

Quantifying microstructural evolution in moving magma

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32 **Keywords: volcanology, multi-phase, rheology, synchrotron radiation, *in situ*, bubbles,**
33 **magmatic processes, flow, X-ray tomography**

34 **Abstract**

35 Many of the grand challenges in volcanic and magmatic research are focused on understanding the
36 dynamics of highly heterogeneous systems, and the critical conditions which enable magmas to
37 move, or eruptions to initiate. From the formation and development of magma reservoirs, through
38 propagation and arrest of magma, to the conditions in the conduit, gas escape, eruption dynamics and
39 beyond into the environmental impacts of that eruption: we are trying to define how processes occur,

40 their rates and timings, their causes and consequences. However, we are usually unable to observe
41 the processes directly. Here we give a short synopsis of the new capabilities and highlight the
42 potential insights that *in situ* observation can provide. We present the XRheo and Pele furnace
43 experimental apparatus and analytical toolkit for the *in situ* X-ray tomography based quantification of
44 magmatic microstructural evolution during rheological testing. We present the first 3D data showing
45 the evolving textural heterogeneity within a shearing magma, highlighting the dynamic changes to
46 microstructure that occur from the initiation of shear, and the variability of the microstructural
47 response to that shear as deformation progresses. The particular shear experiments highlighted here
48 focus on the effect of shear on bubble coalescence, with a view to shedding light on both magma
49 transport and fragmentation processes; but the XRheo system is intended to help us understand the
50 microstructural controls on the complex and non-Newtonian evolution of magma rheology, and
51 therefore used to elucidate the many mobilisation, transport and eruption phenomena controlled by
52 the rheological evolution of a multi-phase magmatic flows. The detailed, in-situ, characterization of
53 sample textures presented here therefore represents the opening of a new field for the accurate
54 parameterization of dynamic microstructural control on rheological behavior.

55 **1 The need for *in situ* observation in magmatic research**

56 Magmas are generally a constantly evolving mixture of silicic melt, crystals and bubbles; with the
57 magmatic behavior during mobilisation, flow, fragmentation and eruption fundamentally controlled
58 by the evolution of the melt viscosity, and the changing volume of, and interactions between the
59 suspended crystal and bubble phases. Dynamics are key in volcanology. Most of our research focuses
60 on how magmatic systems have, are and will continue to evolve. We use a wide range of field,
61 laboratory, theoretical and numerical approaches, but understanding the dynamic and highly
62 heterogeneous nature of magmatic systems, defining the behavioural tipping points and identifying
63 the markers we should be monitoring remains difficult because we cannot fully interrogate or observe
64 magma at depth. The spatial and temporal heterogeneity of the system means that while empirical
65 relationships can be defined for a given set of conditions, the rheological evolution of magma
66 remains complex and understanding the magmatic and volcanic processes influenced by the rheology
67 requires a knowledge of the interaction microphysics that we currently lack.

68 The samples we collect in the field inherently contain a complex series of textural and chemical
69 overprints acquired during the entire evolution from formation to eruption/emplacement. The
70 behaviour of the magma at any point in time and space will have been influenced by evolution prior
71 to that point. Likewise, when we try to reproduce sub-surface magmatic conditions in the laboratory,
72 experimental charges capture only a small piece of that evolution, usually starting from an ideal
73 condition in order to capture a specific behaviour – a condition unlikely to be found in the natural
74 magmatic system. Our inability to observe processes under realistic conditions is a key limitation in
75 developing our understanding. In the natural system, many complex and interacting processes operate
76 simultaneously, and experiments generally aim to simplify and isolate phenomena. However, even if
77 we manage to capture and describe all the individual phenomena, there will still be interactions in the
78 natural system that are missed. The *in situ* approach with both natural and synthetic materials allows
79 us to see what those interacting phenomena look like in a new way.

80 **1.1 The need for *in situ* observation in rheological research**

81 Recent efforts into developing an understanding of the rheology of magmatic suspensions has
82 focused on deriving models that capture the rheology of magmatic suspensions as a function of
83 intrinsic (melt- and crystal-composition, texture etc.) and extrinsic (temperature, shear rate, etc.)

84 parameters (Campagnola et al., 2016; Chevrel et al., 2015; Ishibashi and Sato, 2007; Morrison et al.,
 85 2020; Pinkerton, 1994; Ryerson et al., 1988; Sato, 2005; Soldati et al., 2016; Spera et al., 1988; Stein
 86 and Spera, 1992; Vona, 2011; Vona and Romano, 2013). We understand that changes in magma
 87 viscosity and the transition from Newtonian to non-Newtonian behavior, observed during
 88 crystallization and vesiculation, depend fundamentally on the crystal and bubble content, shape,
 89 surface texture, and size distribution as well as the imposed strain rates (Mader et al., 2013).
 90 However, characterization of sample textures in rheological studies, which is crucial for the
 91 derivation of these rheological laws is, without exception, restricted to snapshots before and after an
 92 experiment.

93 While the importance of disequilibrium effects on crystal growth has inspired recent studies
 94 investigating the the dynamic rheology of crystallizing silicate melts at conditions close to those
 95 expected during emplacement in nature (Giordano et al., 2007; Kouchi et al., 1986; Ryerson et al.,
 96 1988; Vona and Romano, 2013), studies under disequilibrium conditions are few, and systematic
 97 descriptions of the effects of cooling (Giordano et al., 2007; Kolzenburg et al., 2019; Vetere et al.,
 98 2019), deformation (Kolzenburg et al., 2018b; Tripoli et al., 2019) and oxygen fugacity (Kolzenburg
 99 et al., 2018a), as well as their interdependence are limited. Textural characterization of experiments
 100 performed under these conditions has been impossible because the high rates of diffusion and crystal
 101 growth at high undercooling make it impossible to quench the experimental charges fast enough to
 102 preserve the textures under experimental conditions. As a result, dynamic changes in crystallization
 103 and vesiculation kinetics as well as sample textures during flow or during disequilibrium processes,
 104 and the associated responses in rheology remain unconstrained.

105 In recent years we have seen a dramatic increase in the ability to perform non-destructive observation
 106 *in situ* using the high-speed capability of X-ray tomography synchrotron facilities. The development
 107 has been largely led by the research needs of the material science community (e.g. Maire and
 108 Withers, 2014), but with the development of sample environments suitable for geological materials,
 109 and high temperature apparatus we are now capable of making certain key observations. The X-Rheo
 110 system described here now permits detailed, in-situ, characterization of these sample textures, a
 111 crucial factor for accurate parameterization of the rheological data, and hence opens up an entirely
 112 new field of study in magma rheology.

113 **2 Recent advances in synchrotron capability**

114 Several synchrotron facilities now have the capability to acquire 3D images of geological materials
 115 with moderate (10-5 μm) or high (<5 μm) spatial resolution. Several tomography beamlines can
 116 achieve fast (defined here as < 10 s to collect a single 3D tomographic dataset) tomography of these
 117 samples, and those focusing on imaging dynamic systems can routinely perform ultra-fast or real-
 118 time (defined here as < 1 s to collect a single 3D tomographic dataset) tomography on geological
 119 specimens (e.g. Marone et al., 2020). For 3D tomographic acquisition a frequency of up to 20
 120 tomographies per second (TPS) is now possible (Dobson et al., 2016). These same detector systems
 121 can achieve over 5000 fps when used to collect 2D radiography.

122 Imaging at sub-optimal conditions can lead to a reduction in data quality but can also deliver higher
 123 temporal resolution. This usually takes the form of reducing the number of angles from which the 2D
 124 projections are acquired, and/or under-exposing the images. With large volume onboard camera
 125 storage this can enable 10's or 100's of high frequency images to be acquired in a single
 126 experimental run. Cropped projection image areas can further increase the TPS (reducing data
 127 transfer time) and number of tomographies (smaller data volume per image) that can be collected.

128 Data transfer, rather than experimental speed, is now often the bottle neck. This has led to investment
129 in novel continuous read out detectors with real-time data transfer to network storage (Marone et al.,
130 2020; Mokso et al., 2017), and at some beamlines there is no theoretical limit to the number of 3D
131 datasets that can be obtained.

132 On most systems the sample can be made to rotate continuously in a single direction with acquisition
133 triggered at specific angle(s) by the rotational encoder. For the highest speed experiments the last
134 projection of the first data set (collected from 0-180°) can essentially be the first of the second dataset
135 (collected from 180-360°). Alternatively, acquisition “gapped” (e.g. Dobson et al 2016) can be
136 triggered every time the 0° encoder location is passed, or every nth time it is passed: allowing very
137 high-speed acquisition with a lower image frequency. This method has the added advantage that all
138 the reconstructed tomographies are in the same orientation, removing a registration step from the post
139 processing workflow.

140 **2.1 Recent applications to volcanology**

141 Areas within volcanology where real-time *in situ* synchrotron X-ray tomography could prove
142 transformative include (but are not limited to) bubble nucleation, growth and coalescence;
143 crystallization, crystal alignment and physico-chemical interaction, multiphase magma deformation
144 and rheology (incl. strain localisation reactive transport, and fragmentation), magma-rock-fluid
145 interaction, rock mechanics, etc. Work has now begun in some of these areas: bubble growth and
146 permeability (Baker et al., 2019; Baker et al., 2012; Colombier et al., 2018; Pleše et al., 2018), gas
147 driven filter pressing (Pistone et al., 2015), deformation (Okumura et al., 2013), sintering
148 (Wadsworth et al., 2016; Wadsworth et al., 2019), and crystallisation (Arzilli et al., 2019; Polacci et
149 al., 2018).

150 When coupled to advances in furnace capability (presented here) (Bai et al., 2008; Fife et al., 2012;
151 Kudrna Prasek et al., 2018; Polacci et al., 2018), and apparatus for tensional and compressional
152 deformation (Kareh et al., 2012; Philippe et al., 2016), experiments that elucidate how volcanic and
153 magmatic processes occur are now possible: we can observe and quantify their rates, we can
154 systematically explore the relationships that control the timing of these events, and start to define the
155 critical causes and consequences.

156 One area where this is critical, is in understanding magmatic mobilisation, mobility and flow.
157 Magmas are complex multiphase suspensions with non-Newtonian rheological properties. Testing
158 rheological behaviour of natural and synthetic samples is challenging in the laboratory (Caricchi et
159 al., 2007; Chevrel et al., 2015; Cimorelli et al., 2011; Lavallée et al., 2007; McBirney and Murase,
160 1984; Picard et al., 2011; Pistone et al., 2012; Pistone et al., 2013a; Pistone et al., 2013b; Vona et al.,
161 2011; Vona et al., 2016), but understanding how the microscopic behaviour of the suspended phases
162 are controlling the macroscopic behaviour remains even more so. As such, analogue materials offer
163 an excellent solution to scrutinise key rheological controls properties (Llewellyn et al., 2002; Mader
164 et al., 2013; Truby et al., 2014), not permissible during experiments on magma which is commonly
165 enclosed in furnaces and pressure vessels. However, even when working with well controlled
166 analogue materials, samples are generally opaque, and traditional methods cannot observe the nature,
167 or 4D distribution of the suspended phase interactions. Hence the microphysics controlling fluid
168 behaviour is not constrained at the scale of the processes, and rheological models remain incomplete.
169 *In situ* X-ray tomography is ideally suited to meet this kind of challenge, as it reveals previously
170 unrecognised phenomena and their temporal relationships.

171 *In situ* imaging of deformation experiments is a relatively new technique within the volcanological
 172 community, and performing these experiments has many unique challenges. We therefore give an
 173 overview of the considerations that must be made when designing or planning volcanological
 174 research using *in situ* imaging. We then present the XRheo system and associated high temperature
 175 furnace for performing *in situ* rheological experiments on magmatic samples: providing cutting edge
 176 capabilities for understanding real-time evolution of microstructures and how this affects rheological
 177 behaviour. We show the quality of the data that can be obtained and highlight the kind of salient
 178 features that can be observed, before presenting an implementation of new volume correlation
 179 analysis methods that enable the extraction of quantitative information on the displacements within
 180 deforming samples.

181 **3 Designing *in situ* experiments**

182 *In situ* imaging can be demanding, especially when working at high temperatures, or with complex
 183 experimental set-ups. Careful design to allow integration between synchrotron imaging systems and
 184 the *in situ* apparatus is critical to experimental success.

185 **3.1 Sample holders and encasement**

186 In the volcanic and magmatic context, the high temperatures needed for experiments under natural
 187 conditions is a key challenge. Many environmental cells used in standard laboratory experiments are
 188 made of substantial thicknesses of Pt or stainless steel. These materials have high X-ray attenuation
 189 coefficients and cannot easily be used with tomographic imaging. Many materials with low X-ray
 190 attenuation have limited thermal stability; cannot accommodate loading; or will react with the melt,
 191 changing the composition and behaviour of the sample during the experiment. Optimal imaging may
 192 result in deviation from optimal experimental conditions, or vice versa.

193 The need for encasement of molten samples using materials with low X-ray attenuation and good
 194 thermal properties makes boron nitride (BN) and Al₂O₃ the widely used materials of choice for high
 195 temperature experiments. BN is easy to machine and is stable to temperature of ~1000°C. Al₂O₃ can
 196 be used to higher temperatures (1200-1400°C) and has better mechanical properties under
 197 compression but is harder to machine and can still be susceptible to thermal shock. Graphite could
 198 also be used under reducing experimental conditions. For heating schedules typical for volcanic and
 199 magmatic experiments (10-20°C/min) the chance of fracture through thermal shock is low. With
 200 some natural melt compositions and environments, reaction between the BN or Al₂O₃ sample holder
 201 and the enclosed melt may still occur. Where off-line testing is able to demonstrate that some
 202 container-sample reaction is tolerable, the progression can usually be observed in real-time during
 203 acquisition, or defined *a posteriori* in the reconstructed 3D images. Non-reactive high (or low)
 204 temperature magmatic analogues may have to be substituted. In such cases, other experimental
 205 parameters (temperature, deformation rate) may also have to be adjusted to allow for investigation of
 206 specific phenomena. Other encasement materials such as thin Pt foil (Polacci et al., 2018) and boro-
 207 epoxy (Berg et al., 2017) are also possible, but make sample preparation more complex and can limit
 208 the range of mechanical experimental conditions that can be tested.

209 **3.2 Imaging moving magma**

210 When working with changing multiphase materials, easy quantification requires sufficient contrast
 211 between melt and the different mineral phases, while maintaining low enough overall attenuation to
 212 ensure high X-ray fluxes reaching the detector and with good signal quality. In all X-ray tomography
 213 the attenuation differences between the mineral phases should be carefully checked prior to scanning

214 (Hanna and Ketcham, 2017) to ensure the features of interest are visible, and with fast synchrotron
 215 tomography it is usually possible to use a monochromatic beam at an energy chosen to maximise the
 216 attenuation differences, or extract data at two energies bridging an absorption edge. For some
 217 minerals it is possible to determine mineral composition from calibrated reconstructed 3D data, even
 218 from a polychromatic “white” or “pink” beam (Pankhurst et al., 2018), but this should be considered
 219 before starting experimental work so beamline conditions and calibrations standards can be
 220 identified.

221 To achieve the fastest projection acquisition rates for real-time or ultra-fast synchrotron tomography,
 222 magmatic sample densities may limit users to a white or pink beam on the lower energy hard X-ray
 223 imaging beamlines (e.g. TOMCAT, Swiss Light Source; Marone et al., 2020), but still allow
 224 monochromatic imaging on higher energy beamlines (e.g. i12-JEEP, Diamond Light Source;
 225 Drakopoulos et al., 2015). The attenuation differences between most rock-forming minerals are
 226 greatest at lower energies (Hanna and Ketcham, 2017).

227 When working with dynamic systems the changes in total attenuation that will occur during the
 228 experiment (nucleation of crystals and associated changes in melt composition can cause a dynamic
 229 redistribution and therefore local and bulk attenuation) should also be considered, and the imaging
 230 parameters adjusted accordingly to allow adequate image quality at all stages of the experiment.
 231 Imaging crystal precipitation from a melt (Polacci et al., 2018), or crystals with a similar density
 232 (Arzilli et al., 2016) can be improved if phase contrast retrieval methods are used, but quantitative
 233 data analysis is still likely to be challenging and can be time-consuming and hard to automate. The
 234 enhanced phase boundaries achieved in phase contrast enhanced topography can make the small
 235 volumes of melt between crystals and/or bubbles harder to observe and quantify, and after
 236 preliminary testing was not deployed in this study.

237 **4 The XRheo and high temperature apparatus**

238 The XRheo system is designed to be compatible with almost all large enclosure laboratory scanning
 239 systems, and to function on all synchrotron imaging beamlines. Its development was possible by long
 240 term access to both the TOMCAT beamline at the Swiss Light Source and the I12-JEEP beamline at
 241 Diamond Light Source. The XRheo takes a standard high precision low-torque rheometric testing
 242 apparatus (Di Genova et al., 2016) and integrates it with the high temperature sample environment,
 243 and the imaging beamline.

244 The XRheo can perform the same suite of experiments that can be achieved during rotational
 245 rheological testing in the laboratory, but does so during 3D X-ray tomography acquisition; allowing
 246 the internal structural evolution to be captured *in situ* during the experiment while “traditional”
 247 rheological data are being collected. In standard rheological testing the lower plate or cup is a fixed,
 248 reference frame and the upper plate or spindle rotates, with the torque required to rotate at a given
 249 speed recorded. The XRheo uses the rotation needed for 3D image acquisition (i.e. a rotating cup) as
 250 the reference frame. In this study we use the XRheo in a wide gap concentric cylinder (Couette)
 251 configuration, but all other rotational rheological testing configurations are possible. The system
 252 (Figure 1A) is conceptually and operationally simple, but technically challenging to design and set
 253 up. To acquire 3D data during the deformation the XRheo uses the beamline rotation stage to control
 254 the cup. During the initial period the rheometer head is set to keep a constant value of zero load, and
 255 the coupled spindle therefore accelerates with and then rotates at the same velocity as the cup (to
 256 prevent development of any shear stress and internal deformation ahead of the experiment).

257 Once the system rotation is established and the desired experiment temperature reached, the beamline
258 rotation stage and cup continue at the same rotation speed, while the rheometer is initiated and the
259 spindle accelerated at a user-defined rate to a differential velocity with respect to the cup. As soon as
260 the acceleration starts the sample begins to experience internal deformation. Image acquisition from
261 before the start of deformation through to steady state deformation or a given number of relative
262 spindle rotations allows the microstructural evolution to be tracked.

263 Most rheometer systems allow the torque to be recorded throughout the experiment (both zero
264 deformation and acceleration/constant deformation parts of the sequence) and allow a series of
265 programmable accelerations and decelerations as required by the user. All equipment is safeguarded
266 by an over torque failure condition set at a load well below the maximum operating conditions of the
267 rheometer and beamline rotation stages. Should the load hit this threshold, the rheometer
268 automatically moves back to operating at zero load, returning to the zero-torque condition i.e. with
269 the spindle rotation the at the velocity as the sample.

270 The precise set-up is somewhat beamline-specific and gives flexibility to adapt to different sample
271 and equipment mounting systems, the size of the sample and the temperature range over which the
272 system will be operated. Installation and operation of the XRheo system can vary depending on need
273 and the auxiliary furnace equipment available, and the set-ups used at both TOMCAT and I12-JEEP
274 are described here.

275 **4.1 The XRheo**

276 The XRheo is highly adaptable and can operate with the rheometer head most suitable for any given
277 experiment, requiring only alteration of the fixing plate. Here we present data acquired using both
278 Anton Paar and Thermofisher air bearing (nNm sensitivity) rheometric testers. The head is mounted
279 on a set of high load manual translation stages that allows adjustment of the location of the rheometer
280 spindle over the centre of the beamline rotation stage. A second set of stages can be used to level the
281 spindle. The rheometer is controlled in the beamline control room via a USB connection to a laptop
282 and remote desktop control.

283 At TOMCAT (Figure 1A,B), the rheometer head and translation stages are mounted on a rigid
284 horizontal support bar, in turn mounted onto a vertical support bar, which is bolted to the base table.
285 A manual vertical translation stage is included at the bar intersection to allow the head and horizontal
286 bar assembly to be raised for sample changing and then lowered back into imaging position. At I12-
287 JEEP (Figure 1C, D) the rheometer mounting block is closer, attached to a split optical breadboard on
288 a frame around the sample stage for mounting custom sample environments. The same breadboard
289 supports the Pele furnace (see Figure 1C, D).

290 Alignment of the rotation axis of the spindle with that of the imaging stage is critical to both imaging
291 and rheological data quality in standard testing. While the tolerances in machining and the alignment
292 of the manufacturer supplied components may be adequate for a spindle or bob on a larger volume
293 laboratory rheological tester, it should be checked before use on a high-resolution imaging beamline.
294 A disposable BN or Al₂O₃ tip is attached to a steel spindle and attached the rheometer head (Figure
295 1D). The spindle tip diameters (and potentially geometry) can be varied according to the experiment.
296 The spindle is lowered to the imaging position with the tip close to the bottom of the field of view,
297 and its alignment checked. Initial alignment is achieved using the beamline rotation stage mounting
298 pin, raised to just below the spindle tip, as a visual reference. After that initial alignment, the X-ray
299 hutch is closed, and radiographs taken to quantify any alignment adjustment needed and assessed by
300 eye using multiple micrometer accuracy dial gauges.

301 4.1.1 Samples

302 Here we use two samples that have already been employed in experimental volcanology. Synthetic
303 magmatic samples formed from sintered soda lime glass beads of known rheological properties
304 (Wadsworth et al., 2017; Wadsworth et al., 2016; Wadsworth et al., 2019) and low temperature
305 magmatic analogues (oils & syrups). Other analogue materials could also be used, when appropriate
306 to address a specific experimental question. When needed, inert marker crystals are mixed into the
307 fluid. Here we use rutile crystals in the synthetic magma because they have significantly higher
308 density (attenuation) than the molten glass and can be considered nonreactive with the melt under
309 these experimental conditions over the time scale of these experiments. We use olivine or plagioclase
310 as the particles in the low temperature analogues (Figure 2). Initial tests showed that there is no
311 measurable settling of the crystals or rise of the bubbles on the timescale of the experiments.

312 4.1.1.1 Sinter in place

313 Some of the synthetic magmas were sintered in place. The beamline mounting pin (fixed to the
314 rotation stage) is then lowered and the sample cup mounted. The cups used here were BN, 5-8 mm
315 internal diameter, 0.5 mm wall thickness, and a 10 mm depth (Figure 3). Below the cup the ceramic
316 extended ~20 mm to a standard pin for attaching to the rotation stage. The cup is raised and aligned
317 in the field of view of imaging such that the base of the cup is ~2 mm below the bottom of the image,
318 and the rotation axis of the cup and spindle are aligned. Powdered samples were then loaded into the
319 cup using a pipette tip, taking care not to knock the spindle and to fill the cup evenly from all sides.
320 Once filled to the rim the cup was tapped gently to allow settling and increase packing tightness.

321 4.1.1.2 Pre-sintered or natural samples

322 Sintered-in-place, powder loaded samples at 1 bar typically exhibit a minor and variable amount of
323 porosity. Although the porosity can be carefully characterised prior to any deformation, pre-sintered
324 synthetic samples of known crystal and bubble content can often be produced with higher
325 reproducibility *ex situ*. Small cores cut from natural samples or larger blocks of pre-sintered material
326 can also be used in the XRheo set-up; however, it can be harder to ensure good thermal and
327 mechanical coupling between the cup, sample and spindle. Cores were ground slightly smaller than
328 the inner diameter of the cup to allow for thermal expansion of the melt during heating without
329 promoting large stresses which may fracture the cup. Finely ground (<60 μm) material of the melt
330 composition was then added to the cup to bridge the sample to the cup and encourage mechanical
331 coupling. In some samples smaller diameter drills were used to excavate a channel into the centre of
332 the core (Figure 1D) to accommodate the spindle, and additional ground material added over the
333 centre of the sample once the spindle was in place.

334 The best results were achieved when the spindle was raised 10 mm from the spindle imaging position
335 and the cup aligned as for the powdered samples (above). For the drilled samples, the spindle can
336 then be lowered back to image position prior to heating. For the un-drilled samples, the sample is
337 heated prior to the spindle being lowered into the molten material. Immersing the spindle into the
338 sample is done slowly to allow outward movement of the melt and accommodation of the spindle
339 without breaking the tip. Some samples fail to form continuous bubble free contacts between the
340 melt, the cup and the spindle. In these cases, large bubbles are usually visible in the reconstructed
341 data and the sample can be changed prior to performing the experiment. In some cases (including
342 some test samples which were loaded into the cups and heated *ex situ* before loading), the samples
343 appear to form a good mechanical contact but decouple from either the cup or spindle as soon as
344 deformation begins. The precise reason for this is unclear, especially when acceleration to
345 deformation speed is low and the sample is thermally equilibrated. It should be noted that the need

346 for rotation during heating means it is not possible to run solid undrilled samples at TOMCAT using
347 the laser system and the current XRheo set-up.

348 **4.1.1.3 Low temperature analogues**

349 The XRheo can also be used to perform experiments on low temperature magmatic analogues (e.g.
350 oils syrups, waxes). Loading these samples follows the pre-sintered protocol, with the fluids loaded
351 into the sample cup in the laboratory prior to transfer onto beam. While Al_2O_3 and BN cups and
352 spindles can be used for the high temperature experiments, BN cups and spindles are preferred when
353 using low temperature analogues because of the lower attenuation of the fluid phase. Use of an
354 encapsulation material that is denser than the sample generally makes image processing more
355 challenging. BN is also lower cost and easier to machine so more suitable for the shorter experiment
356 times (no long heating periods) when working with analogue systems.

357 **4.2 Working at magmatic temperatures on synchrotron beamlines**

358 Maintaining thermal stability across the sample volume can also be difficult in the configurations
359 required on a beamline. Heat losses are largely governed by the geometry of the experimental
360 apparatus and the heating system used. Some beamlines have their own furnaces for direct integration
361 with the beamline controls, and this integration permits accurate and complete capture of the
362 processes under investigation if they are strongly temperature dependent. In most cases furnace
363 volumes, hot zones or focused heating spots are relatively small and samples can experience a
364 substantial heat sink through the mounting pin. This effect can be seen, and quantified in some tests
365 such as during sintering experiments, where the powder material in contact with the BN mounting
366 plate sinters at a slower rate than the material in the central part of the furnace. The volume of the hot
367 zone and the effect of the heat loss on the experiment can sometimes be tested off beam, but in some
368 experimental geometries, rotation and rotation speed can have an effect that is challenging to test in
369 the laboratory.

370 At TOMCAT (SLS) the XRheo is compatible with the laser heating system (Figure 1) (Fife et al.,
371 2012; Marone et al., 2020). This system has 2 lasers mounted on opposite sides of the sample, each
372 generating a 6 mm x 4 mm portrait spot on the sample holder. The lasers are initiated after rotation
373 has begun, and rotation continues while the lasers are heating. The system works well with small
374 volume samples but high thermal gradients are present away from the laser spot (i.e. above and
375 below the imaging field of view, and may effect behaviours for larger samples. The furnace is
376 operated through the beamline control system.

377 **4.2.1 The Pele Furnace**

378 When working at I12-JEEP (DLS) existing beamline furnaces are incompatible with the XRheo and
379 instead we use the bespoke Pele furnace (Figure 3), fabricated by Severn Thermal Solutions, who
380 adapted a concept previously developed for a uniaxial deformation press (Lamur et al., 2018) and a
381 rotary-shear apparatus (Wallace et al., 2019). This rail mounted 1kW furnace has an operating range
382 of 200-1250°C, and the split design allows easy sample loading and a large hot zone (Figure 3A)
383 suitable for samples up to 20 mm in diameter. Thermal variability across the 30 mm (vertical) x 30
384 mm (across split) x 50 mm (along split) hot zone is better than 2°C and thermal stability is better than
385 $\pm 2^\circ\text{C}$. The furnace is powered and controlled by a Eurotherm programmable control unit (at present
386 operated independently of the beamline control system).

387 The split design allows the upper and lower insulation panels to be removed (Figure 3B) and replaced
388 according to the apertures required for any given experiment. The soft ceramic also allows for

389 insertion of additional thermocouples for monitoring internal temperatures if needed, or to allow
390 fitting of a gas line to allow heating under a controlled atmosphere. In the XRheo set-up, the upper
391 and lower insulation plates were drilled to give a central circular aperture, with the diameter of the
392 hole kept to ~1 mm larger than that of the cup (lower) or spindle (upper) (Figure 3B). To minimise
393 heat losses a second pair of insulation plates were placed on the closed furnace at 90 degrees to the
394 furnace closure (not shown in figure). The insulation pieces that allow X-ray transmission, holding
395 the X-ray windows upstream and downstream from the sample are equally versatile and can be
396 square cut (for tomography) or fanning outward (The Pele furnace has been used for high
397 temperature diffraction studies on i12-JEEP). It is also possible to insert a solid high transmission
398 window into the ceramic insulation should heat losses need to be minimised or atmospheric
399 conditions within the furnace volume controlled.

400 The height (55 mm), length (173 mm) and width (155 mm) of the furnace enclosure makes it highly
401 versatile. The system can be mounted above or below (as shown here, Figure 1 and Figure 3) the rail,
402 and the rail can be located upstream or downstream (as shown here, Figure 1 and Figure 3) of the
403 sample. When mounted below the rail, the thermocouple control and power cables for the furnace
404 connect to the bottom of the unit, allowing space above for user access or additional equipment. The
405 furnace rail can be mounted on the optical breadboard that supports the XRheo or can be integrated
406 with other *in situ* apparatus. The furnace has been successfully mounted on the Deben open frame 10
407 kN deformation cell. It can also be used in a vertical configuration, hanging from the rail or a custom
408 mounting system. This makes it compatible with laboratory scanning systems. In the vertical
409 orientation the ram apertures and X-ray windows are switched. This enables source-sample and
410 sample-detector distances to be reduced (minimum of 30 mm) to the distances necessary for high
411 resolution laboratory scanners, although additional heat protection for the source and detector would
412 be needed and would increase this distance slightly.

413 With the heat sinks of the XRheo cup and rams in place, the sample and spindle rotating, and
414 additional air cooling focused on the upper part of the spindle to prevent overheating of the
415 rheometer head air bearing (Figure 1) there is a reproducible differential between control loop
416 temperature and sample temperature of ~50°C for sample temperature of 500-1200°C. This was
417 calibrated from the rate of sintering in the boro-silicate bead pack, as we have thorough
418 understanding of, and accurate control on, sintering (Wadsworth et al., 2017). This differential will
419 vary with sample and cup materials as should be calibrated for a different experimental set-up.

420 **4.2.2 Heating Schedules with the XRheo**

421 On both beamlines the hot zone of the furnace system is centered on the imaging field of view. The
422 samples were heated to the chosen dwell temperature at 10-20°C/min and held there until the end of
423 the experiment. In cases where solid samples were loaded the sample was heated to temperatures
424 1000°C, usually much higher than target experimental temperatures to increase the speed of
425 equilibration around the spindle and promote coupling. At the end of the experiments the temperature
426 can be ramped back down or quenched by turning off the furnaces and allowing the samples to cool
427 in place before sample change. All cups and ceramic spindle tips are single use consumables for all
428 high temperature experiments. The two high temperature experiments described here (Sections 6.1
429 and 6.2) were performed at a dwell temperature of 900°C.

430 **5 Image acquisition**

431 While the tomography acquisition condition can be changed for each sample, this increases the
432 chance of human error, especially when calculating rheometer control in terms of relative difference

433 in rotation speed. Therefore, for simplicity we define image acquisition parameters (rotation speed,
 434 exposure time and project number) at the start of the experiment, using a high-density sample.
 435 Reducing the exposure time, or even under-exposing, generates more tomographies per second (TPS)
 436 and less motion blur, but this will eventually cause a level of signal to noise that prevents post-
 437 processing and accurate analysis of the data. Reducing the number of projections acquired per
 438 tomography while keeping slightly longer exposure times can also increase the tomographies per
 439 second that can be collected, but again, can lead to degradation of the data quality (although this may
 440 be preferable with the latest iterative reconstruction algorithms, see below). The acceptable limit of
 441 quality reduction must be defined on a case by case basis.

442 Typical acquisition schedules showing the link between the beamline, deformation control and
 443 heating is shown in Figure 4. This structure was used on the TOMCAT and i12-JEEP beamlines to
 444 acquire the data presented below.

445 We present data collected at TOMCAT for which optimal conditions for high temperature
 446 borosilicate melt plus up to 30 vol % rutile crystal cargo were 1s per tomography, collecting 1000
 447 projections per tomography (180° rotation) with an exposure time of 1ms. The data were acquired on
 448 the GigaFRoST detector (Mokso et al., 2017), with a 3.7 μm reconstructed voxel (3D equivalent to a
 449 pixel) resolution. Up to 150 tomographies were collected in each run. At 1TPS there was negligible
 450 motion blur in the reconstructed data except where coalescence was occurring. Faster rotation is
 451 therefore unnecessary for the lower deformation speeds from 0.01 - 1 rotations per minute (rpm)
 452 differential speed (equivalent to strain rates of 10⁻³ to 10⁰ s⁻¹ at the spindle surface, within the range
 453 of those reported for natural systems and experimental studies (Chevrel et al., 2015; Kendrick et al.,
 454 2013; Kolzenburg et al., 2016; Mueller et al., 2011; Pistone et al., 2012; Vona et al., 2011). To avoid
 455 unnecessary post-experiment data processing steps the projections were always collected over the 0-
 456 180° arc of the rotation, no data were collected from 180-360°. The magnitude of the deformation at
 457 the lower strain rates will require less frequent, but not slower image acquisition. The tomography
 458 acquisition rate (TAR) for the data presented (0.15 and 0.3 rpm differential speed) was 0.07/s (1
 459 tomography every 7 rotations, every 14 seconds, ~4 tomographies per minute). The data were
 460 reconstructed using the standard TOMCAT pipeline.

461 We also present data collected at i12-JEEP, for which optimal conditions for high temperature
 462 borosilicate melt plus rutile crystal cargos were 0.25 to 0.5 s per tomography (depending on rutile
 463 content). We collected 720 projections per tomography (180° rotation) with an exposure time of 150-
 464 300μs. The data were acquired using the MIRO 310M camera using module 3 magnification giving a
 465 7 μm reconstructed voxel resolution to increase maximum FPS. For the low temperature analogue
 466 materials acquisition times reduced to 0.125 per tomography, giving a maximum of 8 TPS, and 360
 467 projections per tomography. This enabled imaging of the lower viscosity analogues with deformation
 468 speeds of up to 10 rpm differential speed. The projections were always collected over the 0-180° arc
 469 of the rotation. The scanning was performed at a minimum sample-detector distance to minimise the
 470 phase contrast which tended to make the bubble bearing samples harder to process. The gap used
 471 depended on the experiment and deformation rate. Data were downloaded between each run, and in
 472 each run a maximum of between 43 and 75 tomographies could be collected depending on the height
 473 of the cropped image. The TAR for the high temperature data presented was 0.07/s (as for
 474 TOMCAT) for the same sample deformation conditions. For the low temperature analogues, and
 475 faster deformation rates, the maximum TAR was 1/s (1 tomography every 4 rotations, every 1
 476 seconds). The data were reconstructed using the standard i12 pipeline.

477 **6 4D observations: deformation and strain localisation during bubbly magmatic flow**

478 The focus of this paper is technical and methods innovation, and the application of these methods to
 479 experiments relevant to volcanological research. Thus, the XRheo system as used here is designed to
 480 simulate the conditions of volcanic flow. Detailed analysis of the data and subsequent interpretation
 481 will be presented elsewhere. A selection of images and features that show the capability and
 482 sensitivity of the method are highlighted in the following sections. 3D and 4D image data are
 483 challenging to present on a page and here we show a series of 2D slices through the reconstructed
 484 volume, along with the relationship of that slice to the data (Figure 5). No difference is seen in the
 485 general behaviour across the height of the sample, and the spindle tip extends well below the field of
 486 view. In subsequent images all 2D slices are the mid-point slice (relative equivalent to the red slice
 487 Figure 5), and all images have been overlain with a 7x7 grid to aid description of the feature locations
 488 and visualisation of the displacements.

489 All data have been visualised and processed in Avizo© 2019.2 (ThermoFisher™). The data shown
 490 have received minimal post processing. The data in Figure 5, Figure 6, and Figure 7 have been
 491 cropped, down sampled (voxel size now 7.4 µm) and filtered using the Anisotropic Diffusion filter (5
 492 iterations) with the default settings as defined by the algorithm from the greyscale distribution. The
 493 data in Figure 8 and Figure 9 have not been processed. These are the raw reconstructed data.

494 **6.1 Bubble Bearing Magma (high temperature)**

495 Figure 6 shows the evolution and displacement of the bubbles in the middle slice through time. The
 496 images are taken from the acceleration phase of the deformation cycle. Deformation starts in the 5th
 497 frame collected (#5) and it reaches 0.3 rpm differential speed at #80. The spindle is moving in a
 498 clockwise direction, faster than the cup, which is stationary in the image reference frame. Three
 499 bubbles have been picked out with coloured labels (red, blue, yellow). These highlight different
 500 behaviours. The red bubble, close to the cup (B6, #40) is displaced along a broadly circular path (C7,
 501 #47), but as deformation increases, we start to observe an additional component of outward radial
 502 displacement (D7, #54). The blue bubbles (E1, E2, #41) undergo coalescence between #42 and #43,
 503 then remain nearly stationary for a few frames while the volume around them responds to the change
 504 in local stress following the coalescence. The bubble then starts to move and elongate (#42-48) then
 505 expands possibly by vertical translation through the plane of the image, or by coalescence with
 506 bubbles out of the plane of the image as the bubbles surrounding the blue bubble are not captured.
 507 The yellow bubble is almost stationary until frame #46, but then follows a circular displacement path,
 508 but the bubbles around it showing coalescence and significantly different displacement rates (not the
 509 differences in bubble structure in the region between the red and yellow bubbles (#50-#56).
 510 Substantial microstructural changes are occurring even at the onset of deformation.

511 Figure 7 shows the same sample much later in the deformation, during the final part of the
 512 acceleration, and the initial stages of the constant rate deformation. The bottom three frames are
 513 during the constant deformation stage. Much larger bubbles have formed through coalescence (as
 514 seen relative to Figure 6), and coalescence continues through the experiment. The higher deformation
 515 rate is reflected in the greater elongation of the bubbles, and we see more evidence for shear
 516 localisation.

517 Five bubbles have been identified in #60 and tracked through the time series, although the same
 518 behaviours can be seen throughout the volume. The blue bubble increases in area from #60 to #65 by
 519 movement through the plane of the image and not by coalescence. The blue and red bubbles coalesce,
 520 and in #69 we can see the coalesced bubble undergoing relaxation. A higher TAR would have
 521 enabled capture of the coalescence events in more detail. The yellow bubble coalesces with the

522 bubble trailing it (not coloured) between frames #65 and #66, but does not coalesce with the red, or
523 red-blue bubble pair despite the film thickness between the yellow-red pair and the red-blue pair
524 being similar and tall three bubbles being close in volume. The green and pink bubbles are at the
525 same radial distance but are travelling at different speeds, although both are travelling faster than the
526 yellow-red-blue bubble train. The green bubble “catches up” to the pink, and then both approach a
527 third bubble of similar size and radial position that is moving at an even slower rate. Throughout the
528 volume we can see short lived apparent slip surfaces developing, and the formation of a “nested
529 annular” structure.

530 Similar to the yellow-red-blue bubble train, two of the three large bubbles in C1, D1 and E1 at #60,
531 coalesce between #62 and #63, but despite having a very thin film between them, do not coalesce
532 until the sample moves from the accelerating to the constant deformation stage. These bubbles are in
533 much closer contact than the red-blue pair that undergo coalescence (#68-#70). Many of the bubbles
534 also show deformation that gives an apparent rotation inward toward the spindle, and in some
535 instances the bubbles develop a concave geometry on the spindle-facing bubble surface.

536 The 2D time series data displayed in Figure 6 and Figure 7 demonstrate how much deformation can
537 occur at even low strain rates and low total strain. No coalescence was observed when the sample
538 was left undeformed over the same timescales. The bubble number and bubble size distribution at the
539 start of the constant deformation (e.g. #84, Figure 7) is significantly different than at the start of the
540 experiment (cf. Figure 6). Images from later during the constant deformation continue to show
541 coalescence and elongation throughout the experiment.

542 This sample will be used to illustrate the more complex quantitative methods applied to the XRheo
543 data set (Section 7).

544 **6.2 Three-Phase Magma (high temperature)**

545 An example of the complexity in 3 phase (melt + crystals + bubbles) deformation can be seen in
546 Figure 8. This sample contains a 30 vol % volume crystal load. In this small subset of the
547 acceleration phase the spindle has accelerated to approximately half the final differential speed of
548 0.07 rpm. Translation and rotation of crystals is apparent throughout the volume, and the bubbles are
549 both more highly deformed, and undergoing less coalescence than in the crystal-free, bubble bearing
550 magma (Figure 6 and Figure 7). The high density of the rutile grains causes some beam hardening
551 and minor streak artefacts visible in some parts of the data. The formation, deformation and breaking
552 of chains of crystals can be seen. In the region surrounding B6 and C6 crystals are “rolled” along
553 (migrate with the bubble and show some rotation) one surface of the bubble while the crystals on the
554 opposite surface remain un-affected.

555 **6.3 Low temperature analogue (room temperature)**

556 Low temperature analogues are widely used to allow more systematic testing of parameter space, and
557 they often allow more control in the sample preparation or exploration of systems that are unstable in
558 laboratory conditions. The XRheo has been tested with the Canon N2700000 high viscosity standard
559 mineral oil, working at higher deformation rates than used for the magmatic samples above. The data
560 presented in Figure 9 show a short part of a steady state 5 rpm differential speed deformation step
561 within a longer experiment. These data do show a small, and variable amount of motion blur but this
562 does not prevent qualitative, and some quantitative analysis. Like in Section 0, this sample has 30
563 volume % crystals, but the lower viscosity of the sample and the more irregular shapes of the olivine

564 crystals allow observation of crystal rotations. Strain localisation is observed but it is generally less
565 focused and less long-lived than seen in the magmatic samples.

566 7 Quantifying displacements and heterogeneity in flowing magma

567 The spatial heterogeneity within magma, and the highly dynamic (e.g. changes in crystal and bubble
568 fraction, melt viscosity) nature of the system during storage, transport and eruption makes it
569 extremely challenging to understand how the microstructural evolution effects the macroscopic
570 properties. Understanding this control is the objective of many experiments in the volcanic and
571 magmatic studies that work to quantify how features move, how deformation is accommodated, how
572 bubbles and crystals interact and how textures evolve through time. Ultimately, many of the key
573 magmatic and volcanic processes we wish to understand, such as mobilisation, mixing, localisation,
574 flow, degassing, fragmentation, welding, and many more, are all controlled by the microstructure
575 (and how it is changing) at the time that the process is operating.

576 Understanding these processes therefore requires thorough *in situ* analysis while those processes are
577 operating under realistic conditions. The experiments presented here focus on using the XRheo to
578 understand bubble deformation, flow localisation and bubble interactions during magmatic flow,
579 focussing on the processes of coalescence and degassing. The data show the processes occurring, but
580 if these data are to be used to build models to predict real world behaviours a quantitative
581 understanding of the microphysics of the evolving magma is needed. Here we use the 4D XRheo
582 XCT data to develop a new implementation of recently developed Digital Volume Correlation
583 (DVC) algorithms that enable the local displacements and strains to be calculated, and the causes and
584 consequences of the structural and rheological heterogeneity within samples to be quantified. DVC is
585 an optical flow-based image matching technique that can measure displacement and strain fields
586 between two 3D image datasets, and can measure displacements smaller than the voxel size (Bay,
587 2008; Buljac et al., 2018; Hild et al., 2016; Sloof et al., 2016; Tozzi et al., 2014).

588 7.1 Different DVC Approaches

589 Here we apply a novel DVC approach to the bubble bearing magma discussed in Section 6.1, and
590 shown in Figures 6 and 7.

591 The established “local” subset-based DVC approach divides the reference and the deformed images
592 into smaller volumes that are then individually correlated (Bay et al., 1999; Buljac et al., 2018; Madi
593 et al., 2013; Smith et al., 2002). This method generally yields displacement fields with small random
594 errors (high precision) but high systematic errors (low accuracy), because the sub-volumes are treated
595 independently and there is a loss of displacement continuity at the boundaries. The more recently
596 developed “global” DVC approach assumes that a measured displacement field is continuous, and so
597 yields more accurate and robust results close to the solution, but this approach is more
598 computationally costly. Here we employ a protocol which uses the “local” approach to define an
599 initial displacement field, which is then used as the initialisation for the more robust global DVC.

600 Although the data are collected over the same angular interval, and are therefore pre-registered, the
601 location of the cup and spindle drift slightly within the image volume over the time series. To
602 interpret the displacement field within the melt, this rigid body motion was removed using the
603 *Registration* algorithm in Avizo©, prior to running a DVC analysis. A single melt volume was then
604 defined, and this “sample mask” used to both remove the cup and spindle from the images, and
605 generate the 3D meshes that define the individual elements for the DVC.

606 7.2 Quantitative analysis of displacements on XRheo 4D data

607 The overall workflow is presented in Figure 10. First, the “local” approach was applied to a simple
 608 cubic mesh defined within the DVC algorithm in Avizo©. The DVC algorithm then calculates the
 609 displacement field, and this is used to initiate the “global” approach across a coarse tetragonal mesh
 610 (Global Course (G_C) mesh elements 79 voxels) generated from the sample mask (see Figure 10 for
 611 detail of the steps undertaken in this method). The displacement field output from that correlation
 612 (global course mesh, images 1 & 2 = G_C^{1-2}) is then used as the initialisation for G_C^{2-3} . The output
 613 from G_C^{2-3} used as the initialisation for G_C^{3-4} and on through the entire time series (Figure 10).

614 The heterogeneity observed by this process suggests that the G_C mesh was smoothing some of the
 615 smaller scale displacements, and so a second iteration using a finer mesh (G_M) with mesh elements of
 616 55 voxels was then completed. Here G_C^{1-2} was used to initialise G_M^{1-2} , and then G_M^{1-2} used to
 617 initialise G_M^{2-3} and on through the data (Figure 10). Throughout this process the G_C and G_M outputs
 618 are broadly similar and show no evidence for divergence. At several time points the $G_M^{X-(X+1)}$
 619 generated from $G_M^{(X-1)-X}$ was compared to that generated from $G_C^{X-(X+1)}$. While there are some small
 620 differences, the overall variations in the displacement field are small, confirming the validity of
 621 applying the G_M (Figure 10). At the end of the processing the DVC outputs were rendered as a
 622 volume colour wash over the greyscale volumes. This comparison showed that some small-scale
 623 deformations were still not captured. A finer global mesh (G_F) iteration with mesh elements of 35
 624 voxels was then completed (up to #69 when the displacements become too large for the finer mesh
 625 and the data diverge). Finer meshes were not generated as these would generate a high number of
 626 mesh elements without textural information preventing accurate tracking. For each DVC output the
 627 displacement fields can be visualised as a 3D colour wash volume, or as a vector field (Figure 11,
 628 Figure 12) shown here with the vectors colour coded by displacement magnitude. Tangential and
 629 radial magnitudes can also be extracted.

630 7.3 Observing heterogeneity in rheological test samples

631 During the initial acceleration (G_F^{8-9} - G_F^{15-16} , Figure 11) we see a gradual radial propagation of the
 632 displacement field. As the deformation continues, we start to see the development of heterogeneity in
 633 the displacement field. Initially the heterogeneity is angular, with some sectors showing faster motion
 634 than others, but as this develops some vertical variability also appears. At this point in the series (see
 635 also Figure 6), coalescence and bubble deformation are limited, but are observed. The sectoral nature
 636 of the heterogeneity might suggest that the cause may be a misalignment between the rotation axes of
 637 the spindle and the cup. However, if this was the case, we should expect the sector showing the
 638 maximal deformation to rotate with the spindle, tracking the narrowest gap distance. This is not
 639 observed, and the region showing the greatest displacement shows no systematic larger scale
 640 migration. In fact, the high deformation region is generally fixed at one part of the sample volume, or
 641 fluctuates between 3 o'clock and 12 o'clock regions (as the data are displayed in Figure 11 and
 642 Figure 12). From G_F^{25-26} we start to see an increase in the magnitude of the displacement in two
 643 regions. From G_F^{38-39} there is a rapid increase in displacement in one sector of the sample, with a
 644 noticeable lack of a radial position dependence on the displacement speed. Once coalescence
 645 becomes more common in the volume from about G_F^{45-46} onwards (see also ~#45 onwards, Figure 6)
 646 we see another rapid increase in the magnitude of the displacements. There is substantial fluctuation
 647 however, especially from G_F^{48-49} to G_F^{52-53} , and G_F^{60-61} to G_F^{63-64} part of the series, where we see a
 648 maximum displacement (~50 voxels) well away from the spindle. This displacement is substantially
 649 higher than the velocity of the spindle surface, and corresponds to the major coalescence event shown
 650 in #61-#64 (Figure 7).

651 As the 0.3 rpm differential speed is reached at #80 (Figure 7), the system might be expected to
652 equilibrate and the displacements to become more homogeneously distributed. However, the ongoing
653 coalescence and the established heterogeneity and focused displacements continue (Figure 12). There
654 is little change immediately after the acceleration ceases (#79 to #90) despite the ongoing
655 coalescence events and changes in the bubble distribution that can be seen (Figure 7). Also, at this
656 time we start to see larger local decelerations in the region opposite the highest displacements (e.g.
657 #69, #74, #104) suggesting there may be some far field effects becoming visible once the
658 deformation rate is constant. Clearly the data are showing an extremely complex behaviour, and
659 suggest that any bulk viscosity measurement on this sample may contain contributions that are
660 intrinsically tied to the heterogeneity and localisation that is present.

661 The DVC approach presented here has been successfully applied to all three sample types presented
662 in this study (crystal bearing, bubble bearing and three-phase), and has the potential to be useful for
663 most *in situ* magmatic analysis, provided that the geometrical changes (bubble deformation) and
664 displacements are not extreme between the frames being correlated. Better understanding of
665 rotational crystal motions is achieved with non-spherical particles. More traditional versions of the
666 DVC method have been used in other geological systems (e.g. McBeck et al., 2018) but approach
667 taken here could also make the approach invaluable to experimental volcanology as well, where
668 heterogeneity and local environment can be critical controls on behaviour.

669 **8 Broader challenges and future directions**

670 There are always improvements to be made, and in tomographic analysis, better quality data can
671 significantly reduce processing and analysis duration. In these experiments motion blur still exists in
672 some regions of some images, showing that we may need yet faster FPS and similar overall
673 experiment durations (meaning an increase in data volume) if we want to investigate behaviours
674 higher strain rates or lower viscosities.

675 Faster exposure times result in less motion blur, but will generate projection data with high levels of
676 noise; making data harder to analyse. Advances in iterative reconstruction methods can provide
677 increased data quality from reduced number of projections, and therefore mitigate this challenge.
678 Methods that use “structural prior” to enhance the reconstruction processes can yield data of
679 comparable quality from 10x fewer projections. These methods take a high spatial-low temporal
680 resolution data set collected prior to the dynamic experiment to provide additional structural
681 information and tune the reconstruction of the under-sampled (low projection numbers) and high
682 noise data from the high TPS part of the experiment (Kazantsev et al., 2014; Kazantsev et al., 2015a;
683 Kazantsev et al., 2015b). Where areas of the image volume remain largely unchanged (e.g. fluid flow
684 in porous media experiments and some crystallisation experiments) data from non-changing areas in
685 the high TPS part of the experiment can be used instead of that from a high spatial-low temporal
686 resolution data (Eyndhoven et al., 2015). Where contrasts are high and the features of interest are
687 large relative to the voxel size, this could result in data of sufficient quality for quantitative analysis
688 from as little as 18 projections (Eyndhoven et al., 2015). Applying iterative algorithms can be used to
689 achieve three objectives: i) increase data quality for a given FPS/projection number acquisition
690 protocol, ii) enable higher FPS to reduce motion blur and experiments with higher strain rates, or iii)
691 allow longer duration experiments where there are hardware-based limitations on the total number of
692 projections that can be acquired. These iterative methods are becoming routinely available at high-
693 speed imaging beamlines, but may not improve data quality in all experimental scenarios.

694 Another way to improve data quality is to increase the signal to noise in the projection data, usually
695 by increasing the X-ray flux reaching the detector. Many synchrotron facilities are planning or
696 implementing upgrades to existing infrastructure that will deliver this, with increasing beam
697 brilliance (X-ray flux) key among many of the development plans. More brilliance could be used to
698 increase the TPS rate maintaining current data quality, or to obtain better quality data for the same
699 TPS.

700 The experiments shown here are of shorter duration than some *ex situ* laboratory rheological tests,
701 and shorter than is necessary for rheological analysis during crystallisation: so longer experiments
702 may be needed. DVC algorithms and other particle tracking techniques require relatively small
703 displacements between images, so we cannot achieve longer experimental durations by reducing the
704 TAR. More data are required. Even with current data collection protocols, high speed acquisition can
705 generate vast amounts of data (TB/day) especially when the processes under investigation are poorly
706 understood. Experiments need to start collecting images before the start of the processes, and in some
707 cases the occurrence of the processes themselves are hard to identify until after the reconstruction
708 step. This demonstrates the need to reconstruct ever increasing amounts of data before leaving the
709 synchrotron, the need to transfer ever larger data volumes, and to develop more automated processing
710 algorithms.

711 The final challenge to *in situ* experiments is always in the amount of post processing and data
712 analysis required to interpret the results. The DVC approach applied here removes the need for
713 segmentation or labelling of individual phases, particles or bubbles, but in most experiments there
714 will be a need for an additional image processing workflow to extract quantitative information. Rapid
715 advances in automated and machine learning algorithms for application to tomographic image
716 analysis, but these can still be challenging to apply to low contrast geological materials, and can be
717 more difficult to implement when changes in sample structure (e.g. clustering of crystals, growth of
718 bubbles) affect local or bulk attenuation. Further work in this area is needed before fully automated
719 image processing workflows can be achieved on most geological materials.

720 The data presented here showcase the capabilities we currently possess, and shows how real-time *in*
721 *situ* tomography can provide useful and extensive information about volcanic and magmatic
722 processes of flow and deformation. The data here are far from optimal by the standards typical for
723 high (spatial) resolution laboratory or synchrotron acquisition but are still more than sufficient to
724 perform complex quantitative analysis of dynamic processes. As multi-scale and correlative imaging
725 techniques (linking to other 2D, 3D, and 4D imaging modalities) become more widely available, the
726 incorporation of data from phase contrast imaging (grain boundaries/crystallisation), k-edge imaging
727 (chemical information), simultaneous CT and XRD, and diffraction imaging and diffraction contrast
728 tomography will continue to increase the detailed understanding of microstructural evolution that can
729 be extracted from such *in situ* studies. This, coupled to the continued widening of laboratory scanner
730 access, a diversification and expansion of the *in situ* apparatus available within the community, and
731 plans for beams of higher brilliance at several of the European synchrotron facilities in years to come,
732 means the future of *in situ* magmatic and volcanic flow research looks bright.

733 **9 Conclusions**

734 The XRheo is one example of the rapid development of new experimental apparatus that has been
735 driven by the expansion of *in situ* and real-time X-ray computed tomography. When coupled to the
736 Pele furnace, it enables *in situ* observation of the internal microstructural evolution of samples during
737 rheological experiments. By using the appropriate analogue, synthetic and natural materials as

738 experimental charges, and operating under a range of temperature and deformation conditions, it
739 allows us to interrogate the processes occurring in magmas moving in the shallow crust and on the
740 Earth's surface. Studies assessing the microstructural controls on rheological behaviours, and the
741 critical microstructural conditions that enable, accelerate or arrest flow are now being undertaken. By
742 implementing the latest Digital Volume Correlation methods to allow quantitative assessment of
743 displacements and strains and the heterogeneity in the displacement fields, we provide an exemplar
744 workflow for the level of quantification that can now be achieved. These tools provide the ability to
745 significantly improve our understanding of the key processes that control the behaviour of natural
746 magmas. Specifically, they permit the quantitative understanding of the microstructural changes *in*
747 *situ* while also recording traditional rheological data. This gives us the ability to define, for the first
748 time, how dynamic microstructures (bubble volume and geometry, bubble coalescence or collapse,
749 changes in local crystal content, growth or dissolution of crystals, and the formation and
750 fragmentation of load bearing crystal networks and other localisation features) effect key
751 rheologically controlled processes such as percolation, mobilisation, flow, degassing, and
752 fragmentation.

753 While computed tomography has become an established technique in earth sciences, it is the recent
754 advances at synchrotron beamlines and the environmental cells and apparatus used to perform *in situ*
755 experiments that now allows the volcanology community to exploit the technique to its maximum
756 potential: tackling some of the greatest challenges in the field and addressing the dynamics of
757 magmatic processes in new ways.

758 **10 Conflict of Interest**

759 The authors declare that the research was conducted in the absence of any commercial or financial
760 relationships that could be construed as a potential conflict of interest.

761 **11 Author Contributions**

762 K.J.D developed the XRheo and the Pele Furnace, managed the beamline experiments, the data
763 acquisition and analysis and led the writing of manuscript. All team members (listed alphabetically)
764 contributed to acquiring and processing the data on which this work is based, and contributed to the
765 preparation of the manuscript.

766 **12 Funding**

767 KJD was supported by NERC M018687/1 and ERC 2009 ADV Grant 247076 (EVOKES). DBD
768 acknowledges the support of ERC 2009 ADV Grant 247076 (EVOKES) and ERC 2018 ADV Grant
769 834225 (EAVESDROP) during the conduction of these experiments and the preparation of the
770 manuscript, respectively. RC, JEK, YL, JS, FWvA were supported by ERC 2012 StG Grant (SLiM)
771 406388. SK acknowledges the support of a H2020 Marie Skłodowska-Curie fellowship
772 DYNAMOLC – No.795044.

773 **13 Acknowledgments**

774 We acknowledge the Paul Scherrer Institut, Villigen, Switzerland for provision of synchrotron
775 radiation beamtime at the TOMCAT beamline X02DA of the SLS under proposal 20150413, and
776 Diamond Light Source for time on i12-JEEP under proposal EE15898; and thank all the staff at both
777 beamlines who provided additional support to the beamline activity. We are also grateful to Severn

778 Thermal Solution and everyone who advised and helped with the design and construction of the
779 technologies presented in this work.

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992 **15 Data Availability Statement**

993 The datasets for this manuscript are not publicly available at this time because they are the subject of
 994 other manuscripts in preparation. Once published they will be made available through the NERC
 995 National Geoscience Data Centre. Access will be granted where possible and requests should be
 996 directed to Katherine Dobson: katherine.dobson@strath.ac.uk

997 **16 Figure Captions**

998 Figure 1 - Schematics of the XRheo (A, D) and its installation at the TOMCAT (Swiss Light Source)
 999 (C) and at i12 (Diamond Light Source) (B) beamlines. This mounting system can be easily adapted to
 1000 other synchrotron or laboratory systems, and the system can be used *ex situ* for laboratory bench
 1001 testing.

1002 Figure 2 – 2D slices through 3D datasets showing the contrast between A) rutile in borosilicate
 1003 (filtered pink beam, SLS-TOMCAT, max energy ~40keV), B) olivine in the borosilicate (DLS-i12-
 1004 JEEP, monochromatic beam 53keV), and C) olivine in Canon N2700000 viscosity standard mineral
 1005 oil (DLS-i12-JEEP, monochromatic beam 53keV). The low contrast between olivine and borosilicate
 1006 (shown in B) makes data processing and quantitative analysis extremely challenging when working
 1007 at short detector distances (minimizing phase contrast to allow accurate rendering of thin bubble wall
 1008 films). Experiments that have used quartz, wollastonite and feldspars in synthesised haplogranitic and
 1009 dioritic melts have even lower contrasts, and the particles cannot be observed even with the lower
 1010 energy at SLS-TOMCAT. The crystals in all three samples are in the 90-180µm size fraction.

1011 Figure 3 – A) The sample cups and spindles used with the XRheo for high temperature experiments.
 1012 B-C) The Pele *in situ* furnace (see also Figure 1) as mounted and used with the XRheo. D) The Pele
 1013 *in situ* furnace with the upper insulation plates removed, this also shows the apertures set to accept
 1014 the spindles/cups (as shown) could also be used as the X-ray window, mounting the sample from the
 1015 bottom of this image for high temperature (no deformation) scanning in laboratory scanning systems.

1016 Figure 4 – Acquisition schedules typically used with the XRheo, showing how single scans of the
 1017 sample pre- and post-experiment are collected, and how the timing of the repeated multi-scan
 1018 acquisitions can fit in a deformation cycle. The full deformation cycle can either be captured as a

1019 single scan, or as a series of condition specific data sets for the acceleration (A), constant
1020 deformation (C) and deceleration stages (D). The maximum TPS used in an experiment would not
1021 usually change, but TAR can be adjusted according to the deformation, as can the total number of
1022 tomographies being collected. Some beamlines require time between scans to download the data to
1023 network storage. This can be achieved during periods of zero torque (i.e. spindle moving with
1024 sample, no deformation) or during extended periods at any of the constant deformation stages.
1025 Experiment duration is dependent only on the acquisition capability of the beamline. Heating rates of
1026 20°C/min were used to best control glass behaviour (Wadsworth et al., 2019). Cup rotation speed is
1027 fixed by image acquisition, and rotation speed of the spindle calculated accordingly with a minimum
1028 differential determined by the sensitivity of the rheometer control (0.01 rpm differential at the highest
1029 cup rotation speed used in these experiments). The strain rates within the sample are therefore
1030 controlled by the spindle rotation speed and the cup and spindle diameters.

1031 Figure 5 – 2D greyscale slices perpendicular to the rotation axis showing the same slices from the 9th
1032 and 10th tomographies at different heights within the sample. The upper right schematic shows the
1033 location of the slices within the sample volume: light blue = slice 50, dark blue = slice 100, red =
1034 slice 150 (mid-point), yellow = slice 200, green = slice 250, as shown in the image border colour. #9
1035 and #10 are at the onset of acceleration when spindle speed = 0.02 rpm clockwise. In the image data,
1036 the black regions are bubbles (or air outside the cup), the mid grey is the ceramic of the spindle and
1037 cup, and the pale grey is the borosilicate melt. A 7x7 grey grid has been overlaid on all XY images to
1038 assist in identifying the displacements and the features discussed in the text.

1039 Figure 6 – 2D greyscale slices perpendicular to the rotation axis showing the deformation and
1040 coalescence in the mid slice tomographies between #40 to #54 during the acceleration to 0.3 rpm
1041 differential rotation. Three bubbles are highlighted (red, yellow, blue) see text for details. Note that
1042 bottom row of images shows every other image.

1043 Figure 7 – 2D greyscale slices perpendicular to the rotation axis showing the deformation and
1044 coalescence in the mid slice tomographies between #60 to #94 during the end of acceleration and the
1045 start of steady state deformation at 0.3 rpm (#80 onwards). Five bubbles are highlighted (red, yellow,
1046 blue, green, pink) see text for details. Note that the bottom row shows every 5th tomography over the
1047 initial part of the constant deformation.

1048 Figure 8 – 2D greyscale slices perpendicular to the rotation axis showing the deformation and
1049 coalescence in a selection of slices through a three-phase system. The sample was sintered from the
1050 same glass bead pack with 30% rutile crystals. See text for discussion.

1051 Figure 9 – 2D greyscale slices perpendicular to the rotation axis showing the deformation and
1052 coalescence in a selection of slices through a low temperature analogue system. N2700000 + 30% by
1053 volume olivine crystals. See text for discussion.

1054 Figure 10 – The local-global integrated DVC approach showing how the outputs from one iteration
1055 and scale can be used to initiate the calculation at the next spatial or temporal step. G = Global, C =
1056 Coarse, M = Medium, F = fine. The size of the regular and irregular meshes should be adjusted to the
1057 sample structure and material texture under consideration. The meshes shown were those applied to
1058 the experiment discussed in Section 6.1 and shown in Figures 6 and 7.

1059 Figure 11 – DVC outputs from the start of the acceleration phase. Note the discontinuous time steps
1060 shown. Workflow for iterations as shown in Figure 10. The outputs shown were those generated on
1061 the experiment discussed in Section 6.1 and shown in Figures 6 and 7.

1062 Figure 12 – DVC outputs from the end of the acceleration and constant deformation phase. Note the
1063 discontinuous time steps shown. The first two frames show the last G_F and the equivalent G_M . After
1064 this the large displacements meant the G_F iteration was unnecessary. Workflow for iterations as
1065 shown in Figure 10. The outputs shown were those generated on the experiment discussed in Section
1066 6.1 and shown in Figures 6 and 7.