped (b) devices. (c) The device to device distribution for 25 N-doped devices.

The HRS variability is ascribed to variation in diameter of the critical filament region which affects the length of the ruptured part of a filament and hence the HRS resistance. Evolution of a conductive filament and its rupture and restoration have been studied by a number of analytical and numerical models<sup>13-16</sup> which have shed light on the device switching characteristics and the stability of states. The filament rupture and hence the HRS resistance after reset depends on the electric field distribution in the filament and in particular in the critical filament region, the narrowest part of the filament or the neck of its commonly described hourglass shape. The diameter of a filament in the critical region depends on the initial distribution of O vacancies within the film and their migration by electric field. When it is narrow, the electric field in this region is stronger and more O vacancies are emigrated, resulting in a larger HRS resistance. Local Joule heating due to the LRS current before reset can also make the diameter larger.<sup>13, 14</sup>

The MLC by set current is preferable because of good variability performance of the LRS resistance. Fig. 3 (a) shows 3-bit MLC using only this method where the lower part of memory window remains unused. At very low set currents, the devices fail to switch (set and reset) and even if they do, the stability is very poor. This gap can be filled with high resistive states by changing  $V_{\max}$ . Figure 4 (a) shows the MLC with 6 states using this method. Combining the two MLC methods can provide 3-bit (8 states) or more, with large separation of states, as is shown in Fig. 4 (b). By this method and with further improvement in device stability, 4-bit MLC is feasible, but this is unlikely to be achieved by setting either current or voltage individually.

We have previously reported the properties of  $O_{\text{vac}}$  in the  $\lambda$  phase of  $Ta_2O_5$  crystal<sup>17</sup> and in the amorphous  $Ta_2O_5$  network<sup>18</sup> through sX density functional simulations. It is found that the 2-fold  $O_{\text{vac}}$  in both networks has a small formation energy, low diffusion barrier, and deep defect states below the conduction band, believed to be strongly related to the resistive switching behavior in the RRAM devices. Here we have applied the sX hybrid functional density functional simulation to calculate the effect of N doping in the amorphous  $Ta_2O_5$  network. The simulation settings are the same as in our previous work<sup>18</sup>. The atomic structure and density of states (DOS) of the amorphous  $Ta_2O_5$  network with a 2-fold  $O_{\text{vac}}$  is shown in Fig. 5 (a, b). The  $O_{\text{vac}}$  gives rise to a deep defect state in the mid bandgap region, with the defect orbital localized on one nearby Ta atom. The effect of N doping has been studied by substituting two O atoms with N atoms near the  $O_{\text{vac}}$  site, as shown in Fig. 5 (c). The original  $O_{\text{vac}}$  defect state in the bandgap is removed as a result of N doping<sup>12</sup>, leaving some N non-bonding states near the valence band of  $Ta_2O_5$  (Fig. 5d), as in previous work.<sup>19</sup>

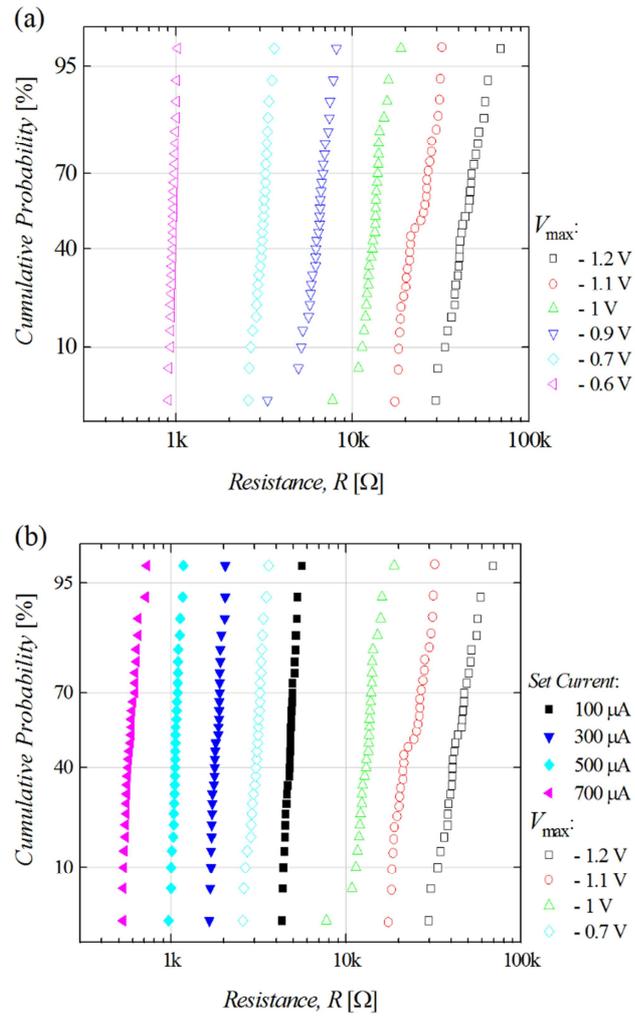


FIG. 4. (a) MLC switching by changing the HRS resistance by varying  $V_{\max}$ . (b) Combination of two methods for 3-bit MLC. The solid and open symbols show LRS and HRS resistances, respectively.

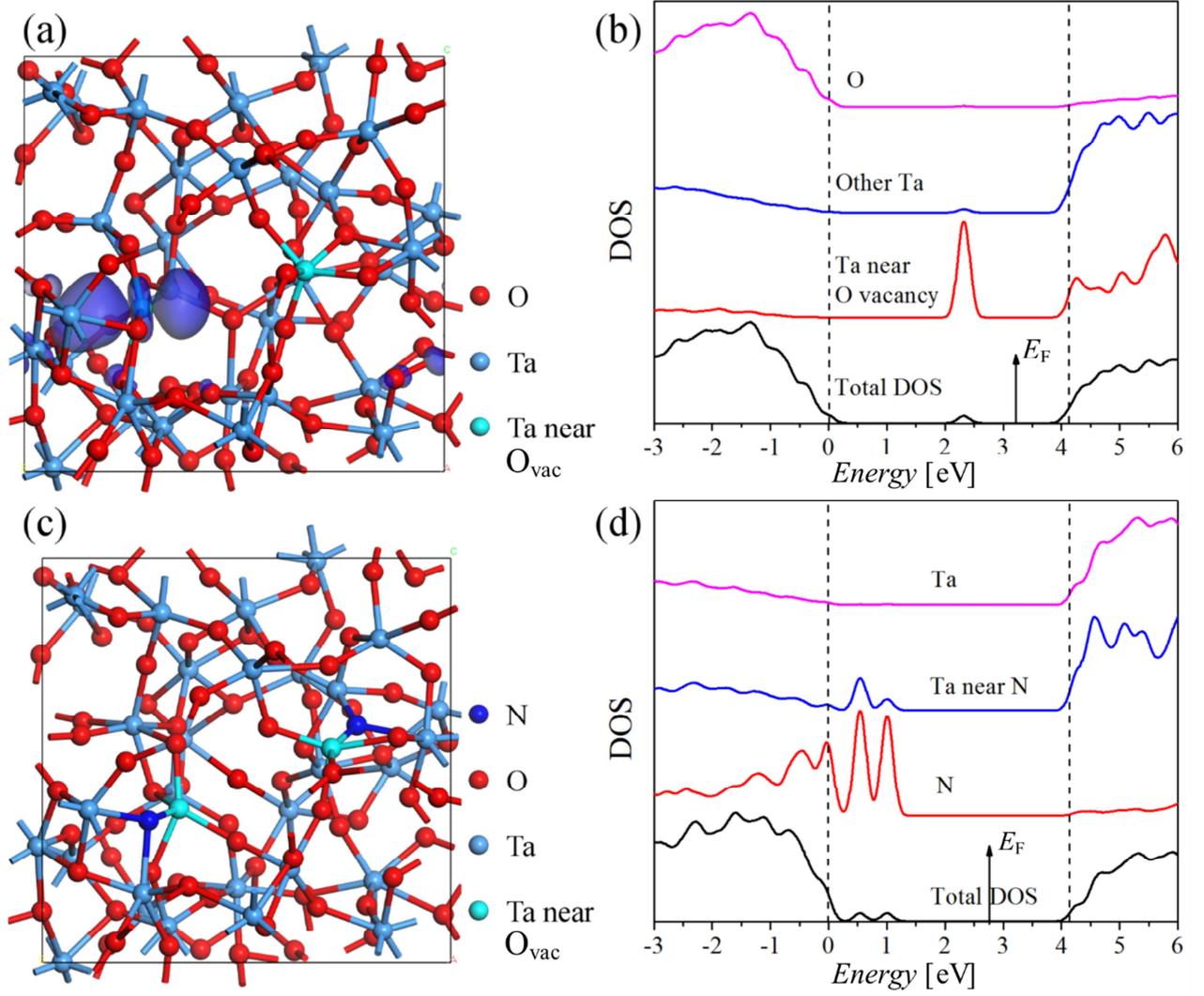


FIG. 5. (a) The atomic model and (b) the density of states of undoped amorphous Ta<sub>2</sub>O<sub>5</sub> structure with an O vacancy. The blue isosurface shows defect orbital in mid bandgap. (c) The atomic model and (d) the density of states of amorphous Ta<sub>2</sub>O<sub>5</sub> structure with two N atoms substituting two O atoms near the O vacancy site. The red, light blue, cyan, and dark blue balls denote O, Ta, Ta near O vacancy, and N atoms, respectively.

We now consider the mechanism that influences the HRS current and increases the resistance window between states due to the N doping process. The original O<sub>vac</sub> introduces two excess electrons localized on the nearby Ta atom. During the doping process, each N atom substitutes one O atom in the network, and creates a valence band hole. The two excess electrons from O<sub>vac</sub> transfer to two nearby N atoms to fill their valence states and form a closed shell electronic configuration. Locally, the O<sub>vac</sub> becomes a V<sup>2+</sup> center, involving a certain lattice distortion, and pushes the original defect state up into conduction band<sup>12</sup>. The energy gain of this process comes from the two electrons falling from mid gap states into valence band holes, which compensates the energy required for local distortion. In this way, the 2-fold O<sub>vac</sub> in the amorphous Ta<sub>2</sub>O<sub>5</sub> network is nullified or cancelled by the presence of two N atoms. The neutralization of the defect states associated with the

$O_{vac}$  eliminates the excess conductive paths<sup>7</sup> and reduces the numbers or densities of potential filaments. This effect has been observed in sputtered  $TaO_x$  films by controlling the oxygen partial pressure during sputtering<sup>4</sup>. It was concluded that the variability of the conductive filament resistance depends on both the filament diameter and the oxygen vacancy density in the filament. The model presented here shows that the effect of adding nitrogen to the ALD tantalum oxide can help to improve these factors and that multilevel switching is achievable in the N-doped film. It reduces the variability in LRS current by preventing the set operation of the device through a different conductive path. It also makes the filament denser which decreases the HRS current due to a larger gap in the ruptured filament at critical filament region.

In conclusion, RRAM devices were fabricated by ALD  $Ta_2O_5$  doped with nitrogen. The dc sweep set-reset measurements show that passivation of  $O_{vac}$  increases the switching memory window and states stability. The stability of states and the separation of states at various set currents enables the device for MLC switching applications. The devices have set and reset voltages of less than 1 V and their variability is improved by doping. Density functional simulation of the atomic structures and calculation of density of states using sX hybrid functional correction show that N dopants substitute O atoms near the  $O_{vac}$ , resulting in nullifying the associated electronic defect states in the mid bandgap of the oxide, and thus eliminates alternative conductive paths.

The work has been funded by the Engineering and Physical Sciences Research Council (EPSRC) UK, project numbers EP/M00662X/1, EP/M009297/1 and EP/M006727/1.

<sup>1</sup>H. Akinaga and H. Shima, Proc. IEEE **98**, 2237 (2010); H. S. P. Wong, H. Y. Lee, S. M. Yu, Y. S. Chen, Y. Wu, P. S. Chen, B. Lee, F. T. Chen, M. J. Tsai, Proc. IEEE **100**, 1951 (2012).

<sup>2</sup>W. Kim, S. Menzel, D. J. Wouters, Y. Guo, J. Robertson, B. Roesgen, R. Waser, and V. Rana, Nanoscale **8**, 17717 (2016); A. Calderoni, S. Sills, C. Cardon, E. Faraoni, N. Ramaswamy, Microelectron. Eng. **147**, 145 (2015).

<sup>3</sup>S. H. Misha, N. Tamanna, J. Woo, S. Lee, J. Song, J. Park, S. Lim, J. Park, and H. Hwang, ECS Solid State Lett. **4**, P25 (2015).

<sup>4</sup>A. Prakash, D. Deleruyelle, J. Song, M. Bocquet, and H. Hwang, App. Phys. Lett. **106**, 233104-1 (2015).

<sup>5</sup>F. Miao, J. P. Strachan, J. J. Yang, M.-X. Zhang, I. Goldfarb, A. C. Torrezan, P. Eschbach, R. D. Kelley, G. Medeiros-Ribeiro, and R. S. Williams, Adv. Mater. **23**, 5633 (2011).

<sup>6</sup>U. Celano, L. Goux, R. Degraeve, A. Fantini, O. Richard, H. Bender, M. Jurczak, and W. Vandervorst, Nano Lett. **15**, 7970 (2015).

- <sup>7</sup>A. Prakash, J. Park, J. Song, J. Woo, E. J. Cha, and H. Hwang, *IEEE Electron Device L.* **36**, 32 (2015).
- <sup>8</sup>M. C. Wu, W. Y. Jang, C. H. Lin, and T. Y. Tseng, *Semicond. Sci. Technol.* **27**, 065010 (2012).
- <sup>9</sup>S. Yu, Yi Wu, and H.-S. P. Wong, *Appl. Phys. Lett.* **98**, 103514-1 (2011).
- <sup>10</sup>Y. Guo and J. Robertson, *Appl. Phys. Lett.* **105**, 223516-1 (2014).
- <sup>11</sup>M. J. Lee, C. B. Lee, D. Lee, S. R. Lee, M. Chang, J. H. Hur, Y. B. Kim, C. J. Kim, D. H. Seo, S. Seo, U. I. Chung, I. K. Yoo, and K. Kim, *Nat. Mater.* **10**, 625 (2011).
- <sup>12</sup>K. Xiong, J. Robertson, and S. J. Clark, *J. Appl. Phys.* **99**, 044105-1 (2006); H. Li, Y. Guo, and J. Robertson, *Appl. Phys. Lett.* **104**, 192904-1 (2014).
- <sup>13</sup>D. Ielmini, *IEEE T. Electron Dev.* **58**, 4309 (2011).
- <sup>14</sup>D. Ielmini, *Semicond. Sci. Technol.* **31**, 063002-1 (2016).
- <sup>15</sup>P. Sun, S. Liu, L. Li, and M. Liu, *J. Semicond.* **35**, 104007-1 (2014).
- <sup>16</sup>S. Kim, S. J. Kim, K. M. Kim, S. R. Lee, M. Chang, E. Cho, Y. B. Kim, C. J. Kim, U. I. Chung, and I. K. Yoo, *Sci. Rep.* **3:1680**, 1 (2013).
- <sup>17</sup>Y. Guo and J. Robertson, *Appl. Phys. Lett.* **104**, 112906-1 (2014).
- <sup>18</sup>Y. Guo and J. Robertson, *Microelectron. Eng.* **147**, 254 (2015).
- <sup>19</sup>G. Shang, P. W. Peacock, J. Robertson, *Appl. Phys. Lett.* **84**, 106 (2004).