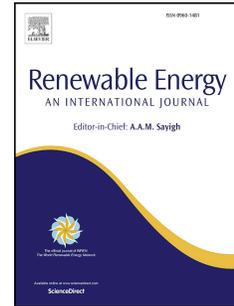


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Modelling impacts of tidal stream turbines on surface waves

Xiaorong Li^a, Ming Li^b, Stuart McLelland^c, Laura-Beth Jordan^c, Laurent Amoudry^d, Qingyang Song^e, Andy Plater^a

^a*School of Environmental Sciences, University of Liverpool, Liverpool, L69 7ZT, U.K.*

^b*School of Engineering, University of Liverpool, Liverpool, L69 3GQ, U.K.*

^c*School of Environmental Sciences, University of Hull, Cottingham Road, Hull, HU6 7RX*

^d*National Oceanography Centre, Joseph Proudman Building, 6 Brownlow Street, Liverpool, L3 5DA*

^e*FTZ-Westkste/Coastal Research Laboratory, Christian-Albrechts-Universitt zu Kiel, Kiel, Germany*

Abstract

A high resolution Computational Flow Dynamics (CFD) numerical model is built based on a laboratory experiment in this research to study impacts of tidal turbines on surface wave dynamics. A reduction of $\sim 3\%$ in wave height is observed under the influence of a standalone turbine located 0.4 m from the free surface. The artificial wave energy dissipation routine ‘OBSTACLE’ within FVCOM is shown to effectively capture the correct level of wave height reduction, reproducing the CFD results with significantly less computational effort.

The turbine simulation system is then applied to a series of test cases to investigate impact of a standalone turbine on bed shear stress. Results suggest an apparent increase in bed stress ($\sim 7\%$) upstream of the turbine due to the inclusion of surface waves. However, in the immediate wake of the turbine, bed stress is dominated by the presence of the turbine itself,

accounting for a $\sim 50\%$ increase, with waves having a seemingly negligible effect up to 9D downstream of the turbine. Beyond this point, the effect of waves on bed shear stress become apparent again. The influence of OBSTACLE on bed stress is also noticeable in the far wake, showing a reduction of $\sim 2\%$ in wave height.

Keywords: Tidal stream energy, Oceanographic model, Wave-current coupling, Bottom shear stress

1 Nomenclature

- 2 \bar{P} The time-averaged static pressure
- 3 \bar{u}_i ($\bar{u}, \bar{v}, \bar{w}$) The time-averaged water velocities in the x_i (x, y, z) directions
- 4 δ_{ij} The Kronecker delta
- 5 μ The molecular viscosity
- 6 ρ The water density
- 7 σ The relative frequency
- 8 θ The wave direction
- 9 \vec{C}_g The group velocity vector
- 10 \vec{V} The ambient water current vector
- 11 C_σ The wave propagation velocity in frequency space
- 12 C_θ The wave propagation velocity in directional space

13	c_d	The drag coefficient
14	c_L	The lift coefficient
15	d	The water depth
16	f	The Coriolis parameter
17	f_d	The drag force
18	F_i	The external body forces in the i directions (x, y, z)
19	f_L	The lift force
20	F_u	The horizontal momentum term in the x direction
21	F_v	The horizontal momentum term in the y direction
22	H	The wave height
23	K_m	The vertical eddy viscosity coefficient
24	K_t	The wave energy transmission coefficient of OBSTACLE
25	L	The wave length
26	N	The wave action density spectrum
27	N_b	The number of blades
28	P_a	The air pressure at sea surface
29	P_H	The hydrostatic pressure
30	q	The non-hydrostatic pressure

31	S_{tot}	The source-sink terms
32	t	Time
33	u	The velocity component in the x direction
34	U_r	The Ursell number
35	v	The velocity component in the y direction
36	V_{tot}	The fluid velocity relative to the blade
37	w	The velocity component in the z direction
38	x	The east axis in the Cartesian coordinate system
39	y	The north axis in the Cartesian coordinate system
40	z	The vertical axis in the Cartesian coordinate system
41	u'_i (u', v', w')	The fluctuating water velocities in the x_i (x, y, z) directions
42	BBL	The Bottom Boundary Layer module
43	BEM	The Blade Element Method
44	CFD	Computational Flow Dynamics
45	FVCOM	The Unstructured Grid Finite Volume Community Ocean Model
46	HATT	Horizontal Axis Tidal Turbine
47	RANS	The Reynolds-averaged Navier-Stokes equations
48	ROMS	Regional Ocean Modelling System

- 49 SWAN Simulating Waves Nearshore
- 50 TbM (BBL) A TbM case with bottom shear stress calculated through BBL,
51 otherwise bottom shear stress is calculated through Equations de-
52 scribed in section 2
- 53 TNO Wave-current FVCOM case without obstacle (for model verification)
- 54 TNO15 Wave-current FVCOM case without obstacle (for impact identifica-
55 tion)
- 56 TSR Tip Speed Ratio
- 57 TYO Wave-current FVCOM case with obstacle activated at the turbine
58 location (for model verification)
- 59 TYO15 Wave-current FVCOM case with obstacle activated at the turbine
60 location (for impact identification)
- 61 VBM The Virtual Blade Model
- 62 VOF The Volume of Fluid method

63 **1. Introduction**

64 As a very promising clean, non-carbon alternative to traditional fossil
65 fuels, tidal stream energy has been gaining significant attention. However,
66 despite the growing interest in this sector of renewable energy, our under-
67 standing of the impacts of tidal stream energy devices on the surrounding
68 environment is still limited, largely due to the lack of data collected from
69 on-site projects.

70 Alternatively, laboratory experiments and numerical simulations are widely
71 adopted to investigate such impacts. For example, porous actuator disc sim-
72 ulators [1, 2, 3] and down-scaled turbine prototype models [4, 5] have been
73 used in laboratories to study turbine-caused impacts on passing flows and
74 turbulence. Also, [6] carried out laboratory experiments to study changes of
75 wake recovery of a turbine subjected to opposing waves. As a complement
76 to laboratory experiments, Computational Flow Dynamics (CFD) modelling
77 is also commonly applied. Similarly, works with turbines approximated as
78 porous discs [7, 8, 9] and with realistic turbine geometry resolved in the com-
79 putational mesh [10, 11, 12] have been published to reveal how flow patterns
80 and turbulent mixing are changed by the turbine in near-field scale.

81 To study the far-field hydrodynamic changes caused by the operation
82 of turbines and turbine arrays, numerical oceanographic models, such as
83 Regional Ocean Modelling System (ROMS) [13] and The Unstructured Grid
84 Finite Volume Community Ocean Model (FVCOM) [14], have also been used.
85 Modifications have been made to such models in order to simulate the effect
86 of tidal stream turbines on the flow motion. These modifications are mostly
87 based on either the additional bottom friction approach [15, 16, 17] or the
88 turbine-induced body force concept [18, 19, 20, 21, 22, 23, 24].

89 In an effort to account for turbine-caused impacts on turbulence in large
90 scale oceanographic models, [25] added three terms to the $k-\epsilon$ closure within
91 ROMS to model turbine related turbulence generation, dissipation and tur-
92 bulence length-scale interference. These three terms were later adapted ac-
93 cordingly to accommodate the theory around which the MY-2.5 turbulence
94 closure is based and applied in FVCOM by [26].

95 In terms of interactions between surface waves and tidal turbines, current
96 research focus has been mainly put on the impact of waves on the performance
97 of turbines due to its immediate industry relevance [27, 28, 29, 30, 31, 32, 33].
98 However, there is a lack of emphasis on the effects of turbines on surface waves
99 in both physical experimental studies and numerical modelling. Because tidal
100 turbines are normally expected to be installed in relatively shallow coastal
101 waters due to difficulties in device installation and operation that would oc-
102 cur otherwise [2], they are likely to have a close proximity to the free surface
103 and hence interfere with the propagation of surface waves. Also, the altered
104 three-dimensional flow structure due to the presence of tidal turbines could
105 also have influence on surface waves through wave-current interaction mech-
106 anisms. Surface waves, particularly in shallow coastal areas, can influence
107 sediment transport dynamics significantly. For instance, vertical mixing in
108 the water column due to wave activities can keep sediment in suspension for
109 longer, inhibiting sediment deposition in the downstream areas of the turbine
110 [34]. Also, wave actions can increase bottom shear stress, leading to enhanced
111 sediment resuspension and erosion [35]. Further, through wave-current in-
112 teractions, waves can drive longshore currents, contributing to long-term
113 shoreline evolution [36, 37]. Therefore, changes in wave dynamics caused by
114 tidal turbines are of high importance in terms of fully understanding impact
115 of tidal turbines on local and regional geomorphology.

116 Due to the aforementioned interactions, the primary objectives of the
117 work documented in this paper are to first explore the potential impacts
118 of tidal turbines on surface waves with the help of high resolution CFD
119 simulations, and second, to develop a Horizontal Axis Tidal Turbine (HATT)

120 simulation system that could implement the impacts of tidal stream turbines
121 on surface waves with a realistic spatial scale.

122 This paper details one high resolution CFD model for tidal turbine im-
123 pact assessment on surface waves. Understandings obtained from the CFD
124 modelling then advise turbine parameterization in large scale oceanographic
125 models. The high resolution modelling is based on a CFD solver — AN-
126 SYS FLUENT. The implementation of effects of turbine operation on sur-
127 face waves is an extension of the turbine simulation platform reported in
128 [26], which parameterized tidal turbines in the current and turbulence clo-
129 sure modules of FVCOM. Impacts of tidal turbines on surface waves are
130 considered in this new model by modification of wave energy flux across the
131 device. A thorough validation study is also presented in which the turbine
132 representation and operation in the CFD models is validated against labora-
133 tory data collected from an experiment conducted at the University of Hull
134 using their ‘Environment Simulator Laboratory Flume’ [5] and the FVCOM
135 model is verified utilizing the CFD simulated results.

136 The structure of the paper is provided as follows for clarity. Firstly in
137 Section 2 ANSYS FLUENT and the FVCOM model are introduced. The in-
138 tegration of turbine simulation within these two frameworks is also discussed
139 in this section. Next, Section 3 introduces the exploratory CFD models which
140 aim to reveal the impacts of turbines on surface waves. A set of experimental
141 data was used for CFD model validation in this section. Section 4 details the
142 verification study for the turbine implementation in FVCOM which considers
143 surface waves. Note that as the experimental data available was considered
144 insufficient for comprehensive validation, verification in this section is based

145 on data generated via the CFD modelling detailed in Section 3. In Section 5,
 146 the turbine simulation system developed based on FVCOM is applied to test
 147 cases in order to reveal impacts of a standalone turbine on its surroundings
 148 which incorporate wave-current interaction processes. A set of discussion is
 149 presented in Section 6, followed by concluding remarks given in Section 7 to
 150 summarise important results from sections 4 and 5, along with suggestions
 151 for potential future developments.

152 2. Modelling system

153 2.1. ANSYS FLUENT — a CFD solver

154 FLUENT solves the three-dimensional Reynolds-averaged Navier-Stokes
 155 (RANS) equations which can be written in tensor form as follows:

$$\frac{\partial \rho}{\partial t} + \frac{\partial \rho \bar{u}_i}{\partial x_i} = 0 \quad (1)$$

$$\frac{\partial(\rho \bar{u}_i)}{\partial t} + \frac{\partial(\rho \bar{u}_i \bar{u}_j)}{\partial x_j} = -\frac{\partial \bar{P}}{\partial x_i} + \frac{\partial}{\partial x_j} \left[\mu \left(\frac{\partial \bar{u}_i}{\partial x_j} + \frac{\partial \bar{u}_j}{\partial x_i} \right) - \frac{2}{3} \mu \frac{\partial u_j}{\partial x_i} \delta_{ij} \right] + \frac{\partial}{\partial x_j} (-\rho \overline{u_i' u_j'}) + F_i \quad (2)$$

157 where ρ is the water density; t is time; μ is the molecular viscosity; δ_{ij} is the
 158 Kronecker delta and F_i are external body forces in the i directions (x, y, z).
 159 \bar{u}_i ($\bar{u}, \bar{v}, \bar{w}$) and u_i' (u', v', w') are the time-averaged (mean) and fluctuating
 160 water velocities in the x_i (x, y, z) directions, respectively. The combination
 161 of these two velocity components forms the instantaneous (exact) velocities:

$$u_i = \bar{u}_i + u_i' \quad (3)$$

162 Likewise, \bar{P} is the time-averaged static pressure and for all scalar vari-
 163 ables:

$$\phi = \bar{\phi} + \phi' \quad (4)$$

164 where ϕ denotes a scalar quantity such as pressure and $\bar{\phi}$ and ϕ' are the mean
 165 and fluctuating components of a scalar variable.

166 The Reynolds stress terms, $-\overline{\rho u_i' u_j'}$, which appear on the right hand side
 167 of Equation 2 represent the effects of turbulence and are modelled based
 168 on the Shear Stress Transport (*SST*) $k - \omega$ turbulence closure [38] in this
 169 research.

170 To simulate the wind-wave-induced free surface effects, the Volume of
 171 Fluid (VOF) method is used in FLUENT. The formulation of the VOF model
 172 relies on the fact that the modelled phases are not immiscible. It calculates
 173 the fractions (α_i , $0 < \alpha_i < 1$) of the simulated phases (water and air in
 174 the present research) in each computational cell and in each control volume.
 175 The volume fractions of all phases sum to unity. Based on the local value of
 176 α_i , the appropriate properties and variables will be assigned to each control
 177 volume within the domain.

178 A single momentum equation which is dependent on the volume fractions
 179 of all phases through the properties ρ and μ is solved throughout the calcu-
 180 lation domain, and the computed velocity field is shared among the phases.
 181 The momentum equation is given by

$$\frac{\partial}{\partial t}(\rho \vec{v}) + \nabla \cdot (\rho \vec{v} \vec{v}) = -\nabla p + \nabla \cdot [\mu(\nabla \vec{v} + \nabla \vec{v}^T)] + \rho \vec{g} + \vec{F} \quad (5)$$

182 where ρ is the volume-fraction-averaged density $\rho = \sum \alpha_i \rho_i$ and μ the
 183 volume-fraction-averaged viscosity calculated in the same manner.

184 A continuity equation for the volume fraction of one (or more) of the
 185 phases helps to track the interface(s) between the phases. For the i^{th} phase,

186 this equation takes the form of the following:

$$\frac{\partial \alpha_i}{\partial t} + \vec{v} \cdot \nabla \alpha_i = 0 \quad (6)$$

187 Additional scalar equations, such as those solving turbulence quantities,
 188 are also processed applying the shared-fields approach; i.e. only a single/a
 189 single set of transport equations is solved and the variables (e.g., k and ω)
 190 are shared by the phases throughout the domain.

191 A wave boundary condition is applied to the velocity inlet of the VOF
 192 model to enable the simulation of wave propagation. FLUENT provides
 193 a good variety of wave theories such as first order linear wave theory and
 194 second/higher order Stokes wave theories. The choice of wave theory is
 195 made based on Ursell number ($U_r = \frac{HL^2}{d^3}$) and wave steepness (H/L), where
 196 H , L and d are wave height, wave length and water depth, respectively.
 197 Linear wave theory is suitable when $U_r < 40$, given $H/L < 0.04$ and sec-
 198 ond/higher order Stokes wave theories are more appropriate when $U_r < 40$
 199 and $H/L > 0.04$ [39]. The wave theories are fully coupled with the continuity
 200 and momentum equations of FLUENT. Details of the wave theories and the
 201 wave-current coupling can be found in [38, 40].

202 2.2. Representation of HATT in FLUENT

203 The Virtual Blade Model (VBM) is adopted in this research to simulate
 204 HATT in FLUENT. In VBM, the actual blades are not directly present.
 205 Instead, the rotor is simulated inside a rotor disk fluid zone across which the
 206 virtual blades swipe. The virtual blades are achieved through adding a body
 207 force in the x , y and z directions. This method is an application of a built-in
 208 blade simulating scheme — Blade Element Method (BEM) — within ANSYS

209 FLUENT. In BEM, each blade is divided into small sections from root to tip.
 210 The lift and drag forces exerted on each segment are calculated based on the
 211 blade design as well as the lift and drag coefficients of each section:

$$f_{L,D} = c_{L,D} \cdot c(r/R) \cdot \frac{\rho \cdot V_{tot}^2}{2} \quad (7)$$

212 where $c_{L,D}$ is lift/drag coefficient specified by the user; $c(r/R)$ is the chord
 213 length; ρ is the fluid density and V_{tot} is the fluid velocity relative to the blade.

214 The lift and drag forces are then averaged over a full turbine rotation to
 215 calculate the force on each cell in the discretized domain:

$$F_{L,D_{cell}} = N_b \cdot \frac{dr \cdot d\theta}{2\pi} \cdot f_{L,D} \quad (8)$$

216

$$\vec{S}_{cell} = -\frac{\vec{F}_{cell}}{V_{cell}} \quad (9)$$

217 where N_b is the number of blades and V_{cell} is the volume of a grid cell.

218 2.3. Three-dimensional FVCOM

219 To model the impacts of tidal stream energy devices on coastal regions,
 220 FVCOM, which is a three-dimensional, free surface, terrain-following oceanographic
 221 model [14], is used in this research. The momentum and continuity
 222 equations of FVCOM are presented in Equations 10-13. FVCOM includes
 223 fully coupled wave-current-sediment modules and, therefore, is particularly
 224 useful for modelling coastal processes. Also, it uses an unstructured triangular
 225 mesh to discretize computational domains horizontally, which allows
 226 for high resolution around individual turbines whilst maintaining a smooth
 227 transition to a relatively large mesh size far from the turbines. Such a treatment
 228 of spatial discretization provides a good balance between accuracy and

229 computational effort.

$$\frac{\partial u}{\partial t} + u \frac{\partial u}{\partial x} + v \frac{\partial u}{\partial y} + w \frac{\partial u}{\partial z} - f v = -\frac{1}{\rho} \frac{\partial (P_H + P_a)}{\partial x} - \frac{1}{\rho} \frac{\partial q}{\partial x} + \frac{\partial}{\partial z} \left(K_m \frac{\partial u}{\partial z} \right) + F_u \quad (10)$$

$$\frac{\partial v}{\partial t} + u \frac{\partial v}{\partial x} + v \frac{\partial v}{\partial y} + w \frac{\partial v}{\partial z} + f u = -\frac{1}{\rho} \frac{\partial (P_H + P_a)}{\partial x} - \frac{1}{\rho} \frac{\partial q}{\partial y} + \frac{\partial}{\partial z} \left(K_m \frac{\partial v}{\partial z} \right) + F_v \quad (11)$$

$$\frac{\partial w}{\partial t} + u \frac{\partial w}{\partial x} + v \frac{\partial w}{\partial y} + w \frac{\partial w}{\partial z} = -\frac{1}{\rho} \frac{\partial q}{\partial z} + \frac{\partial}{\partial z} \left(K_m \frac{\partial w}{\partial z} \right) \quad (12)$$

$$\frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} + \frac{\partial w}{\partial z} = 0 \quad (13)$$

233 where x , y , and z are the east, north, and vertical axes in the Cartesian
 234 coordinate system; u , v , and w are the three velocity components in the x ,
 235 y , and z directions respectively; P_a is the air pressure at sea surface; P_H is
 236 the hydrostatic pressure; q is the non-hydrostatic pressure; f is the Coriolis
 237 parameter and K_m is the vertical eddy viscosity coefficient. F_u , F_v represent
 238 horizontal momentum terms.

239 Extensive work has been done by the authors to enable the prediction of
 240 complete three-dimensional velocity profiles and mixing in the wake of tur-
 241 bines by making modifications to the current and turbulence closure modules
 242 of FVCOM [26]. The current research further extends the turbine simula-
 243 tion platform reported in [26] in terms of proposing a way to incorporate the
 244 effects of turbines on surface waves in the model.

245 For completeness, the basic theory surrounding surface waves and wave-
 246 current coupling in FVCOM is given as follows. More details of the model
 247 can be found in [41].

248 To simulate surface wave propagation, Simulating Waves Nearshore (SWAN)
 249 [42] is integrated with FVCOM. The governing equation of the wave action

250 density spectrum is given as:

$$\frac{\partial N}{\partial t} + \nabla \cdot \left[\left(\vec{C}_g + \vec{V} \right) N \right] + \frac{\partial C_\sigma N}{\partial \sigma} + \frac{\partial C_\theta N}{\partial \theta} = \frac{S_{tot}}{\sigma} \quad (14)$$

251 where N is the wave action density spectrum, \vec{C}_g is the group velocity vector,
 252 \vec{V} is the ambient water current vector, σ is the relative frequency, θ is the wave
 253 direction, C_σ and C_θ are the wave propagation velocities in the frequency
 254 domain and directional space respectively and S_{tot} is the source-sink term
 255 considering wind-induced wave growth, nonlinear transfer of wave energy due
 256 to three-wave interactions, nonlinear transfer of wave energy due to four-wave
 257 interactions, wave decay due to white capping, wave decay due to bottom
 258 friction and wave decay due to depth-induced wave breaking. More details
 259 are available in the SWAN technical manual [42].

260 Due to the presence of surface waves, the bottom boundary layer is af-
 261 fected and the shear stress is much higher than that due to current alone
 262 [35]. To take this into account, a special treatment is needed close to the
 263 bed, which is implemented in the bottom boundary layer module (BBL).
 264 BBL calculates the bottom shear stresses under the condition of combined
 265 waves and currents. The calculation of bottom shear stress is important as
 266 it influences the flow field as well as sediment transport patterns. The BBL
 267 module developed by [43] based on the theory proposed by [44] was con-
 268 verted into an unstructured-grid finite-volume version and implemented in
 269 FVCOM. It is, hence, used in the present research. Details of BBL can be
 270 found in [43].

271 FVCOM includes a wave-current-sediment fully coupled system. After
 272 initialization, the wave module starts to solve the wave dynamics, providing
 273 information of surface waves. The interactions between the current and wave

274 modules are achieved through radiation stress terms according to Mellor's
275 theory [45, 46, 47]. Results from the current module, velocities and surface
276 elevation in particular, provide the wave module feedback for the next time
277 step calculation. Results from the current and wave modules are then sent
278 to the BBL module to calculate the bottom stresses under the combined
279 influence of waves and current. These stresses are then used to solve the
280 momentum equations.

281 *2.4. Representation of HATT in FVCOM*

282 As will be demonstrated by CFD experiments in Section 3, surface wave
283 height is affected by the inclusion of turbines. To represent this effect, one
284 of the built-in features of SWAN — “OBSTACLE” is applied in the present
285 study. The OBSTACLE routine absorbs wave energy along a finite line
286 (defined between two locations) and dissipates it according to a constant
287 transmission coefficient K_t . A detailed implementation of the OBSTACLE
288 routine in this context can be found in [48].

289 To model the effect of turbines on waves, the OBSTACLE energy absorp-
290 tion line length in the model is set to the diameter of the simulated turbine.
291 Note however that the impact of the line length upon the simulation is not
292 continuous, as it absorbs energy only where it intersects with the mesh. In
293 other words, two energy absorption lines of different length but with ends
294 lying in the same respective triangle segments would have equal effect. The
295 line is positioned in a way that it passes through the centre and crosses two
296 sides of the triangles selected to house the turbine (see Figure 1). It should
297 be pointed out that the turbine parameterization in the current and turbu-
298 lence closure modules of FVCOM reported in [26] are utilized in this research

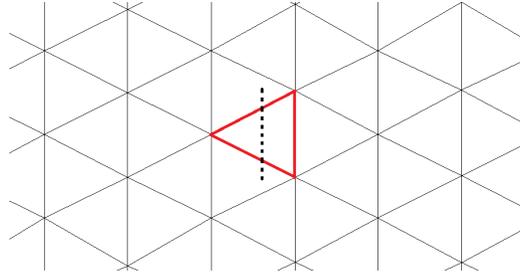


Figure 1: Illustration of the turbine position in the x-y plane on the mesh. The red triangle indicates the mesh element in which the turbine is implemented. The black dotted line illustrates the application of OBSTACLE.

299 when a turbine is present.

300 3. The CFD model

301 A CFD model is built in this research to study the impacts of tidal tur-
 302 bines on surface waves. It is based on an experiment carried out at the
 303 University of Hull using their ‘Environment Simulator Laboratory Flume’
 304 [5]. The flume is 11 m in length, 1.6 m wide and 0.8 m deep. The water
 305 depth was 0.6m throughout the experiment. The flow rate at the inlet was
 306 0.3 m/s. A surface wave propagating in the direction of the flow was imposed
 307 upon the inlet. The wave height and wave period were 0.15 m and 1 s, re-
 308 spectively. A horizontal axis rotor with a diameter of 0.2 m was located 0.2
 309 m above the bed and the tip speed ratio (TSR) of the rotor was constantly
 310 5.5. Measurements of velocity were taken along the centreline from 1D to 4D
 311 downstream of the rotor (where D is the turbine diameter).

312 Although a wide range of data was collected, the measurements did not
 313 include free surface variations which are the main focus of this research.
 314 Therefore, a CFD model replicating the experimental conditions was set up

315 to capture the impacts of the rotor on surface waves. The CFD model was
316 validated by recreating the conditions of the experiments for which measure-
317 ments were available.

318 In the CFD model, the flume length was, instead of 11m, 3.1 m for ease
319 of simulation. The velocity at the inlet was 0.3 m/s. A following wave with
320 wave height of 0.15 m and wave period of 1 s was imposed at the inlet. The
321 computation of wave propagation is based on the 2nd-order wave theory. To
322 reduce the wave energy being reflected back into the flume from the exit,
323 three porous zones, with thickness of 0.2m, 0.2m and 0.1m, were set at the
324 outlet boundary, with porosity declining from 0.95 to 0.9 to 0.8. Essential
325 configurations of VBM, i.e. geometrical setup and running parameters of the
326 rotor are specified according to [49].

327 Figure 2 compares the ensemble average of stream-wise flow velocity pro-
328 files predicted by the CFD model against that measured in the laboratory at
329 1D, 2D, 3D and 4D downstream of the rotor. It should be noted that there
330 are overlaps in the measured profiles. This is because in the laboratory, the
331 centreline slice on which the velocities were measured was divided into 9 sub-
332 slices and each of these sub-slices overlaps with its neighbour sub-slices. The
333 overlaps provide a way to ensure the sub-slices are aligned correctly.

334 It can be seen from Figure 2 that the computed velocity profiles at all 4
335 locations agree well with the measurements at the rotor swiping layers with
336 the exception of location 1D specifically above the rotor hub. This is due
337 to the fact that the rotor housing and supporting structure (suspending the
338 turbine from above) in the laboratory flume interfere with the flow at 1D.
339 As these additional structures are not accounted for in the model, the result

Table 1: NSME for the CFD case against the experimental data

1D	2D	3D	4D
0.88	0.93	0.91	0.91

340 differs in this area. Further, the velocities in the region below the rotor are
 341 over-estimates. This over-estimation is likely due to a slightly over-predicted
 342 near bed wave boundary layer effect. To quantify the agreement between the
 343 predictions and measurements, the Nash Sutcliffe Model Efficiency (NSME)
 344 is calculated based on Equation 15 for each location for the rotor swiping
 345 layers and provided in Table 1. The NSME has been widely used to quantify
 346 the accuracy of model prediction, and the model performance is considered
 347 as excellent for NSME in between 0.65-1, very good for 0.65-0.5, good for
 348 0.5-0.2, and poor for less than 0.2 (e.g. [50, 51, 52]). Therefore, the agree-
 349 ment between FLUENT based CFD model results and measured data are
 350 considered to be satisfactory at all sites.

$$NSME = 1 - \frac{\sum_{i=1}^n (q_i - q_{iest})^2}{\sum_{i=1}^n (q_i - \bar{q})^2} \quad (15)$$

351 where n is the number of records in the validation data; q_i is the validation
 352 data; q_{iest} is the calculated result; \bar{q} is the average of the validation data.

353 After being validated, the CFD model predicted free surfaces are studied
 354 to investigate the impacts of tidal turbines on surface waves. For this purpose,
 355 an undisturbed case (i.e. no turbine) was run to provide baseline surface
 356 wave profiles. The computed free surfaces at the two time instants when the
 357 trough and peak pass the turbine location are presented in Figure 3 (A) and
 358 3 (B) respectively. It can be seen from Figure 3 that the inclusion of the

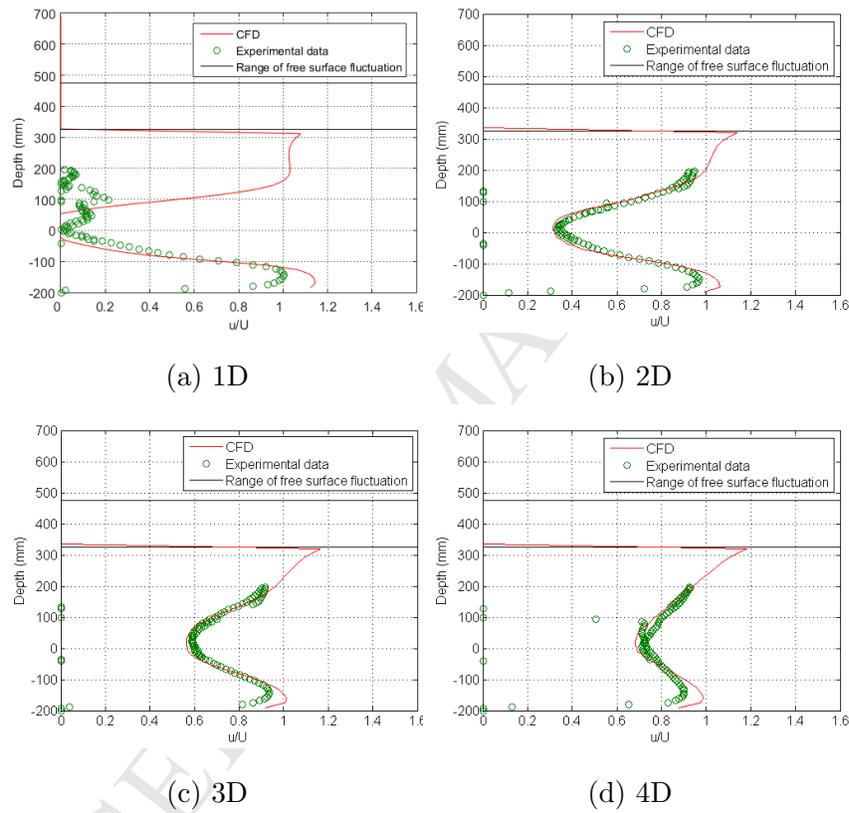


Figure 2: Normalized velocity profiles of the wave-current CFD case against those measured in the laboratory at 1D, 2D, 3D and 4D downstream of the rotor.

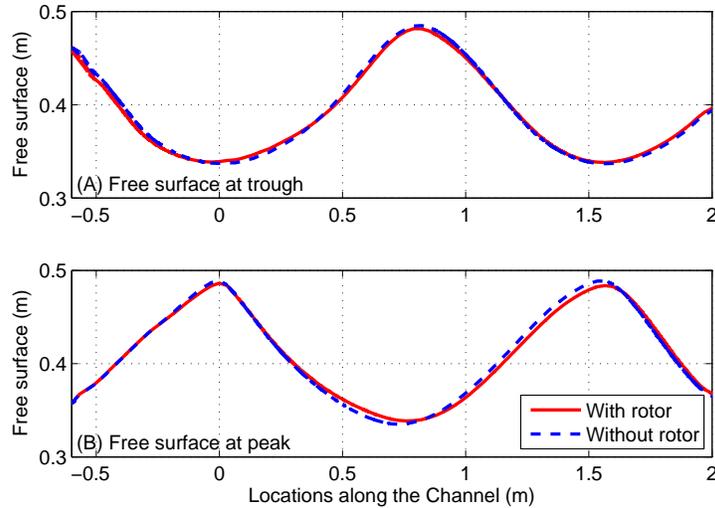


Figure 3: CFD predicted free surfaces at the wave trough (A) and peak (B) with and without the rotor. The rotor is positioned at 0 m along the channel.

359 rotor reduces the wave height; The wave height drops by $\sim 2.5\%$ when the
 360 rotor is present. It is also observed from Figure 3 that the wave length is
 361 increased due to the inclusion of the rotor.

362 The deformation of surface waves observed above, i.e. wave height drop
 363 and wave length increase, is likely to be caused by wave-current interactions.
 364 The obstruction effect of the rotor in motion forces the passing water to flow
 365 around the device, causing the velocity near the free surface to be increased.
 366 The accelerated flow at the surface results in a faster transport of wave energy
 367 and, consequently, reduced wave height and increased wave length.

368 4. Verification of the FVCOM model

369 This section explores the possibility of using the OBSTACLE mentioned
 370 above to represent the observed rotor-caused wave height drop. Hence, a

371 FVCOM based model was set up according to the above-mentioned experi-
372 mental conditions. The mesh of the model has a uniform spatial resolution of
373 0.2 m (i.e. 1D) throughout the computational domain. Vertically, the water
374 column is evenly divided into 50 sigma layers to accommodate the turbine
375 representation in the current and turbulence modules recorded in [26].

376 The turbine effects on surface wave propagation is represented by sub-
377 tracting a certain amount of energy from the energy conservation equation
378 (Equation 14) as discussed in Section 2.4. In particular, the wave energy
379 transmission coefficient K_t needs to be estimated. For this purpose, three
380 cases are tested: baseline case where turbine is absent and the hydrodynam-
381 ics resemble those of the undisturbed experimental conditions, case TNO
382 where the turbine is present but OBSTACLE is deactivated, and case TYO
383 where both the turbine and OBSTACLE are implemented. In case TYO, the
384 wave energy transmission coefficient of OBSTACLE, K_t , is 0.98.

385 To verify the choice of K_t , Figure 4 compares the drop of wave height in
386 percentage along the channel of the two FVCOM cases, TYO and TNO, and
387 that of one of the CFD models (rotor positioned at 0.2 m above the bed).
388 Wave height drop in percentage (hereafter wave height drop) is defined as the
389 ratio between the decrease in wave height and the background wave height.
390 It is obvious that the wave height drop at the turbine location predicted by
391 TNO is $\sim 1.0\%$ less than that predicted by the corresponding CFD case. This
392 difference is quite significant given that the correct drop is $\sim 2.5\%$ at the
393 turbine location. The result of case TYO shows that the wave height drop is
394 increased to the correct level by activating OBSTACLE; it is increased by \sim
395 0.9% at the turbine location due to the introduction of OBSTACLE. Hence,

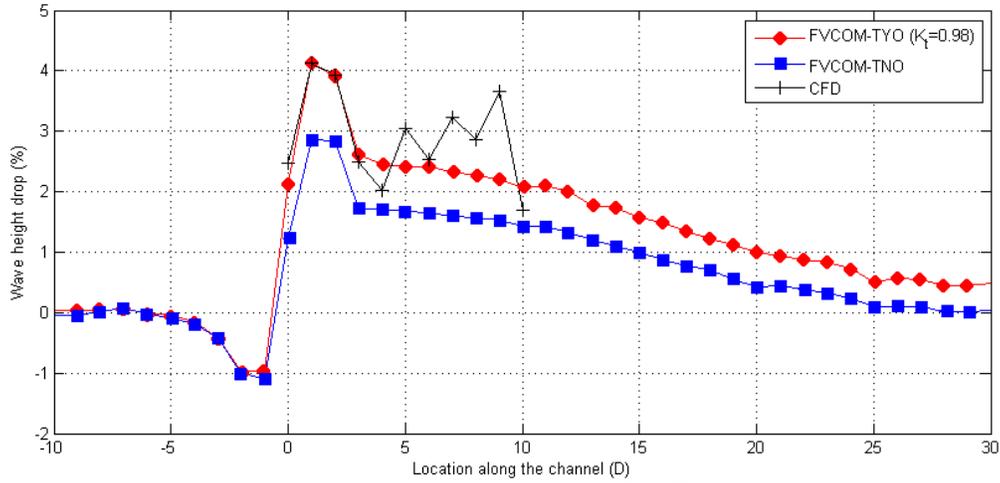


Figure 4: Wave height drop in terms of percentage along the channel for two FVCOM cases, TYO and TNO, and for the wave-current CFD case (the turbine is positioned at 0D).

396 the built-in feature OBSTACLE provides an effective way to simulate the
 397 turbine-caused wave height reduction.

398 The consistency between the CFD and FVCOM simulated wave heights
 399 in the wake of the turbine is obtained through calibrating the wave energy
 400 transmission coefficient K_t mentioned in Section 2.4 according to the results
 401 of the CFD model. However, it should be noted that the two models are
 402 based on different wave theories: the CFD model uses linear wave theory
 403 while the wave model in FVCOM (i.e. SWAN) is a spectral wave model.
 404 The reason the above-mentioned match is achievable despite different wave
 405 theories are applied is that the action balance equation of SWAN (Equation
 406 14) is in fact an energy transfer equation derived based on the linear wave
 407 theory used in the CFD model. The spectrum which contains information
 408 of wave energy in different directions and frequencies can be regarded as a

409 superposition of independent waves following the linear wave theory.

410 **5. Application — Standalone turbine tests**

411 This section investigates the effects of the inclusion of waves and activa-
412 tion of OBSTACLE upon the bottom shear stress based on a series of tests
413 carried out using a prototype 15 m diameter turbine model as the test bed
414 [26]. Water depth of these cases is 45 m and the turbine hub is located at
415 a depth of 22.5 m. The flow and wave conditions are set to reflect those
416 of the Anglesey coast, North Wales, UK, which is identified as one of the
417 potential locations for tidal energy exploitation [53]. The water velocity is
418 1.0 m/s. The significant wave height is 2.4 m and wave period is 7 s: typical
419 conditions of storms observed along the Anglesey coast [54].

420 The results of a current-only case (case TbM (BBL)) and a wave-imposed
421 case without OBSTACLE (case TNO15) are compared to reveal the impact
422 of surface waves on bottom shear stress. Another wave-current coupled case
423 with OBSTACLE activated (case TYO15) is also tested in this section to
424 further discuss how OBSTACLE affects the prediction of bottom shear stress.
425 Turbine simulation in the current and turbulence modules is activated in
426 these cases according to [26]. Bottom shear stress of these three cases are
427 calculated through the BBL module [41] mentioned above. In case TYO15,
428 the OBSTACLE wave energy absorption line (Figure 1) is 15m long and K_t
429 is 0.98.

430 The computed significant wave height of cases TYO15 and TNO15 are
431 shown in Figure 5 (A). Figure 5 (B) & (C) show normalized water velocity
432 at the surface and bottom shear stress for cases TYO15, TNO15 and TbM

433 (BBL). It is observed from Figure 5 (A) that the inclusion of the turbine is
434 causing the significant wave height decrease by $\sim 4.7\%$ beyond 10D down-
435 stream of the turbine and the inclusion of OBSTACLE further reduces the
436 significant wave height by 0.6%.

437 In Figure 5 (B), velocity at the surface increases due to the implemen-
438 tation of the turbine; In this case a peak increase of $\sim 23\%$ is observed for
439 TYO15 1D downstream of the turbine. Further, velocity at the surface for
440 TNO15 is $\sim 4\%$ higher than Tbm (BBL). This is due to the Stokes drift
441 caused by the waves [55]. Note that waves propagating in the same direction
442 of the carrying current are reported to cause a reduction of the flow velocity
443 near the surface [56]. The inclusion of OBSTACLE leads to a reduction in
444 wave height and hence an increase in flow velocity near the surface. This
445 leads to a surface velocity increase of $\sim 3\%$ for TYO15 over TNO15.

446 In Figure 5 (C), it is observed that the inclusion of surface waves increases
447 bottom shear stress by an average of $\sim 7\%$ (for both TYO15 and TNO15)
448 in the regions upstream of the turbine and $>9D$ downstream of the turbine.
449 Difference in bottom shear stress caused by the waves from the turbine within
450 9D downstream of the turbine is relatively small (compared to outside this
451 region). The retarding force which represents the turbine operation is playing
452 the major role within this region, increasing the bottom shear stress by $\sim 50\%$
453 of all three cases. This is a result of the flow acceleration near the bed
454 being identified by a three-dimensional model [26]. Also, the wave bottom
455 boundary layer is likely to be dissipated by the strong mixing caused by the
456 turbine. In the far wake region, as expected, the inclusion of OBSTACLE
457 slightly reduces bottom shear stress compared to TNO15 ($\sim 2\%$ reduction).

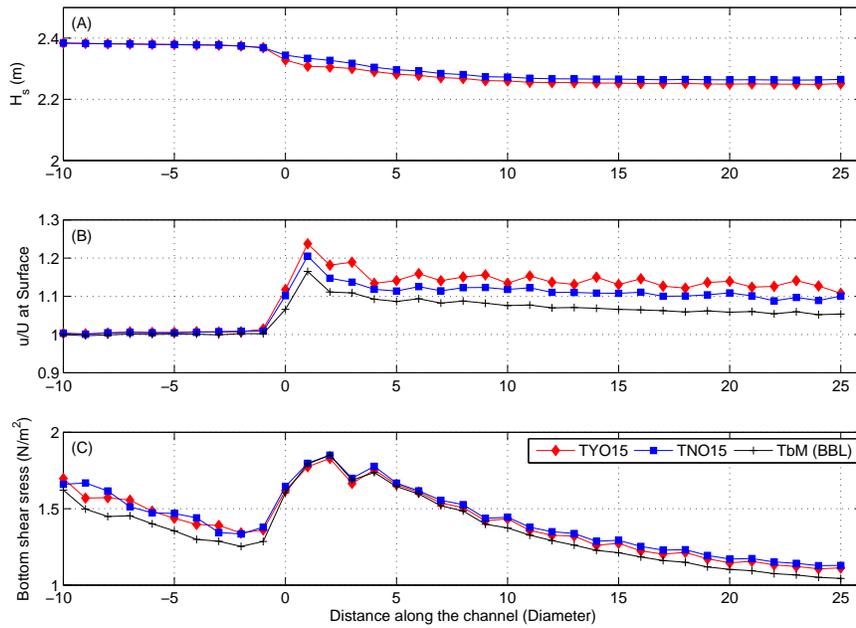


Figure 5: (A) Significant wave height (B) Normalized water velocity at the surface and (C) Bottom shear stress, all calculated under three different scenarios: TYO15 - Retarding force + turbulent terms + waves + obstacle, TNO15 - Retarding force + turbulent terms + waves and TbM (BBL) - Retarding force + turbulent terms with bottom shear stress calculated through BBL. (The turbine is positioned at 0D)

458 6. Discussions

459 6.1. Choice of turbine simulation method in FLUENT

460 Apart from VBM, there are a number of other methods that are widely
461 used to model tidal turbines in CFD simulations, such as the Actuator
462 Disc Method (ADM) which provides a momentum sink in the rotor disk
463 fluid zone without the BEM [57], and the Moving Reference Frame (MRF)
464 method which explicitly simulate the structure and the rotational motion
465 of the turbine [58]. Compared to the fully resolved MRF, VBM has two
466 well-documented limitations: 1) The mechanical turbulence caused by the
467 turbine blades in the form of tip and hub vortex and the blade trailing edge
468 wake is not accounted for [59], leading to under-predicted turbulence level
469 behind the turbine [26]. 2) The lift and drag forces are annularly averaged
470 over a full rotation circle, hence the VBM does not account for transient flow
471 characteristics [10]. This could result in skipping of wave loadings on tur-
472 bines due to the fact that waves can have higher frequencies than the blade
473 passing frequency. Further, large shear can exist across the rotor depend-
474 ing on the vertical flow structure (especially when waves are present as the
475 effect of waves vary significantly with depth), suggesting that the annularly
476 averaged forces could be potentially invalid and a full multi-blade simulation
477 is required to resolve the loadings more realistically. These disadvantages
478 of VBM can result in fallacious power and fatigue analysis, which can ulti-
479 mately lead to inaccurate prediction of design, build and maintenance costs
480 [33]. However, considering that the main focus of this research is the impact
481 of turbines on waves, instead of waves on the performance of turbines, and
482 that the coefficients of VBM can be calibrated against measured data to en-

483 sure acceptable predicted flow conditions in the wake (e.g. [11, 26]), VBM is
484 a viable choice for the purpose of this research. It is also worth noting that
485 the integration of surface waves in CFD simulations can significantly increase
486 the computational effort required, hence VBM which is comparably less com-
487 putationally demanding can serve as a more feasible choice for wave-current
488 simulations, especially in cases where multiple devices are presented.

489 *6.2. Effect of static turbine simulation coefficients*

490 By using VBM to simulate turbines, the lift/drag coefficients ($c_{L,D}$) of
491 the turbine in the CFD simulations are assumed to be static despite the flow
492 conditions. This could be incorrect as surface waves can cause time-varying
493 loadings on turbines which in turn lead to time-dependent effective $c_{L,D}$ [33].
494 In terms of impact assessment, the fixed $c_{L,D}$ used in the CFD simulations
495 could lead to under-/over-estimated instantaneous flow deceleration, turbu-
496 lence generation, wave height modulation and bottom bed shear change.
497 Similarly, the coefficients related to turbine simulation in FVCOM (those in
498 current and turbulent mixing modules [26], as well as K_t in the wave mod-
499 ule mentioned above) are static. Hence, the FVCOM model could also lead
500 to the above-mentioned inaccurate instantaneous predictions. However, it is
501 worth noting that the assessment of turbine-driven local/regional morpholog-
502 ical evolution, which depend highly on the above-mentioned hydrodynamic
503 factors, should take into consideration the life span of tidal turbine arrays
504 which could be up to 100 years [60]. Therefore, the mean overall morpho-
505 logical evolution when considered over such a long time scale could become
506 insensitive to the individual predictions.

507 7. Conclusions

508 The impact of turbines on surface waves is investigated in this study in
509 light of the importance of surface waves on local/regional geomorphology and
510 also as a response to the lack of attention on turbine-induced wave dynamic
511 alternation in the literature. A CFD simulation with a turbine (blockage
512 ratio 3.3% and TSR 5.5) located 0.4 m from the free surface revealed a $\sim 3\%$
513 reduction in wave height as well as a slight increase in wave length. To
514 simulate the wave height drop in FVCOM, the OBSTACLE energy dissi-
515 pation routine of the wave module (SWAN) was activated, and it captured
516 the behaviour to a large extent (Figure 4). However, there are two obvious
517 shortcomings with the modelling method. First, by simply using OBSTA-
518 CLE which subtracts energy from the propagating surface waves, the model
519 does not fully resolve the mechanism of turbine-wave interaction. In this
520 regard, further work is recommended into the investigation of how turbines
521 and surface waves interact. Second, only one turbine configuration is tested
522 at a single depth. However, the specific value of K_t may in fact need to be
523 defined as a function of depth which would also serve as an interesting avenue
524 for investigation.

525 Impacts of tidal turbines on bed shear stress are also studied under wave-
526 current fully coupled scenarios. It is found that although the inclusion of
527 waves increased bed shear stress in the upstream area by an average of $\sim 7\%$,
528 its influence on the bottom shear stress within the near wake zone, i.e. 0D-
529 9D downstream of the turbine, is negligible. The turbine is the dominant
530 factor within this region that increases the bottom shear stress by $\sim 50\%$, as
531 the blockage effect of the turbine forces the water to flow around the device

532 which increases the water velocity near the bed and subsequently increases
533 the bottom shear stress. Impacts of waves on bottom shear stress resume
534 in the far wake, i.e. $>9D$ downstream of the turbine. The influence of
535 OBSTACLE on bottom shear stress is also noticeable in the far wake. The
536 OBSTACLE implemented in this work reduced bottom shear stress by $\sim 2\%$.

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- 546 [1] L. Myers, A. Bahaj, An experimental investigation simulating flow ef-
547 fects in first generation marine current energy converter arrays, *Renew-*
548 *able Energy* 37 (1) (2012) 28–36.
- 549 [2] L. Myers, A. Bahaj, Experimental analysis of the flow field around hor-
550 izontal axis tidal turbines by use of scale mesh disk rotor simulators,
551 *Ocean Engineering* 37 (2) (2010) 218–227.
- 552 [3] F. Maganga, G. Germain, J. King, G. Pinon, E. Rivoalen, Experimental
553 characterisation of flow effects on marine current turbine behaviour and

- 554 on its wake properties, *IET Renewable Power Generation* 4 (6) (2010)
555 498–509.
- 556 [4] S. Tedds, I. Owen, R. Poole, Near-wake characteristics of a model hori-
557 zontal axis tidal stream turbine, *Renewable Energy* 63 (2014) 222–235.
- 558 [5] L. B. Jordan, S. Simmons, S. McLelland, B. Murphy, D. Parsons, L. Vy-
559 bulkova, The impact of tidal stream turbines on 3D flow and bed shear
560 stress measured with particle image velocimetry in a laboratory flume,
561 in: *Proceedings of the 11th European Wave and Tidal Energy Confer-*
562 *ence, Nantes, France, 2015*, pp. 654–660.
- 563 [6] A. Olczak, T. Stallard, P. Stansby, Tidal turbine wake recovery due to
564 turbulent flow and opposing waves, *Proceedings of the 2nd Oxford tidal*
565 *energy workshop*.
- 566 [7] X. Sun, J. Chick, I. Bryden, Laboratory-scale simulation of energy ex-
567 traction from tidal currents, *Renewable Energy* 33 (6) (2008) 1267–1274.
- 568 [8] M. Harrison, W. Batten, L. Myers, A. Bahaj, Comparison between CFD
569 simulations and experiments for predicting the far wake of horizontal
570 axis tidal turbines, *IET Renewable Power Generation* 4 (6) (2010) 613–
571 627.
- 572 [9] L. Bai, R. R. Spence, G. Dudziak, Investigation of the influence of array
573 arrangement and spacing on tidal energy converter (TEC) performance
574 using a 3-dimensional CFD model, in: *Proceedings of the 8th European*
575 *Wave and Tidal Energy Conference, Uppsala, Sweden, 2009*, pp. 654–
576 660.

- 577 [10] X. Bai, E. Avital, A. Munjiza, J. Williams, Numerical simulation of a
578 marine current turbine in free surface flow, *Renewable Energy* 63 (2014)
579 715–723.
- 580 [11] R. Malki, I. Masters, A. J. Williams, T. N. Croft, Planning tidal
581 stream turbine array layouts using a coupled blade element momentum–
582 computational fluid dynamics model, *Renewable Energy* 63 (2014) 46–
583 54.
- 584 [12] A. Goude, O. Ågren, Simulations of a vertical axis turbine in a channel,
585 *Renewable energy* 63 (2014) 477–485.
- 586 [13] A. F. Shchepetkin, J. C. McWilliams, The regional oceanic model-
587 ing system (ROMS): a split-explicit, free-surface, topography-following-
588 coordinate oceanic model, *Ocean Modelling* 9 (4) (2005) 347–404.
- 589 [14] C. Chen, H. Liu, R. C. Beardsley, An unstructured grid, finite-volume,
590 three-dimensional, primitive equations ocean model: application to
591 coastal ocean and estuaries, *Journal of atmospheric and oceanic tech-
592 nology* 20 (1) (2003) 159–186.
- 593 [15] I. G. Bryden, S. J. Couch, ME1 marine energy extraction: tidal resource
594 analysis, *Renewable Energy* 31 (2) (2006) 133–139.
- 595 [16] R. Karsten, J. McMillan, M. Lickley, R. Haynes, Assessment of tidal
596 current energy in the Minas Passage, Bay of Fundy, *Proceedings of
597 the Institution of Mechanical Engineers, Part A: Journal of Power and
598 Energy* 222 (5) (2008) 493–507.

- 599 [17] I. Walkington, R. Burrows, Modelling tidal stream power potential, *Ap-*
600 *plied Ocean Research* 31 (4) (2009) 239–245.
- 601 [18] Z. Defne, K. A. Haas, H. M. Fritz, Numerical modeling of tidal currents
602 and the effects of power extraction on estuarine hydrodynamics along
603 the Georgia coast, USA, *Renewable Energy* 36 (12) (2011) 3461–3471.
- 604 [19] R. Ahmadian, R. Falconer, B. Bockelmann-Evans, Far-field modelling
605 of the hydro-environmental impact of tidal stream turbines, *Renewable*
606 *Energy* 38 (1) (2012) 107–116.
- 607 [20] D. R. Plew, C. L. Stevens, Numerical modelling of the effect of turbines
608 on currents in a tidal channel–Tory Channel, New Zealand, *Renewable*
609 *Energy* 57 (2013) 269–282.
- 610 [21] D. Fallon, M. Hartnett, A. Olbert, S. Nash, The effects of array configu-
611 ration on the hydro-environmental impacts of tidal turbines, *Renewable*
612 *Energy* 64 (2014) 10–25.
- 613 [22] J. Thiébot, P. B. du Bois, S. Guillou, Numerical modeling of the effect of
614 tidal stream turbines on the hydrodynamics and the sediment transport–
615 Application to the Alderney Race (Raz Blanchard), France, *Renewable*
616 *Energy* 75 (2015) 356–365.
- 617 [23] R. Martin-Short, J. Hill, S. Kramer, A. Avdis, P. Allison, M. Piggott,
618 Tidal resource extraction in the Pentland Firth, UK: Potential impacts
619 on flow regime and sediment transport in the Inner Sound of Stroma,
620 *Renewable Energy* 76 (2015) 596–607.

- 621 [24] P. E. Robins, S. P. Neill, M. J. Lewis, Impact of tidal-stream arrays in
622 relation to the natural variability of sedimentary processes, *Renewable*
623 *Energy* 72 (2014) 311–321.
- 624 [25] T. Roc, D. C. Conley, D. Greaves, Methodology for tidal turbine rep-
625 resentation in ocean circulation model, *Renewable Energy* 51 (2013)
626 448–464.
- 627 [26] X. Li, M. Li, S. J. McLelland, L.-B. Jordan, S. M. Simmons, L. O.
628 Amoudry, R. Ramirez-Mendoza, P. D. Thorne, Modelling tidal stream
629 turbines in a three-dimensional wave-current fully coupled oceanographic
630 model, *Renewable Energy* 114 (2017) 297–307.
- 631 [27] G. McCann, Tidal current turbine fatigue loading sensitivity to waves
632 and turbulence—a parametric study, in: *Proceedings of the 7th European*
633 *Wave and Tidal Energy Conference*, Porto, Portugal, 2007.
- 634 [28] G. McCann, M. Thomson, S. Hitchcock, Implications of site-specific
635 conditions on the prediction of loading and power performance of a
636 tidal stream device, in: *2nd International Conference on Ocean Energy*,
637 Brest, France, 2008.
- 638 [29] C. Faudot, O. G. Dahlhaug, Prediction of wave loads on tidal turbine
639 blades, *Energy Procedia* 20 (2012) 116–133.
- 640 [30] E. E. Lust, L. Luznik, K. A. Flack, J. M. Walker, M. C. Van Ben-
641 them, The influence of surface gravity waves on marine current turbine
642 performance, *International Journal of Marine Energy* 3 (2013) 27–40.

- 643 [31] L. Luznik, K. A. Flack, E. E. Lust, K. Taylor, The effect of surface waves
644 on the performance characteristics of a model tidal turbine, *Renewable*
645 *energy* 58 (2013) 108–114.
- 646 [32] T. de Jesus Henriques, S. Tedds, A. Botsari, G. Najafian, T. Hedges,
647 C. Sutcliffe, I. Owen, R. Poole, The effects of wave–current interaction on
648 the performance of a model horizontal axis tidal turbine, *International*
649 *Journal of Marine Energy* 8 (2014) 17–35.
- 650 [33] M. A. Holst, O. G. Dahlhaug, C. Faudot, Cfd analysis of wave-induced
651 loads on tidal turbine blades, *IEEE Journal of Oceanic Engineering*
652 40 (3) (2015) 506–521.
- 653 [34] T. Spencer, I. Möller, F. Rupprecht, T. Bouma, B. Wesenbeeck,
654 M. Kudella, M. Paul, K. Jensen, G. Wolters, M. Miranda-Lange, et al.,
655 Salt marsh surface survives true-to-scale simulated storm surges, *Earth*
656 *Surface Processes and Landforms* 41 (4) (2016) 543–552.
- 657 [35] L. C. Van Rijn, L. C. van Rijn, L. C. van Rijn, Principles of sediment
658 transport in rivers, estuaries and coastal seas, Vol. 1006, Aqua publica-
659 tions Amsterdam, 1993.
- 660 [36] M. S. Longuet-Higgins, Longshore currents generated by obliquely inci-
661 dent sea waves: 1, *Journal of geophysical research* 75 (33) (1970) 6778–
662 6789.
- 663 [37] N. C. Kraus, S. Harikai, Numerical model of the shoreline change at
664 oarai beach, *Coastal Engineering* 7 (1) (1983) 1–28.

- 665 [38] I. Fluent, *Fluent users guide* (2006).
- 666 [39] T. Hedges, Regions of validity of analytical wave theories, *Oceanographic Literature Review* 1 (43) (1996) 10.
- 667
- 668 [40] X. Li, 3D modelling of tidal stream energy extraction for impact assess-
669 ment, Ph.D. thesis, School of Engineering, University of Liverpool (9
670 2016).
- 671 [41] C. Chen, G. Cowles, R. Beardsley, An unstructured grid, finite-volume
672 coastal ocean model: FVCOM user manual, *SMASST/UMASSD*.
- 673 [42] N. Booij, R. Ris, L. H. Holthuijsen, A third-generation wave model for
674 coastal regions: 1. model description and validation, *Journal of Geo-
675 physical Research: Oceans* (1978–2012) 104 (C4) (1999) 7649–7666.
- 676 [43] J. C. Warner, C. R. Sherwood, R. P. Signell, C. K. Harris, H. G. Arango,
677 Development of a three-dimensional, regional, coupled wave, current,
678 and sediment-transport model, *Computers & Geosciences* 34 (10) (2008)
679 1284–1306.
- 680 [44] O. S. Madsen, Spectral wave-current bottom boundary layer flows,
681 *Coastal Engineering Proceedings* 1 (24).
- 682 [45] G. Mellor, The three-dimensional current and surface wave equations,
683 *Journal of Physical Oceanography* 33 (9) (2003) 1978–1989.
- 684 [46] G. Mellor, Some consequences of the three-dimensional current and sur-
685 face wave equations, *Journal of Physical Oceanography* 35 (11) (2005)
686 2291–2298.

- 687 [47] G. L. Mellor, M. A. Donelan, L.-Y. Oey, A surface wave model for cou-
688 pling with numerical ocean circulation models, *Journal of Atmospheric*
689 *and Oceanic Technology* 25 (10) (2008) 1785–1807.
- 690 [48] SWANTeam, SWAN Cycle III version 40.51 user manual, Delft Uni-
691 versity of Technology, Faculty of Civil Engineering and Geosciences,
692 Environmental Fluid Mechanics Section (2006).
- 693 [49] S. Sufian, Numerical modeling of impacts from horizontal axis tidal tur-
694 bines, Ph.D. thesis, School of Engineering, University of Liverpool (6
695 2016).
- 696 [50] D. Maréchal, A soil-based approach to rainfall-runoff modelling in un-
697 gauged catchments for england and wales.
- 698 [51] J. Allen, P. Somerfield, F. Gilbert, Quantifying uncertainty in high-
699 resolution coupled hydrodynamic-ecosystem models, *Journal of Marine*
700 *Systems* 64 (1-4) (2007) 3–14.
- 701 [52] L. Xiaorong, A. Plater, N. Leonardi, Modelling the transport and export
702 of sediments in macrotidal estuaries with eroding salt marsh, *Estuaries*
703 *and Coasts* (2018) 1–14.
- 704 [53] A. Iyer, S. Couch, G. Harrison, A. Wallace, Variability and phasing of
705 tidal current energy around the United Kingdom, *Renewable Energy* 51
706 (2013) 343–357.
- 707 [54] M. Lewis, S. Neill, M. Hashemi, M. Reza, Realistic wave conditions and
708 their influence on quantifying the tidal stream energy resource, *Applied*
709 *Energy* 136 (2014) 495–508.

- 710 [55] K. E. Kenyon, Stokes drift for random gravity waves, *Journal of Geo-*
711 *physical Research* 74 (28) (1969) 6991–6994.
- 712 [56] M. Olabarrieta, R. Medina, S. Castanedo, Effects of wave–current inter-
713 action on the current profile, *Coastal Engineering* 57 (7) (2010) 643–655.
- 714 [57] L. Lavaroni, S. J. Watson, M. J. Cook, M. R. Dubal, A comparison of
715 actuator disc and bem models in cfd simulations for the prediction of
716 offshore wake losses, in: *Journal of Physics: Conference Series*, Vol. 524,
717 IOP Publishing, 2014, p. 012148.
- 718 [58] T. ODoherty, A. Mason-Jones, D. ODoherty, C. Byrne, I. Owen,
719 Y. Wang, Experimental and computational analysis of a model hori-
720 zontal axis tidal turbine, in: *8th European Wave and Tidal Energy*
721 *Conference (EWTEC)*, Uppsala, Sweden, 2009.
- 722 [59] T. Burton, D. Sharpe, N. Jenkins, E. Bossanyi, *Wind energy handbook*,
723 John Wiley & Sons, 2001.
- 724 [60] S. Walker, R. Howell, P. Hodgson, A. Griffin, Tidal energy machines: A
725 comparative life cycle assessment study, *Proceedings of the Institution of*
726 *Mechanical Engineers, Part M: Journal of Engineering for the Maritime*
727 *Environment* 229 (2) (2015) 124–140.

Impact of tidal stream energy device on surface wave dynamics are studied.

A 3D wave-current-sediment fully coupled large-scale numerical model is used.

Impact of turbines on surface waves are incorporated in the large-scale model.

Model prediction indicates a 3% turbine-caused drop in wave height.

Impact of the wave height drop on bed stress in the immediate wake is small.

ACCEPTED MANUSCRIPT