1	Effect of powder characteristics and oxygen content on
2	modifications to the microstructural topology during
3	hot isostatic pressing of an austenitic steel
4	S. Irukuvarghula <sup>*1</sup> , H. Hassanin <sup>2</sup> , C. Cayron <sup>3</sup> , M. Aristizabal <sup>4</sup> , M. M.
5	Attallah <sup>5</sup> , and M. Preuss <sup>1</sup>
6	$^{1}$ School of Materials, University of Manchester, U.K, M1 3BB
7	$^2\mathrm{Department}$ of Civil Engineering and Industrial Design, University of
8	Liverpool, Liverpool L69 3BX, UK
9	<sup>3</sup> Ecole Polytechnique Fedéralé de Lausanne (EPFL), Rue de la Maladière
10	71b, 2000 Neuch $\hat{a}$ tel, Switzerland
11	${}^{4}\text{Ceit-IK4}, P\text{ManuelLardizabal15}, 20018\text{Donostia-SanSebastian},$
12	Spain
13	<sup>5</sup> School of Metallurgy and Materials, University of Birmingham,
14	Edgbaston, B15 2TT, UK

#### Abstract

The effect of powder size distribution and oxygen content on the extent of 16 multiple twinning and spatial distribution of oxide inclusions in hot isostatic pressed 17 (HIPed) 316L steels was investigated using powders with different characteristics. 18 Modifications to, and differences in their microstructural topology, were tracked 19 quantitatively by evaluating the metrics related to twin related domains (TRDs) 20 on specimens produced by interrupting the HIPing process at various points in time. 21 Results revealed that powder size distribution has a strong effect on the extent of 22 multiple twinning in the fully HIPed microstructure, with specimens produced using 23 narrow distribution showing better statistics (i.e., homogeneously recrystallized) 24 than the ones produced using broad size distribution. The oxide inclusion density 25 in fully HIPed microstructures increased with the amount of oxygen content in the 26

15

 $<sup>\ ^*</sup> Corresponding \ author: \ sandeep.irukuvarghula@twi.co.uk$ 

powders while prior particle boundaries (PPBs) were only observed in the specimens
that were HIPed using broad powder distribution. More importantly, results clearly
revealed that the spatial distribution of the inclusions was strongly affected by the
homogeneity of recrystallization. Implications of the results are further discussed
in a broader context, emphasizing the importance of utilizing the occurrence of
solid state phase transformations during HIPing for controlling the microstructure
evolution.

Keywords: austenitic steels, powder metallurgy, hot isostatic pressing, re crystallization, twin related domains

## <sup>36</sup> 1 Introduction

Powder hot isostatic pressing (HIPing) is a manufacturing process that is used to produce 37 near net shape components with fine grain size, chemical homogeneity, and improved in-38 spectability [1]. Used in many industries, including oil & gas and aerospace, it is a 39 thermomechanical process in which powder encapsulated in a canister is consolidated to 40 theoretical density by the application of temperature and pressure. The HIPing condi-41 tions, chosen according to the material (see for e.g., Table. 1 in [1]), ensure complete 42 densification of the powder compact by the end of the HIPing cycle. While the densifi-43 cation kinetics during HIPing for various alloys have been studied to considerable detail 44 [2, 3, 4, 5, 6, 7], the effect of powder characteristics on the microstructure development 45 during HIPing, and on fully HIPed microstructure, still remains unclear, and is of signif-46 icant interest. 47

Albeit the applied pressure during HIPing is isostatic, powders deform inhomoge-48 neously depending on their size; i.e., small particles deform more than large particles 49 [5, 6, 8, 9, 10, 11]. In powder HIPed 316L austenitic stainless steels, it has previously 50 been demonstrated that deformation of the powder is a prerequisite for recrystallization 51 to occur during HIPing [11]. This has been linked to the absence of annealing twins ( $\Sigma 3$ ) 52 boundaries in coincidence site lattice formalism) and associated twin chains in powder 53 particles that do not undergo plastic deformation; these non-deformed powders can be 54 discerned by their prior particle boundaries (PPBs) in the fully consolidated microstruc-55 ture [11]. This observation has a more general applicability to low-medium stacking fault 56 energy (SFE) materials processed by powder HIPing, and is not limited to 316L steels 57 (see for e.g., [12, 13] for microstructures of powder HIPed Inconel 718). 58

It is known that inclusions and second phase particles act as void nucleation sites during ductile fracture, and strongly affect material toughness [14, 15, 16]. Compared to conventionally processed materials, powders contain an order of magnitude more oxygen, which, depending on its solubility in the material, can be in the form of oxide inclusions

(or oxycarbides, depending on the chemistry of the alloy) in the microstructure of a fully 63 HIPed component. It has previously been demonstrated that oxide inclusions lower the 64 impact toughness of powder HIPed 316L and 304L steels [17, 18]. In powder HIPed 65 nickel-base superalloys, presence of inclusions at PPBs in the as-HIPed and heat treated 66 state has been linked to the reduction in their ductility and stress rupture properties 67 [12, 13, 19], although post-HIP hot working has been shown to "break" the PPB networks 68 and randomize the spatial distribution of inclusions [20, 21, 22, 23, 24]. In other words, 69 as a result of recrystallization during post-HIP thermomechanical processing, inclusions 70 at PPBs are pushed by moving grain boundaries, and therefore, are homogeneously dis-71 tributed in the bulk. More specifically, spatial redistribution of the inclusions, which are 72 initially present at PPBs, is promoted by the formation and propagation of twin chains 73 (i.e., recrystallization) during hot working. 74

Since the microstructure in low-medium SFE materials evolves by recrystallization 75 during HIPing (provided there is enough stored energy from particle deformation) [11], 76 understanding the effect of powder characteristics on microstructure development, and 77 more specifically, on the extent and homogeneity of recrystallization – and its effect 78 on the spatial distribution of inclusions, will give insights for the development of high 79 value manufacturing of near net shape HIPed components (e.g., Nickel-base superalloys 80 and austenitic steels for aerospace, oil & gas, and nuclear industries). Therefore, in the 81 present study, by HIPing 316L powders (a representative low-medium SFE material) 82 with different size distributions and oxygen content, we aim to address (a) the effect of 83 powder size distribution on recrystallization and the extent of multiple twinning during 84 HIPing (b) the effect of oxygen content in the powders on the inclusion size and density 85 in the fully HIPed microstructure and (c) the effect of recrystallization on the spatial 86 distribution of oxide inclusions during HIPing. 87

Modifications to the microstructural topology of different powders during HIPing were 88 investigated by quantifying multiple twinning using the metrics related to twin related 89 domains (TRDs). The notion of TRDs as a characteristic microstructural dimension in 90 low-medium SFE materials that are prone to recrystallization was first proposed by Gerts-91 man and Henager [25], and later developed by Reed et al. [26, 27] and Cayron [28, 29]. 92 Specifically, TRDs represent clusters of grains that are related by  $\Sigma 3^n$  misorientations, 93 and connected by chains of  $\Sigma 3$  boundaries [25]. Therefore, the entire recrystallized mi-94 crostructure consists of multiple twinned clusters, i.e., TRDs. Since  $\Sigma 3$  boundaries are 95 generally immune to percolation phenomena (such as crack propagation), and because all 96 grains within TRDs are connected by  $\Sigma 3$  boundaries, it was suggested that the material 97 performance is controlled at the length scale of TRDs [25]. Reed et al. [30] and Barr 98 et al. [31], by providing experimental evidence, clearly demonstrated the existence of a 99 correlation between TRD size and material response towards intergranular phenomena. 100

Therefore, in this study, data analysis was focused towards identifying TRDs and 101 quantifying multiple twinning. Using the metrics related to TRDs, we analyze the data 102 from microstructures produced by interrupting the HIP process at various points in time. 103 Inclusion density (i.e., #/unit area) is measured in the fully HIPed microstructures and 104 is linked back to the oxygen content in the as-received powders. Connections between 105 powder size distribution, PPBs, oxygen content, and the spatial distribution of inclusions 106 on recrystallization are made, and are shown to be associated with the powder deforma-107 tion during HIPing. Suggestions in terms of powder size distribution and oxygen content 108 are put forward, that produce homogeneously recrystallized microstructures with low vol-109 ume fraction and random spatial distribution of inclusions (i.e., not on PPBs) in a near 110 net shape component, potentially eliminating the need for post-HIP thermomechanical 111 processing. 112

## **113 2** Materials and methods

#### 114 2.1 Experimental



Figure 1: Particle size distribution of various powders (For interpretation of the references to colour in this figure legend, the reader is referred to the web version of this article).

Six powders of 316L austenitic steel with different characteristics (chemical and physical) were used in this study <sup>1</sup>. Specifically, four nitrogen atomized powders were obtained from commercial vendors (designated as C, D, E, and F) while two powders were gas atomised at TLS technik, Germany (atomised using argon and nitrogen; designated as A

<sup>&</sup>lt;sup>1</sup>Data for one of the six powders (powder C) were presented in our earlier study [11], and we show them here along with the data of five powders to demonstrate better, the correlation between powder characteristics and their microstructure evolution during HIP consolidation.

and B, respectively), from a forged plate supplied by Rolls-Royce. Particle size distribu-119 tion of the six powders is shown in Fig. 1. It can be seen that powders C and D have 120 much narrower size distribution while others have much broader distribution. Statistical 121 quantities such as  $D_{10}$ ,  $D_{50}$ ,  $D_{90}$ , median, and mean for all particle size distributions are 122 given in Table 1 while their chemical composition and some of their characteristics are 123 shown in Tables 2 and 3, respectively. The oxygen content in the powders varies from 124 110 ppm to 600 ppm and the tap density from 4.9 to 5.7 g.cm<sup>-3</sup>. For interrupted HIP 125 runs, powders were first filled in 25 mm diameter, 30 mm height, and 2 mm thick mild 126 steel canisters, vibrated and vacuum degassed at room temperature before hot crimping 127 the evacuation tube. The canisters were then HIPed using simultaneous application of 128 temperature and pressure, at 950 °C, 1000 °C, 1050 °C, 1120 °C, without any dwell time, 129 at 103 MPa. 130

Metric $(\mu m)$	А	В	$\mathbf{C}$	D	$\mathbf{E}$	$\mathbf{F}$
D <sub>10</sub>	25	34	32	21	20	12
$D_{50}$	82	106	76	61	87	146
$\mathrm{D}_{90}$	205	249	184	147	276	153
Median	82	106	76	61	87	146
Mean	126	151	115	91	147	106

Table 1: Statistical quantities associated with the particle size distributions.

Additionally, powders were filled in mild steel canisters with 76 mm outside diameter, 131 2.3 mm wall thickness, 200 mm height and HIPed at 1160 °C and 103 MPa, and held 132 at those conditions for 4 hours. The canisters were then solution annealed at  $1050 \,^{\circ}\text{C}$ 133 for 1.5 hours and water quenched. These are designated as fully HIPed samples (or 134 fully consolidated state) in this study. For microstructural examination, specimens were 135 sectioned, ground and polished using standard metallographic procedures. Final polish-136 ing was performed on a vibratory polisher using colloidal silica. Backscattered electron 137 (BSE), secondary electron (SE) images, energy dispersive spectroscopy (EDS) data for 138 chemical mapping, and Electron Backscatter Diffraction (EBSD) data were acquired on a 139 Field Emission Gun Scanning Electron Microscope. EBSD maps were acquired to study 140 the differences in the evolution of grain boundary network in all specimens during the 141 HIPing process. A step size of  $0.5 \ \mu m$  and  $1 \ \mu m$  were used for partially HIPed and fully 142 HIPed specimens, respectively. For statistical analysis of the data, EBSD maps from 5 143 random locations were acquired. 144

Powder	$\operatorname{Cr}$	Mn	Mo	Ni	Р	Si	С	$\mathbf{S}$	Ν	Ο	Fe
А	17.93	1.83	2.10	10.09	0.036	0.35	0.028	0.001	0.079	0.029	Bal.
В	17.84	1.85	2.08	10.07	0.035	0.39	0.029	0.001	0.129	0.011	Bal.
$\mathbf{C}$	16.44	1.32	2.08	10.14	0.023	0.57	0.018	0.002	0.098	0.021	Bal.
D	16.94	1.48	2.41	10.45	0.017	0.69	0.016	0.007	0.118	0.061	Bal.
Ε	17.9	1.84	2.44	11.78	0.009	0.73	0.02	0.007	0.061	0.022	Bal.
F	17.6	0.64	2.2	11.3	0.012	0.69	0.012	0.010	0.05	0.0149	Bal.

Table 2: Chemical composition (in wt%) of 316L stainless steel powders determined using inductively coupled plasma mass spectrometry and intert gas fusion.

Table 3: Physical properties/characteristics of the powders.

Property	А	В	С	D	Е	F
Apparent density $(g.cm^{-3})$	5.4	5.3	4.5	4.9	5.2	5.1
Tap density $(g.cm^{-3})$	5.7	5.5	4.9	5.2	5.5	5.6
Flowability (FFC)	17	23	10	19	18	16
Satellites	Yes	Yes	Yes	Yes	Yes	Yes

#### <sup>145</sup> 2.2 Data analysis

Differences in the microstructural topology of the fully HIPed specimens were initially 146 studied by quantifying the distribution of the triple junctions (TJs) associated with  $\Sigma 3$ 147 boundaries. Specifically, three types of TJs were identified:  $J_0$  type triple junction con-148 taining three boundaries which are not  $\Sigma 3$ ,  $J_1$  type containing one  $\Sigma 3$  boundary, and  $J_2$ 149 type containing two  $\Sigma$ 3 boundaries. TJs were identified using scripts written in MAT-150 LAB, and executed as part of the freely available MTEX package [32]. This approach 151 allows qualitative understanding of the extent of twinning. For the EBSD data analysis 152 using MTEX, grains were reconstructed with a threshold misorientation of of 5°. A tol-153 erance angle of 3 ° from the exact misorientation was used for identifying  $\Sigma$ 3 boundaries. 154 Non-indexed pixels, which were typically less than 1.5% of the total pixels in the data 155 (hit-rate in most cases was around 99%), were assigned to surrounding grains. Iden-156 tification of higher order twins, the evolution of TRDs and the associated metrics for 157 intermediate HIP states for all specimens was performed using ARPGE software devel-158 oped by Cayron [33]. More specifically, average values of size of TRDs, number of grains 159 per TRD ( $\langle N_q \rangle$ ), length of longest chain ( $\langle LLC \rangle$ ), polysynthetism ( $\langle p \rangle$ ), and twinning 160 anisotropy  $(\langle a \rangle)$  were analyzed to better quantify multiple twinning. 161

<sup>162</sup> Detailed description and their interpretation are discussed in [29]. Briefly, length of <sup>163</sup> longest chain and polysynthetism represent the longest chain of  $\Sigma 3$  twins in each grain <sup>164</sup> and the tendency to form twin chains of type A-B-A-B-A... where the grains A and B are linked by  $\Sigma$ 3 boundaries in a TRD, respectively. Twinning anisotropy, which depends on the choice of the primary grain used for reconstructing the TRD, represents the structure of the TRD. In other words, it describes the morphology of the TRD from its twinning tree. For statistical analysis of the oxide inclusions, several BSE images of the specimens were acquired and processed using imageJ software [34]. The average size of inclusions, their number density (#/unit area), and the nearest neighbour distance were computed with built-in plugins available in imageJ.

## 172 **3** Results

Figure 2 shows representative images highlighting general microstructural aspects of as-173 received powder. Figure 2a shows morphology of the powder; it is seen that the powder 174 particles are spherical and contain smaller particles (referred to as satellites) on their 175 surfaces. Such a morphology was observed for all powders, with powder C having the 176 most number of satellites. Presence of satellites on powder surfaces reduces flowability 177 and affects powder packing. Figure 2b highlights the surface features of the powder. 178 Small particles,  $< 1 \mu m$ , are seen to adhere the powder surface, presumably from the 179 gas-atomization process. Figure 2c shows the SE image along with the chemical maps 180 of nickel and molybdenum obtained on polished surface of a powder particle. Nickel, 181 and to a lesser extent, molybdenum, are seen to be segregated in what appears to be a 182 cellular microstructure. Figure 2d shows the grain boundary misorientations of a parti-183 cle from the as-received powder. The misorientations are colour coded according to the 184 scheme proposed by Patala et al. [35, 36], which allows complete misorientation represen-185 tation (axis and angle) of grain boundaries. It is seen that the boundaries have a jagged 186 appearance and are predominantly high angle. 187

While the grain boundaries in as-received powder principally are random high angle, 188 the fully HIPed specimens on the other hand contain a large fraction of annealing twins 189 (Fig. 3), indicating that the powders undergo recrystallization during consolidation by 190 HIPing. Qualitatively, Fig. 3 also reveals differences in the microstructures of specimens 191 HIPed with different powders. Specifically, it can be seen that the specimens consolidated 192 using powders A, B, and F show an inhomogeneously recrystallized microstructure with 193 large grains surrounded by several small recrystallized grains (Fig. 3a, Fig. 3b, and 194 Fig. 3f, respectively) while the specimens HIPed using powders C, D, and E are more 195 homogeneously recrystallized (Fig. 3c, Fig. 3d, and Fig. 3e, respectively). 196

Differences in their microstructures were evaluated first by analyzing the distribution of  $J_0$ ,  $J_1$ , and  $J_2$  type triple junctions. Figure 4 shows a representative grain boundary misorientation map of the fully HIPed specimen (using powder D) overlayed with different types of triple junctions. It is seen that most of the triple junctions are of  $J_1$  type, followed



(a)

(b)



(c)



Figure 2: Representative images showing (a) the morphology of as-received powder D (b) small particles adhering to the powder surface (c) SE image and corresponding chemical maps of nickel and molybdenum (d) the grain boundary misorientations in a powder particle using electron backscatter diffraction, along with its legend (For interpretation of the references to colour in this figure legend, the reader is referred to the web version of this article).







Figure 4: Representative map showing grain boundary misorientations colour coded according to the legend in Fig. 2b, overlayed with various triple junction types. Here, triple junctions consisting of three non- $\Sigma$ 3 boundaries (i.e.,  $J_0$ ) are marked by  $\blacktriangle$ ; one  $\Sigma$ 3 boundary (i.e.,  $J_1$ ) by  $\blacktriangle$ ; two  $\Sigma$ 3 boundaries (i.e.,  $J_2$ ) by  $\blacktriangle$  (For interpretation of the references to colour in this figure legend, the reader is referred to the web version of this article).

by  $J_0$  and  $J_2$ , respectively. The distribution of types of triple junctions in fully HIPed 201 specimens is tabulated in Table 4. In all specimens,  $J_1$  type dominates, followed by  $J_0$  and 202  $J_2$ . The statistics also confirm the qualitative observations made from the BSE images 203 (Fig. 3). Specifically, compared to A, B, and F, specimens C, D, and E have lower  $J_0$ 204 type junctions and higher fraction of  $J_1$  triple junctions. In other words, specimens A, 205 B, and F have lower number fraction of  $\Sigma 3$  boundaries compared to C, D and E; it then 206 follows that they also contain more grains which have not undergone multiple twinning. 207 This further suggests that the observed differences arise from their contrasting powder 208 characteristics (see Fig. 1, and Tables 1, 2, and 3). 209

It is emphasized that, owing to the crystallographic constraint at triple junctions [37], 210 the third boundary in a  $J_2$  type triple junction can either be  $\Sigma 9$  or  $\Sigma 1$  boundary (i.e.,  $\Sigma 3^2$ 211 and  $\Sigma 3^0$ , respectively), while for the  $J_1$  type junction, the other two boundaries will be 212 higher order twins (i.e., according to the following rule:  $\Sigma 3^n \cdot \Sigma 3^m = \Sigma 3^{n+m-2i}$ , where *i* is 213 an integer between 0 and n) [26, 28]. It is also possible that the  $J_0$  type junctions contain 214 higher order twins (e.g.,  $\Sigma 9 - \Sigma 9 - \Sigma 9$ ) and therefore in that respect, the above analysis 215 only gives a rudimentary picture regarding the differences in grain boundary network 216 topology between the specimens. Therefore, in order to identify higher order twins and 217 to better understand the topological differences, we go beyond the quantification of triple 218 junctions associated with  $\Sigma 3$  boundaries, and evaluate advanced metrics associated with 219 twin related domains. 220

A representative image with TRDs for specimen E reconstructed using ARPGE is shown in Figure 5a. Here, the grain boundaries can be identified using the legend shown in Figure 5b, with the numbers representing n in  $\Sigma 3^n$ . Figure 5c is the largest TRD highlighted in Figure 5a (with an arrow). All grains in this TRD are related by  $\Sigma 3^n$ misorientations, and are connected by  $\Sigma 3$  boundaries; they are shown on a twinning tree for this TRD using a fractal representation (Figure 5d). All parameters associated with the TRDs are automatically calculated using ARPGE and then averaged.

Specimen	$J_0$	$J_1$	$J_2$
А	$36 (\pm 3)$	$56 (\pm 2)$	$8(\pm 1)$
В	$36 (\pm 3)$	$55 (\pm 4)$	$9(\pm 1)$
С	$31 \ (\pm 2)$	$62 \ (\pm 2)$	$7(\pm 1)$
D	$29 (\pm 3)$	$62 \ (\pm 2)$	$9(\pm 1)$
Ε	$28 (\pm 1)$	$61 \ (\pm 1)$	$11 (\pm 1)$
F	$36 (\pm 3)$	$56 (\pm 3)$	$8(\pm 1)$

Table 4: Triple junction distribution in the fully consolidated specimens.

Figures 6a-f show the metrics (i.e., Length fraction of  $\Sigma 3$ ,  $\langle TRD \rangle$ ,  $\langle N_g \rangle$ ,  $\langle LLC \rangle$ ,  $\langle p \rangle$ , and  $\langle a \rangle$ ) quantifying the extent of multiple twinning for the evolutionary microstructural



Figure 5: (a) Representative map showing the reconstructed TRDs for specimen E (b) the legend for identifying various boundaries, with the numbers representing n in  $\Sigma 3^n$  (c) largest TRD identified in the map along with grain numbers (d) fractal representation of twin chains in the TRD identified by the arrow in (c); here the grains numbers occupy red circles while the squares show the operators that connect the grains.  $\Sigma 3$  corresponds to 1a,  $\Sigma 9$  corresponds to 2a and so on (For interpretation of the references to colour in this figure legend, the reader is referred to the web version of this article).



Figure 6: Evolution of various parameters associated with TRDs for the specimens as a function of HIPing temperature: (a) length fraction of  $\Sigma 3$ grain boundaries (b) average sizes of TRDs, (c) number of grains in a TRD, (d) length of longest chain, (e) twinning polysynthetism, and (f) twinning anisotropy. Lines joining the data points are only a guide to the eye. Here FC stands for 'Fully Consolidated'.

states of all specimens, and very clearly reveal the differences between them. It can be seen 230 that the statistics for all powders HIPed at 950 °C are very similar, but begin to diverge as 231 the HIPing temperature increases (Figures 6a-f). This is because, there is incipient plastic 232 deformation of powders at 950 °C, and the corresponding microstructures are similar to 233 those of the as-received powders. Figure 6a shows the evolution of length fraction of 234  $\Sigma$ 3 boundaries as a function of HIPing temperature. While the length fraction in all 235 specimens increases with HIPing temperature, the rate of increase (with temperature 236 and hence, time) is different. In the fully HIPed condition,  $\Sigma 3$  boundary length fraction 237 of C, D, and E is similar ( $\sim 51\%$ ) and higher than those for A, B, and F ( $\sim 45\%$ ). Specimen 238 C has the largest  $\langle TRD \rangle$  size (46 µm), followed by D, E (~31 µm) then by F (23 µm), 239 and A and B (19  $\mu$ m), (Fig: 6b). 240

The average number of grains in the TRDs (i.e.,  $\langle N_q \rangle$ ) are lowest for A, B, and F (~3) 241 per TRD), while those for C, D, and E are much higher (3.5, 4, and 4.5, respectively), 242 Figure 6c. The average length of longest chain (LLC) that represents the twinning 243 order in a TRD, is greatest for E (1.62), followed by D (1.55) and C (1.43), and then 244 by F, A, and B (0.98, 0.9, 0.88, respectively), Figure 6d. The twins in E, D, and C 245 are more polysynthetic ( $\langle p \rangle$  values are 1.29, 1.27, 1.26, respectively) than F, A, and B 246 (1.16, 1.15, 1.13, respectively), Figure 6e. The average values for twinning anisotropy  $\langle \langle a \rangle \rangle$ 247 are in the following order: E, C, D (1.83, 1.66, 1.55, respectively), and F, A, and B (1.44, (1.44, 1.55)) 248 1.39, 1.39, respectively), Figure 6f. Figures 6a-f also reveal that the propagation of twin 249 chains is the most during the dwell time of HIPing cycle. 250



Figure 7: (a) A representative BSE image showing a large number of "holes". Chemical maps for manganese and oxygen (b) and (c), respectively, showing that the inclusions are enriched in manganese and oxygen.

As mentioned previously, powder processed materials usually contain an order of magnitude more oxygen than cast and forged components. Our results show the specimens containing numerous inclusions around 0.5  $\mu$ m in diameter, enriched in manganese and oxygen, although their precise chemical composition was not evaluated (Figure 7). Most

of the inclusions were dislodged during grinding and polishing, leaving "holes" in the 255 matrix, suggesting that they are very weakly bonded to the matrix. This interpretation 256 is consistent with the results from a previous study from our group [18]. Specifically, elec-257 tropolishing a HIPed 316L specimen resulted in selective removal of the matrix material 258 while the inclusions were retained in the matrix. Moreover, a comparison of microstruc-259 tures between two specimens with different oxygen contents prepared using mechanical 260 polishing and electropolishing, respectively, revealed that there were more holes (inclu-261 sions) in mechanically polished (electropolished) specimens with higher oxygen content, 262 conclusively showing that the holes are in fact inclusions dislodged during mechanical 263 polishing [18]. 264

In order to estimate their fraction in the specimens, the "holes" were treated as ox-265 ide inclusions. Their statistical analysis is shown in Table 5. Their density and average 266 size increase with the oxygen content in the powders, while the nearest neighbour (NN) 267 distance decreases. Although the oxygen content in specimens B and F are very low com-268 pared to others (110 ppm and 150 ppm, respectively), the errors associated with inclusion 269 density and NN distance in them are high. This was observed to be the consequence of 270 their presence along prior particle boundaries (PPBs). Specifically, it was observed that 271 if the PPBs were present in the specimens, inclusions decorated them, and were more 272 closely spaced along them. Moreover, PPBs were only observed in specimens with broad 273 powder size distributions (A, B, and F) and virtually nonexistent in specimens C, D, 274 and E. Interestingly, specimen D which has the highest oxygen content (600 ppm) and 275 a narrow powder size distribution, was homogeneously recrystallized and did not con-276 tain any PPBs. In other words, the inclusions were randomly distributed in the bulk. 277 Representative images for specimen F (110 ppm) and D (600 ppm) are shown in Figure 278 8. PPBs are seen in Figure 8a which contains fewer annealing twins than in Figure 8b. 279 The results thus indicate a correlation between (a) powder size distribution and the ex-280 tent multiple twinning (i.e., recrystallization) (b) oxygen content in the powder and the 281 inclusion density, and (c) spatial distribution of oxide inclusions and recrystallization. 282

Table 5: Inclusion distribution in the fully HIPed microstructure.

	А	В	С	D	Е	F
Oxygen content	$290~\rm ppm$	110  ppm	$200~\rm ppm$	600  ppm	$220~\rm ppm$	150  ppm
Inclusion density $(x10^9/m^2)$	$16 (\pm 3)$	$7.5 (\pm 5.5)$	$8.9 (\pm 1)$	$21.5 (\pm 3.6)$	$12.1 \ (\pm 2.8)$	$20 \ (\pm 10.7)$
Nearest neighbour distance $(\mu m)$	$3.64 \ (\pm 0.39)$	$4.82 (\pm 3.9)$	$4.23 \ (\pm 0.52)$	$3.23 (\pm 0.38)$	$4.18~(\pm 0.58)$	$2.11 \ (\pm 0.9)$
Average size $(\mu m)$	$0.46~(\pm 0.06)$	$0.43 \ (\pm 0.09)$	$0.49~(\pm 0.05)$	$0.5~(\pm 0.04)$	$0.41~(\pm 0.07)$	$0.36 (\pm 0.1)$



Figure 8: Secondary electron showing the effect of recrystalization on inclusion distribution in fully HIPed specimen of powder F (a) and D (b).

### 283 4 Discussion

# 4.1 Effect of powder size distribution on recrystallization dur ing HIPing

Depending on the powder size distribution, the extent to which the powder compact 286 is plastically strained during early stages of HIPing varies. Specifically, powders with 287 low tap density experience more strain as they can accommodate more plastic flow than 288 powders with high tap density. In fact, Li and Funkenbusch, in their HIP model, have 289 demonstrated that the deformation of powder in a monosize distribution could be almost 290 twice of that in a bimodal distribution [6]. Since the stored energy increases with the 291 amount of plastic strain experienced, so will the driving force for recrystallization during 292 HIPing. In other words, powder size distribution has a strong effect on the extent of 293 multiple twinning during HIPing. It is pointed out that during high temperature defor-294 mation of the powder particles during HIPing, dislocation recovery mechanisms will also 295 be operative, and will affect the microstructural evolution. However, in order to precisely 296 understand them, various high temperature microstructural states need to be preserved 297 by rapid quenching. This is extremely challenging as the cooling rates during the HIPing 298 process, generally, are slow to capture them. 299

From Fig. 1 and Table 3, it can seen that the powders considered in this study have different size distributions and consequently, varied tap densities (i.e., in the following order: C < D < E < B < F < A). Therefore, the metrics quantifying multiple twinning can be expected to follow a similar trend, although it must be noted that locally, powders experience strains that will be different from the macroscopic strain due to the shrinkage of canisters during HIPing. In other words, while the differences in tap density/powder size
distribution give an indication of the relative strains, the development of multiple twinned
microstructure depends on the energy stored locally. More specifically, it is analogous
to the dependence of applied strain on the differences in the extent of multiple twinning
observed in the grain boundary engineering studies of low-medium SFE materials (e.g.,
[38, 39, 40]).

From Table 4, which shows the distribution of various triple junction types, it can be 311 seen that  $J_1$  type junctions are higher in specimens C, D, and E compared to A, B, and F. 312 In other words, the number of annealing twins is higher in C, D, and E than in A, B, F. 313 The average number of grains,  $\langle N_g \rangle$ , within TRDs is highest in E, followed by D and C, 314 while they are lowest for A, B, and F. It suggests that the twin chains in specimens C, D, E 315 have propagated more than those in A, B, and F. This is further confirmed if we compare 316 (LLC),  $\langle p \rangle$ , and  $\langle a \rangle$ . It is seen that these parameters are nicely grouped for C, D, E (a 317 shade higher), and A, B, F (comparitively lower), clearly revealing that broad powder 318 size distributions, owing to higher tap densities (and consequently, experiencing lower 319 plastic strains during HIPing), have lower driving force for recrystallization compared 320 to powders with lower tap densities (Table. 3). In summary, a narrow powder size 321 distribution with low tap density produces a homogeneously recrystallized microstructure. 322 Such a microstructure is beneficial from the point of view of enhanced resistance towards 323 material degradation mechanisms that propagate intergranularly (e.g., stress corrosion 324 cracking). 325

#### 326 4.2 Mechanism of formation of oxide inclusions

As previously pointed out, during the early stages of HIPing, plastic deformation of 327 powders is initiated at their contact points with other powders. The temperature and 328 pressure at which the deformation starts and proceeds depends on the elevated temper-329 ature strength of the material and the powder size distribution. Hedberg et al. have 330 shown that gas atomised 316L powder surface consists of a homogeneous layer (< 10 nm) 331 of iron/manganese/chromium oxide [41]. Because of the presence of a continuous oxide 332 film, complete bonding between powders during their incipient deformation at elevated 333 temperature (i.e., at  $\sim 950$  °C in this study) cannot occur. The oxide layer is disrupted 334 during HIPing and metal to metal contacts are established. In order to reduce the inter-335 facial energy between the metal and oxide layer, the disrupted oxide layer coalesces to 336 form inclusions ( $\sim 0.2 \ \mu m$ ) at the inter-particle boundaries, Fig. 9a. 337

With increasing temperature, powders experience larger strains and simultaneously begin to recrystallize, Fig. 9b, where the presence of annealing twins can be seen. However, the inclusions are still present at the inter-particle boundaries. At 1050 °C, depending on the local plastic strain experienced by the powders, the recrystallized grains begin





Figure 9: Representative BSE images highlighting the effect of recrystallization on the randomisation of inclusions in the microstructure (HIPed specimens from powder D). Large dark regions between the powders in (a) and (b) correspond to the porosity resulting from partial consolidation.

to grow past the inter-particle boundaries with the inclusions still arranged on them, 342 Fig. 9c. However, after the full HIP cycle (i.e., HIPed at 1160 °C, 103 MPa, 4 hours 343 dwell time), the inclusions are distributed in the bulk, depending on the homogeneity 344 of recrystallization, albeit there are few regions where they are situated at the grain 345 boundaries (Fig. 9d), indicating that inclusions are dragged by moving grain boundaries 346 during recrystallization. Moreover, their average size has also increased to around  $0.5 \ \mu m$ 347 (refer Table 5), further suggesting that the coarsening of inclusions is diffusion and/or 348 coalescence induced. We note that the growth of oxide inclusions during recrystallization 349 has previously been reported for Ni-Cr and Cu-Si alloys, respectively [42, 43]. Once nu-350 cleated, the oxide inclusions must be quite stable at HIPing temperatures (i.e.,  $1160 \,^{\circ}$ C). 351 This is because, if they were to be in solid solution, there is no a priori reason for their 352 re-nucleation exactly at PPBs during the cooling step of HIPing process. We further note 353 the presence of sub-micron sized cellular microstructure in the partially HIPed specimens 354 (evident in Fig. 9a-c), which is the consequence of rapid solidification of the melt during 355 gas atomisation process (i.e., powder production process). However, the microstructure 356 homogenizes during the "dwell time" of the HIPing cycle. 357

#### <sup>358</sup> 4.2.1 Effect of oxygen content on recrystallization

The powders considered in the present study had oxygen content ranging from 110 ppm 359 (powder B) to 610 ppm (powder D), Table 2. We now compare the metrics for specimens 360 B and D (in other words, the lowest and highest oxygen content specimens, respectively) 361 to specifically understand the effect of oxygen content in the powder on recrystallization 362 during HIPing. Firstly, the microstructures qualitatively reveal that specimen B has re-363 crystallized inhomogeneously while specimen D has a more homogeneous microstructure. 364 The triple junction distributions (Table. 4) quantitatively indicate a higher fraction of 365  $\Sigma$ 3 boundaries in D compared to B (fraction of  $J_0$  is lower while  $J_1$  is higher in D). In 366 addition, the advanced metrics related to TRDs (i.e.,  $\langle TRD \rangle$ ,  $\langle N_a \rangle$ ,  $\langle LLC \rangle$ ,  $\langle p \rangle$ , and  $\langle a \rangle$ ) 367 show that the extent of multiple twinning is much larger in specimen D than in B, clearly 368 indicating that the oxygen content has no effect on recrystallization. On the other hand, 369 the powder size distribution of B is much broader than D (tap densities of  $5.5 \text{ g.cm}^{-3}$  and 370  $5.2 \text{ g.cm}^{-3}$ , respectively) due to which the plastic strain (and hence stored energy) in the 371 compact of powder B during HIPing will be lower than powder D. This in turn results in 372 larger driving force for recrystallization in D and hence, better twinning stastistics than 373 in B. 374

Rao et al. studied the effect of oxygen content on the microstructure of powder HIPed Inconel 718 and concluded that the effect of oxygen content in the powders is to retard recrystallization [12]. Specifically, using three powders with different oxygen contents (275 ppm, 180 ppm, and 140 ppm) they observed (qualitatively, from the micrographs)

that the extent of recrystallization decreased with an increase in oxygen content, although 379 the effect of powder size distribution was ignored in their analysis (refer Figures 2c and 380 3c in [12]). More specifically, the tap density was highest for the powder with 275 ppm 381 oxygen (broad size distribution) and lowest for 140 ppm (narrow size distribution), which 382 clearly suggests that the driving force for recrystallization for the powder with lowest tap 383 density will be greater than that for highest tap density. We believe that the results of 384 Rao et al. tacitly demonstrated the dependence of powder size distribution on recrystal-385 lization, but were interpreted otherwise. In summary, our results suggest that the oxygen 386 content in the powders has no effect on the degree of plastic strain experienced during 387 initial stages of HIPing, although in order to prove it conclusively, powders with same 388 particle size distribution but different oxygen content need to be HIPed. For a particular 389 size distribution of powders, the oxygen content, for example, can be increased by heat 390 treating them under controlled conditions. 391

## 4.2.2 Spatial distribution of oxide inclusions during consolidation and the effect of recrystallization

While powder HIPed components can be manufactured in near net shape, one of the 394 challenges in using them in high integrity structural applications is their lower tough-395 ness values (attributed to the presence of oxide inclusions) compared to cast and forged 396 ones, although it is emphasized that the toughness values in many cases are still above 397 the ones stipulated in ASME codes. See for e.g., [17, 18] for the case of powder HIPed 398 316L and 304L. The size, volume fraction, coherency with the matrix, and spatial dis-399 tribution of inclusions have a strong effect on the material toughness. It is emphasized 400 that powder particles, regardless of their size, contain an oxide layer on their surface; 401 the amount of oxygen depends on their size, gas atomization process, and storage con-402 ditions. Presence of PPBs, and therefore, inclusions decorating them, indicates that the 403 corresponding particles have not plastically deformed during HIPing process. Figure 8 404 shows the effect of recrystallization on the spatial distribution of inclusions. Specifically, 405 Fig. 8a, which corresponds to powder F (highest tap density and lowest oxygen content), 406 reveals a particular region where the inclusions are decorated at PPBs while Fig. 8b, 407 which corresponds to powder D (lowest tap density and highest oxygen content), shows 408 a random distribution of the inclusions. The triple junction distribution and twinning 409 statistics reveal that the scale and homogeneity of recrystallization in specimen D is much 410 better than specimen F, clearly demonstrating that recrystallization strongly affects the 411 spatial distribution of the inclusions. 412

<sup>413</sup> During powder HIPing, a general view regarding the effect of inclusions at PPBs is <sup>414</sup> that they act as pinning centres for moving grain boundaries [12, 19, 44, 45, 46]. In-<sup>415</sup> deed, one of the ways suggested to promote the movement of grain boundaries past

them is to perform HIPing at temperatures where the driving force for the boundary 416 movement is greater than the pinning force of the inclusions at PPBs, although it is rec-417 ognized that it results in increased grain size [13, 19]. However, pushing of particles by 418 moving grain boundaries during solid state phase transformations (e.g., recrystallization, 419 allotropic transformation), and diffusional movement of inert particles have previously 420 been reported [42, 43, 47, 48, 49, 50, 51, 52]. These studies focused on the effect of 421 recrystallization in a material with homogeneous dispersion of inclusions. However, in 422 powder HIPing, the opposite holds; i.e., the inclusions nucleate at inter-particle bound-423 aries (or PPBs) where the contact stresses are the highest (provided there is room for 424 plastic flow), and are then pushed away from PPBs as a result of recrystallization during 425 the dwell time of the HIPing cycle. 426

While it is known that, in low-medium SFE materials, recrystallization occurs during 427 HIPing, previous studies have not explicitly correlated the spatial distribution of inclu-428 sions to multiple twinning. In fact, results from this study have clearly demonstrated that, 429 although the inclusions nucleate at the inter-particle boundaries, their spatial distribu-430 tion after the HIPing cycle (i.e., after complete consolidation) is strongly affected by the 431 extent of recrystallization, which is governed by the powder size distribution. Therefore, 432 in a much broader context, the naturally occurring mechanisms of phase transformation 433 and related phenomena during powder processing of certain alloy systems (e.g., recrys-434 tallization in low-medium SFE materials, transformation mismatch plasticity in titanium 435 alloys and martensitic/bainitic steels; see for e.g., [53, 54]) can be used for optimizing the 436 manufacturing process and controlling the microstructural evolution. 437

## 433 4.3 Implications of the present study for powder-HIPed com-439 ponents with low-medium SFE

One of the main results originating from the present study is that, the homogeneity of 440 recrystallization during HIPing is governed by the powder size distribution and not by the 441 amount of oxygen content. More importantly, the spatial distribution of oxide inclusions 442 and their volume fraction are governed by the extent of recrystallization and the amount 443 of oxygen content in the powders, respectively. We believe that these observations are 444 very important and are of high value, and can be used to potentially manufacture near 445 net shape components for critical applications only by powder-HIPing. As an example, 446 we identify its applicability to powder processed nickel-base superalloys. 447

Highly alloyed nickel-base superalloys which are of interest for demanding environments, are processed using powder metallurgy techniques to avoid casting induced segregation of alloying elements. However, presence of PPBs in the as-HIPed condition has been a perennial problem, and has necessitated post-HIP thermomechanical processing

of these alloys, which not only increases the manufacturing cost, but also restricts the 452 efficient use of HIPing as a process for manufacturing near net shape components. Specif-453 ically, processing involves canning of gas-atomized powders and HIPing, followed by hot 454 extrusion and/or isothermal forging [20, 21]. Post-HIP thermomechanical processing is 455 performed in order to "break" the PPB network present in the as-HIPed products (i.e., 456 via recrystallization). Hot-working of these alloys after HIPing has been shown to be 457 effective not only in enhancing the fraction of  $\Sigma 3$  boundaries, but also in breaking the 458 PPB networks [22, 55, 56], clearly suggesting a direct correlation between the extent and 459 homogeneity of recrystallization and the absence of PPBs. 460

In our study on powder HIPed 316L, by considering different powder size distribu-461 tions, and by just HIPing, we clearly demonstrated a correlation between the extent and 462 homogeneity of recrystallization, and the absence of PPBs (Fig. 3, Table. 4, Fig. 6, and 463 Fig. 8). More specifically, a narrow powder size distribution has produced a more ho-464 mogeneously recrystallized microstructure free from PPBs than the broadly distributed 465 ones. While a broad powder size distribution (in other words, distribution with high tap 466 density) provides better control over the shape change of the component after HIPing<sup>2</sup>, 467 modelling can be used to predict the shape change so that the initial canister can be 468 designed in order to obtain near net shape after HIPing even with narrow powder size 469 distribution. The benefit of uniform recrystallization during HIPing is two-fold. First, 470 the increased fraction of  $\Sigma 3$  grain boundaries due to multiple twinning during HIPing can 471 improve high cycle fatigue crack propagation behaviour [57]. Second, the inclusions that 472 nucleate at inter-particle boundaries will more likely be dispersed homogeneously after 473 HIPing (because of the absence of PPBs), potentially improving the ductility, although 474 it must be noted that any improvement in ductility is dependent on the volume fraction 475 of inclusions. 476

While it can be argued that the propensity for twinning depends on the SFE of 477 the material, and that there would be differences in microstructural topology between 478 316L and Ni-base superalloys, we emphasize that hot working has a strong effect on 479 grain boundary network topology, and we also underscore that HIPing is essentially a 480 thermomechanical process. For example, within the context of twinning induced grain 481 boundary engineering in castings, the effect of process variables (e.g., strain, annealing 482 temperature and time, or strain rate and hot deformation temperature) on grain boundary 483 connectivity can be inferred by referring to [58] for a general overview and [59, 60, 61] 484 specifically for Ni-base superalloys. In the case of powder HIPing, the process variables 485 are powder size distribution, temperature, pressure, and time. 486

487

The effectiveness of HIPing as a thermomechanical process can therefore be taken

<sup>&</sup>lt;sup>2</sup>During HIPing, shrinkage of the canister is more uniform for a broad powder size distribution compared to a narrow size because of the lesser extent of plastic flow offered.

advantage of, to reduce/potentially eliminate the problem of PPBs by proper choice of 488 powder size distribution, chemistry, and HIPing parameters. More specifically, the in-489 clusion density can be reduced by using powder with low oxygen content while their 490 spatial distribution can be randomized by optimizing the HIP process variables to pro-491 mote homogeneous recrystallization. That is, in addition to using narrow powder size 492 distribution, the stored energy in the deformed powders can be tuned by adjusting the 493 way in which pressure and temperature are applied in order to promote multiple twinning 494 and eliminate PPBs during HIPing. While the results in the present study were obtained 495 on specimens with relatively simple geometry, getting a homogeneous microstructure in 496 a large complex HIPed part, due to potential variations in strain distribution, could be 497 challenging, and needs to be demonstrated. In this regard, additional studies focusing 498 on effect of powder characteristics and HIP process variables on the spatial variations in 499 microstructural features in large powder-HIPed components are needed. 500

## 501 Conclusions

The aim of the present study was to understand the effect of powder characteristics and oxygen content on the microstructural evolution during hot isostatic pressing of 316L austenitic steel. The main findings are summarized as follows:

- The principal mechanism by which the microstructure evolves during HIPing is recrystallization, the homogeneity of which, is strongly dependent on the powder size distribution.
- Fully HIPed specimens with narrow powder size distribution (i.e., lower tap density)
   were more homogeneously recrystallized than the ones produced using broad size
   distribution (i.e., higher tap density). Prior particle boundaries (PPBs) were more
   frequently observed in specimens HIPed with broad powder size distribution than
   the ones HIPed with narrow distribution.
- The density of oxide inclusions in the fully HIPed specimens increased with the oxygen content in the powders, with their spatial distribution strongly dependent on the homogeneity of recrystallization.
- The oxide inclusions rich in manganese nucleated on inter-particle boundaries (in other words, PPBs) during early stages of HIPing, and depending on whether or not there was sufficient stored energy, were pushed by the moving grain boundaries during recrystallization.
- Finally, as demonstrated here, homogeneous recrystallization during HIPing results in random distribution of oxide inclusions and a microstructure free from PPBs. This is

highly relevant from the point of view of manufacturing PPB free near net shape components for critical applications, since it would eliminate the need for post-HIP thermomechanical processing (in order to break the PPBs) and reduce manufacturing costs.

## 525 Acknowledgements

The authors would like to thank the EPSRC for funding through New Nuclear Manufacturing (NNUMAN): EP/J021172/1, and Manufacture using Advanced Powder Processes (MAPP): EP/P006566/1. Rolls-Royce is also thanked for supplying the 316L forging for this study.

## 530 References

- [1] H. V. Atkinson and S. Davies. Fundamental aspects of hot isostatic pressing: an overview. *Metallurgical and Materials Transactions A*, 31A:2981–3000, 2000.
- [2] D. S. Wilkinson and M. F. Ashby. Pressure sintering by power law creep. Acta
   Metallurgica, 23(11):1277 1285, 1975.
- [3] E. Arzt. The influence of an increasing particle coordination on the densification of spherical powders. *Acta Metallurgica*, 30(10):1883 – 1890, 1982.
- [4] A. S. Helle, K. E. Easterling, and M. F. Ashby. Hot-isostatic pressing diagrams:
   New developments. Acta Metallurgica, 33(12):2163-2174, 1985.
- [5] S. V. Nair and J. K. Tien. Densification mechanism maps for hot isostatic pressing
  (HIP) of unequal sized particles. *Metallurgical Transactions A*, 18(1):97–107, 1987.
- [6] E. K. H. Li and P. D. Funkenbusch. Hot isostatic pressing (hip) of powder mixtures
   and composites: Packing, densification, and microstructural effects. *Metallurgical Transactions A*, 24(6):1345–1354, 1993.
- [7] D. P. Delo and H. R. Piehler. Early stage consolidation mechanisms during hot
   isostatic pressing of Ti-6Al-4V powder compacts. Acta Materialia, 47(9):2841–2852,
   1999.
- [8] W. A. Kaysser, M. Aslan, E. Arzt, M. Mitkov, and G. Petzow. Microstructural
   Development and Densification During Hipping of Ceramics and Metals. *Powder Metallurgy*, 31(1):63-69, 1988.
- [9] R. N. Wright, R. L. Williamson, and J. R. Knibloe. Modelling of Hipping Consoli dation Applied to Ni<sub>3</sub>Al Powders. *Powder Metallurgy*, 33(3):253–259, 1990.

- [10] H. R. Piehler and D. P. Delo. Physical modeling of powder consolidation processes.
   *Progress in Materials Science*, 42(1-4):263–276, 1997.
- <sup>554</sup> [11] S. Irukuvarghula, H. Hassanin, C. Cayron, M. M. Attallah, D. Stewart, and
   <sup>555</sup> M. Preuss. Evolution of grain boundary network topology in 316L austenitic stainless
   <sup>556</sup> steel during powder hot isostatic pressing. Acta Materialia, 133:269 281, 2017.
- [12] G. A. Rao, M. Kumar, M. Srinivas, and D. S. Sarma. Effect of oxygen content
  of powder on microstructure and mechanical properties of hot isostatically pressed
  superalloy Inconel 718. *Mater. Sci. Eng. A*, 435-436:84–99, 2006.
- [13] L. Chang, W. Sun, Y. Cui, and R. Yang. Influences of hot-isostatic-pressing tem perature on microstructure, tensile properties and tensile fracture mode of Inconel
   718 powder compact. *Materials Science and Engineering: A*, 599:186 195, 2014.
- <sup>563</sup> [14] C. F. Tipper. The fracture of metals. *Metallurgia*, 39:133–137, 1949.
- <sup>564</sup> [15] K. E. Puttick. Ductile fracture in metals. *Philosophical magazine*, 4(44):964–969, <sup>565</sup> 1959.
- [16] A. Pineau, A. A. Benzerga, and T. Pardoen. Failure of metals I: Brittle and ductile
   fracture. Acta Materialia, 107:424–483, 2016.
- [17] A. J. Cooper, N. I. Cooper, A. Bell, J. Dhers, and A. H. Sherry. A Microstructural
   Study on the Observed Differences in Charpy Impact Behavior Between Hot Isostat ically Pressed and Forged 304L and 316L Austenitic Stainless Steel. *Metallurgical* and Materials Transactions A, 46(11):5126–5138, 2015.
- <sup>572</sup> [18] A. J. Cooper, N. I. Cooper, J. Dhers, and A. H. Sherry. Effect of Oxygen Con<sup>573</sup> tent Upon the Microstructural and Mechanical Properties of Type 316L Austenitic
  <sup>574</sup> Stainless Steel Manufactured by Hot Isostatic Pressing. *Metallurgical and Materials*<sup>575</sup> Transactions A, 47(9):4467-4475, 2016.
- [19] C. L. Qiu, M. M. Attallah, X. H. Wu, and P. Andrews. Influence of hot isostatic
  pressing temperature on microstructure and tensile properties of a nickel-based superalloy powder. *Materials Science and Engineering: A*, 564:176 185, 2013.
- <sup>579</sup> [20] X. Pierron, A. Banik, G. E. Maurer, J. Lemsky, D. U. Furrer, and S. Jain. Sub<sup>580</sup> Solidus HIP Process for P/M Superalloy Conventional Billet Conversion. *Superalloys*<sup>581</sup> 2000, pages 425–433.
- [21] M. C. Hardy, B. Zirbel, G. Shen, and R. Shankar. Developing damage tolerance
   and creep resistance in a high strength nickel alloy for disc applications. *Superalloys* 2004, pages 83–90, 2004.

- [22] G. A. Rao, M. Srinivas, and D. S. Sarma. Effect of thermomechanical working on
  the microstructure and mechanical properties of hot isostatically pressed superalloy
  Inconel 718. *Materials Science and Engineering: A*, 383(2):201 212, 2004.
- [23] S. L. Semiatin, K. E. McClary, A. D. Rollett, C. G. Roberts, E. J. Payton, F. Zhang,
   and T. P. Gabb. Microstructure Evolution during Supersolvus Heat Treatment of a
   Powder Metallurgy Nickel-Base Superalloy. *Metallurgical and Materials Transactions* A, 43(5):1649–1661, May 2012.
- <sup>592</sup> [24] S. L. Semiatin, K. E. McClary, A. D. Rollett, C. G. Roberts, E. J. Payton, F. Zhang,
  <sup>593</sup> and T. P. Gabb. Plastic Flow and Microstructure Evolution during Thermome<sup>594</sup> chanical Processing of a PM Nickel-Base Superalloy. *Metallurgical and Materials*<sup>595</sup> Transactions A, 44(6):2778–2798, Jun 2013.
- <sup>596</sup> [25] V. Y. Gertsman and C. H. Henager. Grain boundary junctions in microstructure <sup>597</sup> generated by multiple twinning. *Interface Science*, 11(4):403–415, 2003.
- <sup>598</sup> [26] B. W. Reed, R. W. Minich, R. E. Rudd, and M. Kumar. The structure of the cubic
  <sup>599</sup> coincident site lattice rotation group. *Acta Crystallographica Section A*, 60(3):263–
  <sup>600</sup> 277, May 2004.
- [27] W. R. Bryan and M. Kumar. Mathematical methods for analyzing highly-twinned
   grain boundary networks. *Scripta Materialia*, 54(6):1029 1033, 2006. Viewpoint
   set no. 40: Grain boundary engineering.
- [28] C. Cayron. Multiple twinning in cubic crystals: Geometric/algebraic study and its application for the identification of the  $\Sigma 3^n$  grain boundaries. Acta Crystallographica Section A: Foundations of Crystallography, 63(1):11–29, 2007.
- [29] C. Cayron. Quantification of multiple twinning in face centred cubic materials. Acta
   Materialia, 59(1):252-262, 2011.
- [30] B. W Reed, M. Kumar, R. W. Minich, and R. E. Rudd. Fracture roughness scaling and its correlation with grain boundary network structure. *Acta Materialia*, 56(13):3278–3289, 2008.
- [31] C. M. Barr, S. Thomas, J. L. Hart, W. Harlow, E. Anber, and M. L. Taheri. Tracking the evolution of intergranular corrosion through twin-related domains in grain
  boundary networks. *npj Materials Degradation*, 2(1):14, 2018.
- <sup>615</sup> [32] R. Hielscher and H. Schaeben. A novel pole figure inversion method: specification <sup>616</sup> of the *MTEX* algorithm. J. Appl. Crystallogr., 41(6):1024–1037, Dec 2008.

- [33] C. Cayron. ARPGE: a computer program to automatically reconstruct the parent
   grains from electron backscatter diffraction data. Journal of Applied Crystallography,
   40(6):1183–1188, Dec 2007.
- [34] C. A. Schneider, W. S. Rasband, and K. W. Eliceiri. NIH Image to ImageJ: 25 years
  of image analysis. *Nat Meth*, 9(7):671–675, 2012.
- [35] S. Patala and C. A. Schuh. A continuous and one-to-one coloring scheme for misorientations. Acta Materialia, 59(2):554–562, 2011.
- [36] S. Patala, J. K. Mason, and C. A. Schuh. Improved representations of misorienta tion information for grain boundary science and engineering. *Progress in Materials Science*, 57(8):1383–1425, 2012.
- <sup>627</sup> [37] K. Miyazawa, Y. Iwasaki, K. Ito, and Y. Ishida. Combination rule of  $\Sigma$  values at <sup>628</sup> triple junctions in cubic polycrystals. Acta Crystallographica Section A: Foundations <sup>629</sup> of Crystallography, 52(6):787–796, 1996.
- [38] M. Shimada, H. Kokawa, Z. J. Wang, Y. S. Sato, and I. Karibe. Optimization of grain
   boundary character distribution for intergranular corrosion resistant 304 stainless
   steel by twin- induced grain boundary engineering. Acta Materialia, 50:2331–2341,
   2002.
- [39] C. A. Schuh, M. Kumar, and W. E. King. Analysis of grain boundary networks and
  their evolution during grain boundary engineering. *Acta Materialia*, 51:687–700,
  2003.
- [40] M. Michiuchi, H. Kokawa, Z. J. Wang, Y. S. Sato, and K. Sakai. Twin-induced grain
  boundary engineering for 316 austenitic stainless steel. *Acta Materialia*, 54:5179–
  5184, 2006.
- [41] Y. Hedberg, M. Norell, J. Hedberg, P. Szakalos, P. Linhardt, and I. O. Wallinder.
  Surface characterisation of fine inert gas and water atomised stainless steel 316L
  powders: formation of thermodynamically unstable surface oxide phases. *Powder Metallurgy*, 56(2):158–163, 2013.
- [42] V. N. Antsiferov and R. R. Yablonovskaya. Growth of dispersed oxide inclusions
  in 80% Ni 20% Cr alloys during annealing. Soviet Powder Metallurgy and Metal *Ceramics*, 10(6):482–484, Jun 1971.
- <sup>647</sup> [43] A. V. Serebryakov, T. M. Redkova, and V. I. Lobanov. On recrystallization of
  <sup>648</sup> dispersion-hardened alloys. *Physica status solidi (a)*, 14(1):77–81, 1972.

- [44] D. L. Williams. Hot Isostatically Pressed Alloy APK1, A Nickel-Base Superalloy.
   *Powder Metallurgy*, 20(2):84–89, 1977.
- [45] D. Furrer and H. Fecht. Ni-based superalloys for turbine discs. Journal of Metals,
  51(1):14–17, 1999.
- [46] G. A. Rao, M. Kumar, M. Srinivas, and D. S. Sarma. Effect of standard heat treatment on the microstructure and mechanical properties of hot isostatically pressed
  superalloy inconel 718. *Materials Science and Engineering: A*, 355(1–2):114 125,
  2003.
- [47] M. F. Ashby and I. G. Palmer. The dragging of solid particles through metals by
  grain boundaries. Acta Metallurgica, 15(2):420 423, 1967.
- [48] M. F. Ashby and R. M. A. Gentamore. The dragging of small oxide particles by
   migrating grain boundaries in copper. Acta Metallurgica, 16(9):1081 1092, 1968.
- [49] E. F. Koch and K. T. Aust. The movement of solid particles in a solid. Acta
   Metallurgica, 15(2):405 409, 1967.
- [50] G. I. Kokhanchik, A. V. Serebryakov, and Yu. A. Shiyanov. Diffusion movement of
   particles in solids. *Physica status solidi* (a), 23(1):99–103, 1974.
- [51] G. Gottstein and L. S. Shvindlerman. Theory of grain boundary motion in the
  presence of mobile particles. Acta Metallurgica et Materialia, 41(11):3267 3275,
  1993.
- <sup>668</sup> [52] R. K. Shelton and D. C. Dunand. Computer modeling of particle pushing and
   <sup>669</sup> clustering during matrix crystallization. Acta Materialia, 44(11):4571 4585, 1996.
- [53] Y. Oshida. An application of superplasticity to powder metallurgy. J. Jap. Soc. of
   Powder Metall., 22(5):147–153, 1975.
- <sup>672</sup> [54] C. Schuh, P. Noel, and D. C. Dunand. Enhanced densification of metal powders by <sup>673</sup> transformation-mismatch plasticity. *Acta materialia*, 48(8):1639–1653, 2000.
- <sup>674</sup> [55] M. Detrois, J. Rotella, R. L. Goetz, R. C. Helmink, and S. Tin. Grain bound<sup>675</sup> ary engineering of powder processed Ni-base superalloy RR1000: Influence of the
  <sup>676</sup> deformation parameters. *Materials Science and Engineering: A*, 627:95 105, 2015.
- <sup>677</sup> [56] R. C. Buckingham, C. Argyrakis, M. C. Hardy, and S. Birosca. The effect of strain
  <sup>678</sup> distribution on microstructural developments during forging in a newly developed
  <sup>679</sup> nickel base superalloy. *Materials Science and Engineering: A*, 654:317 328, 2016.

- <sup>680</sup> [57] Y. Gao, R. O. Ritchie, M. Kumar, and R. K. Nalla. High-cycle fatigue of nickel <sup>681</sup> based superalloy ME3 at ambient and elevated temperatures: role of grain-boundary
   <sup>682</sup> engineering. *Metallurgical and Materials Transactions A*, 36(12):3325–3333, 2005.
- [58] V. Randle. Twinning-related grain boundary engineering. Acta materialia,
   52(14):4067–4081, 2004.
- [59] M. Detrois, R. L. Goetz, R. C. Helmink, and S. Tin. Modeling the effect of thermal-mechanical processing parameters on the density and length fraction of twin
  boundaries in Ni-base superalloy RR1000. *Materials Science and Engineering: A*,
  647:157 162, 2015.
- [60] M. Detrois, J. McCarley, S. Antonov, R. C. Helmink, R. L. Goetz, and S. Tin.
   <sup>690</sup> Comparative study of high-temperature grain boundary engineering of two powder <sup>691</sup> processed low stacking-fault energy Ni-base superalloys. *Materials at High Temper-* <sup>692</sup> atures, 33(4-5):310-317, 2016.
- [61] J. McCarley, R. Helmink, R. Goetz, and S. Tin. Grain Boundary Engineering of a
   Low Stacking Fault Energy Ni-based Superalloy. *Metallurgical and Materials Trans- actions A*, 48(4):1666–1677, 2017.