

# Exploring Charge Dissociation by Statistical Sample of Active Layer Models of Organic Solar Cell

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## ABSTRACT:

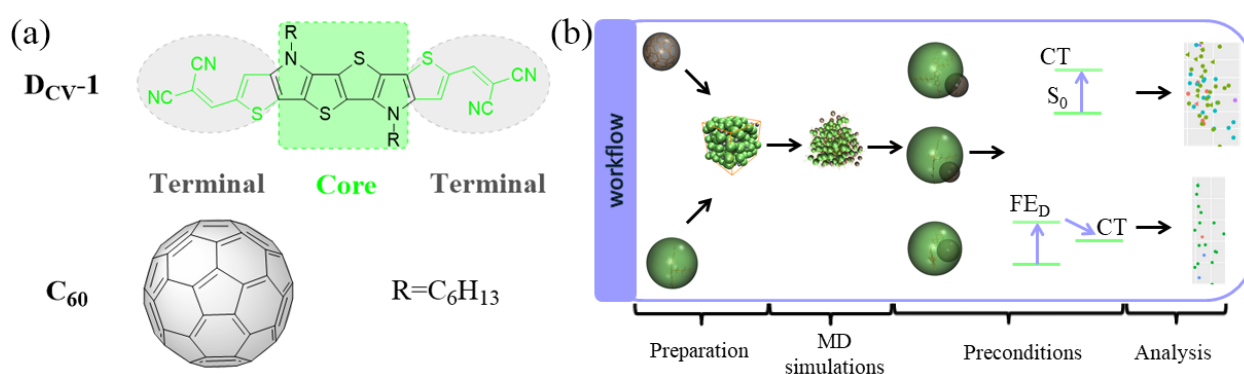
Charge dissociation in active layer is one of the key factors for the power conversion efficiency of bulk heterojunction organic solar cells (OSCs). Numerous charge-transfer mechanisms have been proposed based on one of few microscopic models. Here, we would explore possible charge-transfer mechanisms for 155 models of donor/acceptor (D/A) interfaces, built via materials **Dcv-1** and **C<sub>60</sub>** as donor and acceptor, respectively. After the calculations of the key parameters related to the charge dissociation and a statistical analysis for the correlation between these parameters were carried out, we can obtain a more robust description of the charge dissociation in practical OSCs. The complicated relationship among the key parameters not only illustrates the important correlation between D/A stacking pattern and charge-transfer mechanism, but also suggests that different charge-transfer mechanisms become more likely depending on the specific arrangements of donor and acceptor. Furthermore, the effects of excess energy on the charge-transfer mechanism were preliminarily probed by quantum dynamics simulation, which helps clarifying the much debated role of excess energy on the efficiency of charge generation.

Bulk heterojunction (BHJ) organic solar cells (OSCs) have been developed with continuous breakthrough of efficiency record in OSCs since the initial report in 1995.<sup>1</sup> In recent years, OSCs have obtained considerable attention by virtue of their distinctive superiorities.<sup>2, 3</sup> As the fundamental component of OSC device, the construction of active layer and selection of donor and acceptor materials have been considered to make central contribution to the enhancement of the power conversion efficiency (PCE),<sup>4-14</sup> as they control the microstructure of morphology and even the whole efficiency of device.<sup>15-19</sup>

In fact, based on different active layer materials and morphology of donor/acceptor (D/A) interface, the different charge-transfer mechanisms mediated by “hot” exciton,<sup>20, 21</sup> “direct” excitation,<sup>22-24</sup> built-in and external electric field,<sup>25-28</sup> entropy and disorder,<sup>29</sup> rebound negative charge,<sup>30</sup> etc. have been proposed,<sup>31-35</sup> some of which presented contradictory conclusions and fragmentary information. For instance, just for the role of excess energy, there have been some contradictory conclusions. Tamura et al. proposed a charge-transfer mechanism promoted by vibronically hot CT states in 2013, which shows the free carrier yield ascended with the increment of excess energy but a too large excess energy is unfavorable.<sup>21</sup> In 2014, Sun et al. used the Ehrenfest dynamics to emphasize that a larger excess energy is more conducive to charge separation.<sup>36</sup> Recently, there are several reports of OSCs with high internal quantum efficiencies for charge generation with negligible excess energies.<sup>37-39</sup> Many of such proposals are based on highly idealized geometries of model Hamiltonian and it is not very clear if the results depend on the specific choice of the model. So what if the contradictory conclusions are sourced from different interface models built? In recent years, many researchers have devoted themselves to the study of D/A stacking effect both in experiment and theory.<sup>40, 41</sup> The stacking patterns of donor and acceptor have been explored from the initial distance of 3.5 Å to molecular dynamics (MD) simulation in theory.<sup>42-44</sup> However, as we all know, it is difficult to measure the stacking pattern at molecular level in experiment and most of the selected theoretical D/A stacking models are over-idealized. By comparison, the D/A stacking patterns obtained by MD simulations seem closer to the actual condition,<sup>45, 46</sup> and this technique is widely adopted to characterize the local interface morphologies of D/A blends and to understand the charge-transfer process with the analysis of electronic couplings.<sup>44, 47-52</sup> Meanwhile, the effect of domain size on charge-transfer mechanism of OSCs was investigated based on a large model with PCBM supercell and large polymer domains (a model constructing by 2360 and 2048 atoms for acceptor and donor,

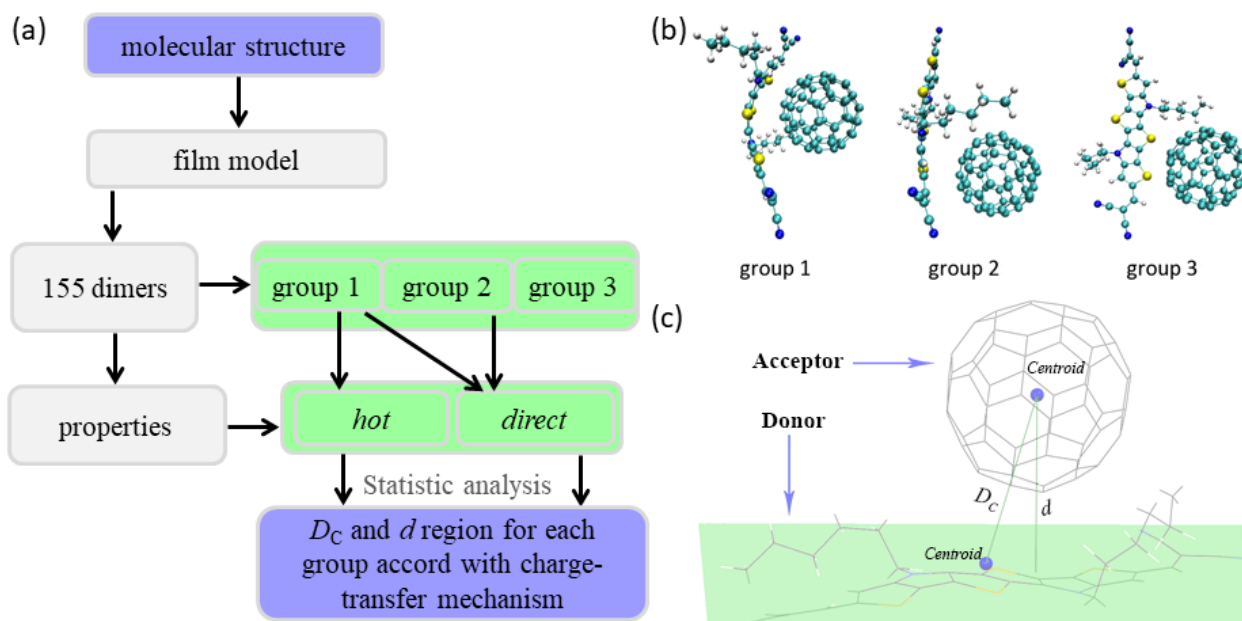
respectively) obtained by MD simulations.<sup>53</sup>

The goal of this work is to obtain charge-transfer mechanisms of different stacking patterns of a specific D/A model. For an active layer model simulated by MD, a statistical analysis was adopted to explore charge-transfer mechanism. And the bulk system is constructed with well-studied active layer materials dicyanovinylene (DCV)-substituted S,N-heterohexacenes-based<sup>54-58</sup> molecule (**Dcv-1**) and **C<sub>60</sub>** shown in Figure 1(a). Through statistical analysis on 155 D/A interfaces, we proposed a workflow visualized in Figure 1(b) to search for promising region of characteristic values for more effective D/A stacking patterns according with different charge-transfer mechanisms.



**Figure 1.** (a) Investigated donor molecule **Dcv-1** and acceptor molecule **C<sub>60</sub>**; (Molecule **Dcv-1** constructed from A-D-A type including one core unit and two terminal units.) (b) Flowchart of the multi-scale simulations.

As described in Figure 2(a), the **Dcv-1/C<sub>60</sub>** interface stacking was firstly obtained by MD simulation, which is performed in Gromacs-5.1.1 software package.<sup>59</sup> The simulation began with a three-dimensional cubic box of sides  $8 \times 8 \times 8 \text{ nm}^3$  with the numbers of D/A molecules being 144/65 according to the experimental proportion.<sup>54</sup> The methods and details of simulation process which can give an intuitive description on the formation of **Dcv-1/C<sub>60</sub>** morphologies are exhibited in the Supporting Information (SI). Then 155 **Dcv-1/C<sub>60</sub>** stacking patterns with the centroid-to-centroid distance less than  $10 \text{ \AA}$  referred to center-of mass (COM) radial distribution functions (RDFs) plotted in the SI were extracted from the final equilibration morphologies. They were further classified manually into three groups shown in Figure 2(b), namely group 1 (**C<sub>60</sub>** prefers to face on the center of backbone of **Dcv-1**), group 2 (**C<sub>60</sub>** tends to face on the terminal backbone of **Dcv-1**) and group 3 (**C<sub>60</sub>** locates at the edge of **Dcv-1**) with the assistance of configurational characteristic parameters  $D_C$  and  $d$  defined in Figure 2(c). The three groups are represented in nearly equal proportions among the 155 configurations with 50, 46, 59 instances for group 1, 2, 3 respectively.

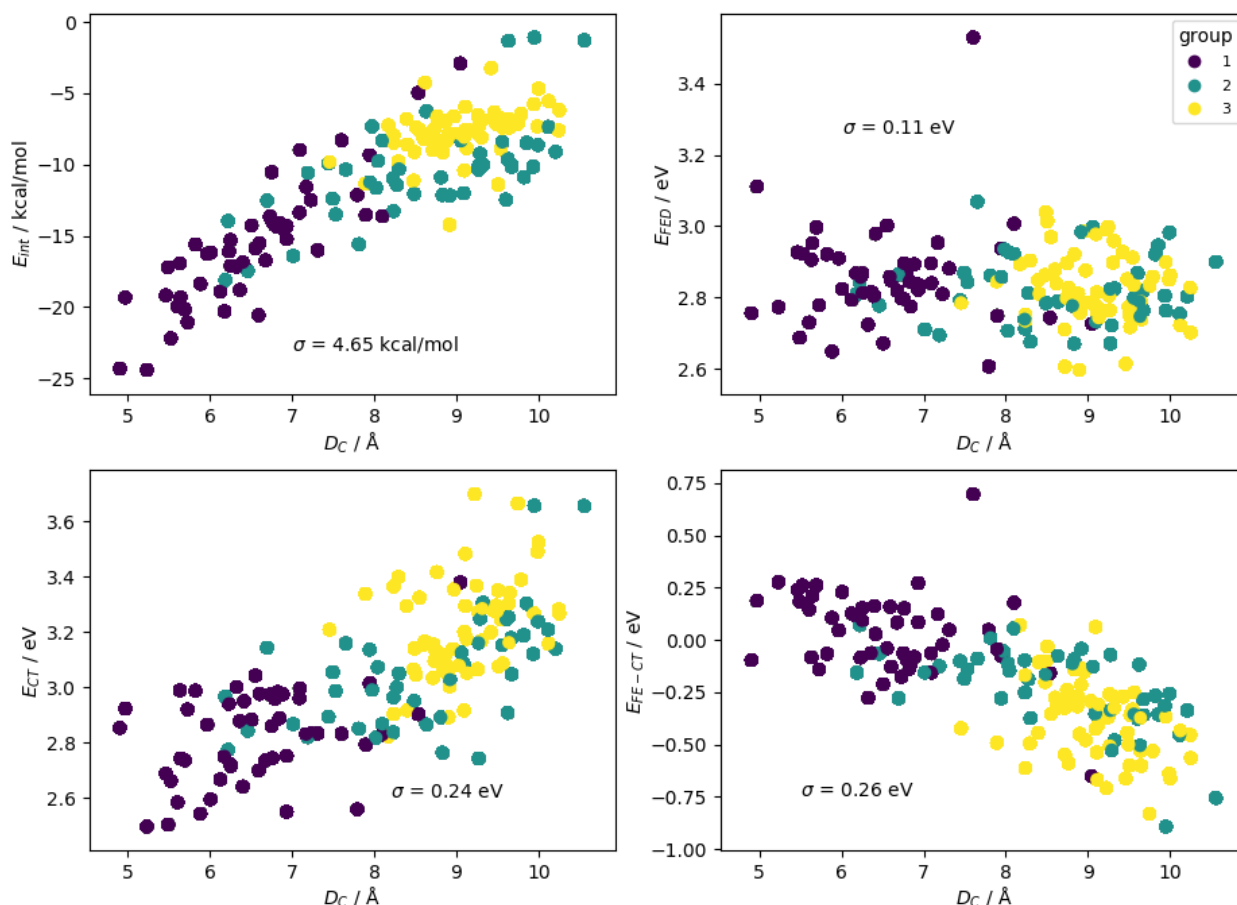


**Figure 2.** (a)Diagram of simulated process. (Input and output are shown in the purple frame. Manual categories are listed in the green frame); (b)Representative configurations extracted from the MD simulations for  $\mathbf{D}_{CV-1}$  and  $\mathbf{C}_{60}$  molecules; (c)Schematic representation of the configurational characteristics for  $\mathbf{D}_{CV-1}/\mathbf{C}_{60}$  dimer. ( $D_C$  represents the distance between centroid of donor and acceptor and  $d$  is the perpendicular distance between the centroid of acceptor and the plane of the conjugated backbone of donor. )

In order to identify the charge-transfer mechanism, we computed a number of relevant properties. Here, the counterpoised-corrected total interaction energies ( $E_{int}$ ) between  $\mathbf{D}_{CV-1}$  and  $\mathbf{C}_{60}$  were firstly calculated to characterize the interaction strength between donor and acceptor. The excited energies of Frenkel exciton (FE) state for donor ( $E_{FED}$ ) and the lowest charge-transfer (CT) state ( $E_{CT}$ ) with their respective oscillator strength  $f_{CT}$  and  $f_{FED}$  were calculated to identify their relative positions and photo-absorption strengths. Meanwhile, the energy difference between FE and CT states, namely the excess energy  $E_{FE-CT}$  was also estimated for every stacking pattern aiming at distinguishing which charge-transfer mechanism it belongs to. Moreover, the electronic couplings between  $\mathbf{D}_{CV-1}$  and  $\mathbf{C}_{60}$  ( $V_{DA-CS}$  for charge-separation process and  $V_{DA-CR}$  for charge-recombination process) were calculated considering their importance in promoting the charge-transfer ability at D/A interface using the generalized Mulliken-Hush (GMH) method<sup>60</sup> as described in the SI. All the calculations mentioned above were carried out at the  $\omega b97xd/6-31G(d,p)$  level, which is not only suitable for the system with weak intermolecular interaction but also for the obvious charge-transfer state.<sup>61, 62</sup>

As a first step to analyze the large set of results, we report some interesting properties such as  $E_{int}$ ,

$E_{\text{FED}}$ ,  $E_{\text{CT}}$  and  $E_{\text{FE-CT}}$ , in a scatter plot as a function of structural characteristic  $D_{\text{C}}$  in Figure 3 for 155 patterns and exhibit them in three groups. (scatter plot matrix for all computed parameters are shown in the SI.) Considering the distribution of excited state energy in Figure 3, one can notice that the stacking patterns have a relatively weak effect on the  $E_{\text{FED}}$  and a stronger effect on  $E_{\text{CT}}$ . As a consequence the excess energy  $E_{\text{FE-CT}}$  varies in the region of  $-0.9 \sim 0.7$  eV largely dictated by the variation of  $E_{\text{CT}}$ . It deviates from our traditional knowledge that a fixed  $E_{\text{FE-CT}}$  is definite when the pair of donor and acceptor is given whether in experiment or theory, which will be focused on in this work. Moreover, making a comparison among the three groups, group 1 has lower  $E_{\text{int}}$  while group 3 seems to have higher  $E_{\text{int}}$ , which implies more stable and favorable stacking configurations for group 1. Besides,  $E_{\text{CT}}$  values of group 1 tend to be in the lower excited-energy region and those of group 3 seem to occupy in the higher excited-energy region. And higher  $E_{\text{FE-CT}}$  of group 1 than those of other two groups is proved by the fact that most  $E_{\text{FE-CT}}$  values are positive for group 1 while only three values are positive for group 2 and only two values for group 3. Hence, we infer that most of stacking patterns in group 1 are assigned to the hot mechanism (exciton dissociation via excited (“hot”) electronic or vibrational levels, namely  $E_{\text{FE-CT}} > 0$ )<sup>20</sup> while rare case occurs for group 2 and group 3. From the analysis above, we can find that  $E_{\text{FED}}$  does not depend too much on the packing because it is an excitation localized on the donor, while  $E_{\text{CT}}$  is influenced by the packing because greater distance between donor and acceptor increase the energy for the CT process.

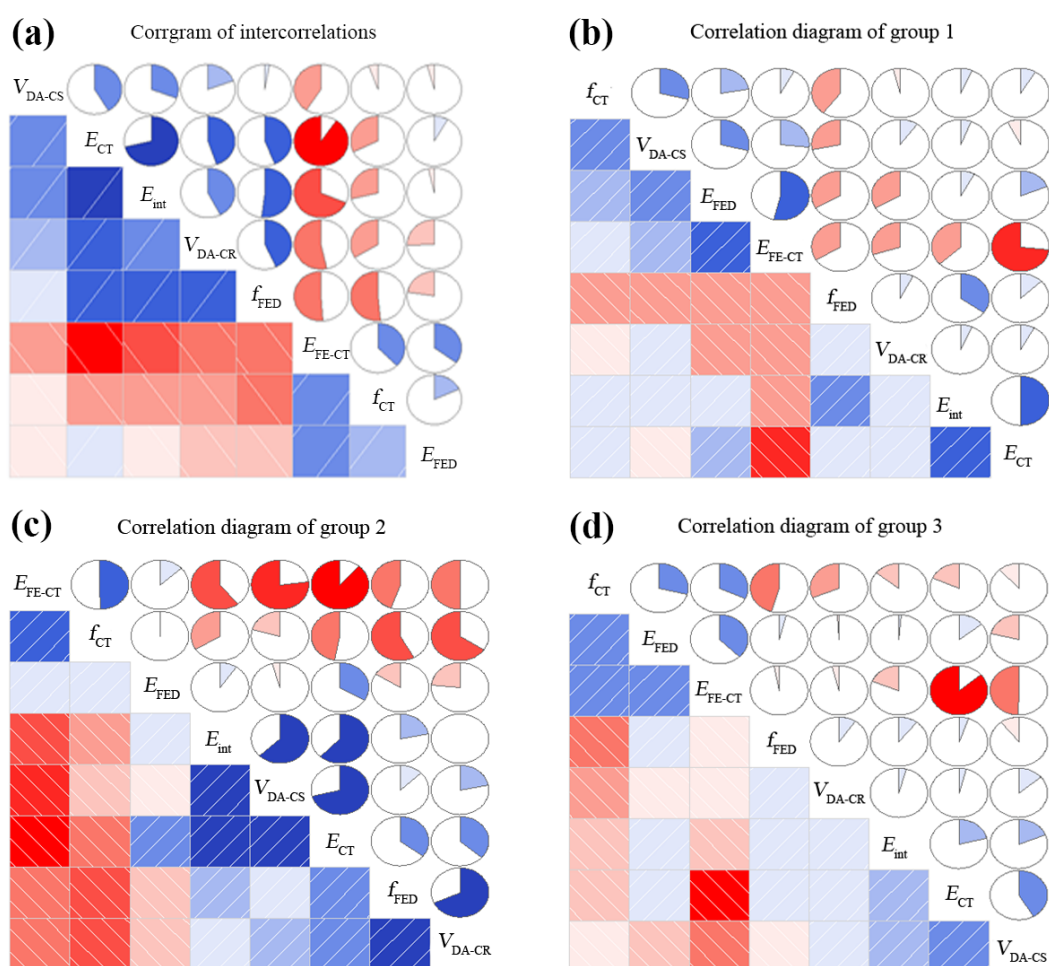


**Figure 3.** Scatter plot of some calculated properties as a function of  $D_C$  ( $E_{\text{int}}/\text{kcal}\cdot\text{mol}^{-1}$ ,  $E_{\text{FED}}/\text{eV}$ ,  $E_{\text{CT}}/\text{eV}$ ,  $E_{\text{FE-CT}}/\text{eV}$ ) of 155 dimers in different groups corresponding with average standard deviation  $\sigma$ .

The correlations for all calculated physical parameters are collected in Figure 4a, which shows the correlation degrees directly. From Figure 4a, largest correlation degree is easily observable between  $E_{\text{CT}}$  and  $E_{\text{FE-CT}}$  in light of their correlation coefficient (-0.90), which is in accordance with their linear relationship presented in scatter plot matrix in the SI. Meanwhile, the correlation coefficient between  $E_{\text{int}}$  and  $E_{\text{CT}}$  is 0.71, whose values are larger among these correlations, implying the correspondence between  $E_{\text{int}}$  and  $E_{\text{CT}}$ . Now we can infer the correlations among them, i.e., the  $E_{\text{int}}$  between donor and acceptor, which is sensitive to the stacking pattern, influences the  $E_{\text{CT}}$  and further  $E_{\text{FE-CT}}$ . In other words, the excess energy strongly depends on the molecular stacking pattern at the interface. This is mostly influenced by the intermolecular conformation of the donor, which explains the weak correlation with other parameter. By comparison, the two parameters  $E_{\text{FED}}$  and  $V_{\text{DA-CS}}$  have weaker correlation with other parameters, suggesting that they are less influenced by the interface stacking pattern and may be mainly determined by our selection of active layer materials. This is weakly correlated because it is well established that the coupling is very sensitive to very small geometry

change.<sup>63</sup>

In Figure 4b-4d, we show correlation diagrams for different groups. The relatively stronger correlation between  $E_{\text{FED}}$  and  $E_{\text{FE-CT}}$  for group 1 than group 2 and group 3 may be originated from the fluctuated  $E_{\text{FED}}$  for group 1. For group 2, the electronic coupling for charge separation shows larger correlation with  $E_{\text{CT}}$ ,  $E_{\text{FE-CT}}$  and  $E_{\text{int}}$ , while the one for charge recombination exhibits larger correlation coefficients with  $f_{\text{FED}}$  (0.68) and  $f_{\text{CT}}$  (-0.66). These phenomena are different from the cases in group 1 and group 3, and it indicates that for the stacking patterns with **C60** facing on the terminal backbone of **Dcv-1**, stronger interaction (more negative  $E_{\text{int}}$ ) is associated, as expected, with (i) lower  $E_{\text{CT}}$  (ii) higher  $f_{\text{CT}}$  and (iii) lower coupling  $V_{\text{DA-CS}}$ . This relation creates indirect correlations, e.g. between  $f_{\text{CT}}$  and  $V_{\text{DA-CS}}$ .



**Figure 4.** (a) Statistic analysis for all calculated properties ( $E_{\text{int}}/\text{kcal}\cdot\text{mol}^{-1}$ ,  $E_{\text{FED}}/\text{eV}$ ,  $f_{\text{FED}}$ ,  $E_{\text{CT}}/\text{eV}$ ,  $f_{\text{CT}}$ ,  $E_{\text{FE-CT}}/\text{eV}$ ,  $V_{\text{DA-CS}}/\text{eV}$  and  $V_{\text{DA-CR}}/\text{eV}$ ) of 155 dimers (a) and of group 1 (b), group 2 (c) and group 3 (d). (Blue and red cells indicate a positive and negative correlation, respectively, between the two variables. The darker color indicates that the variable correlation is greater. The triangular cell shows the same information with a pie chart. Correlation

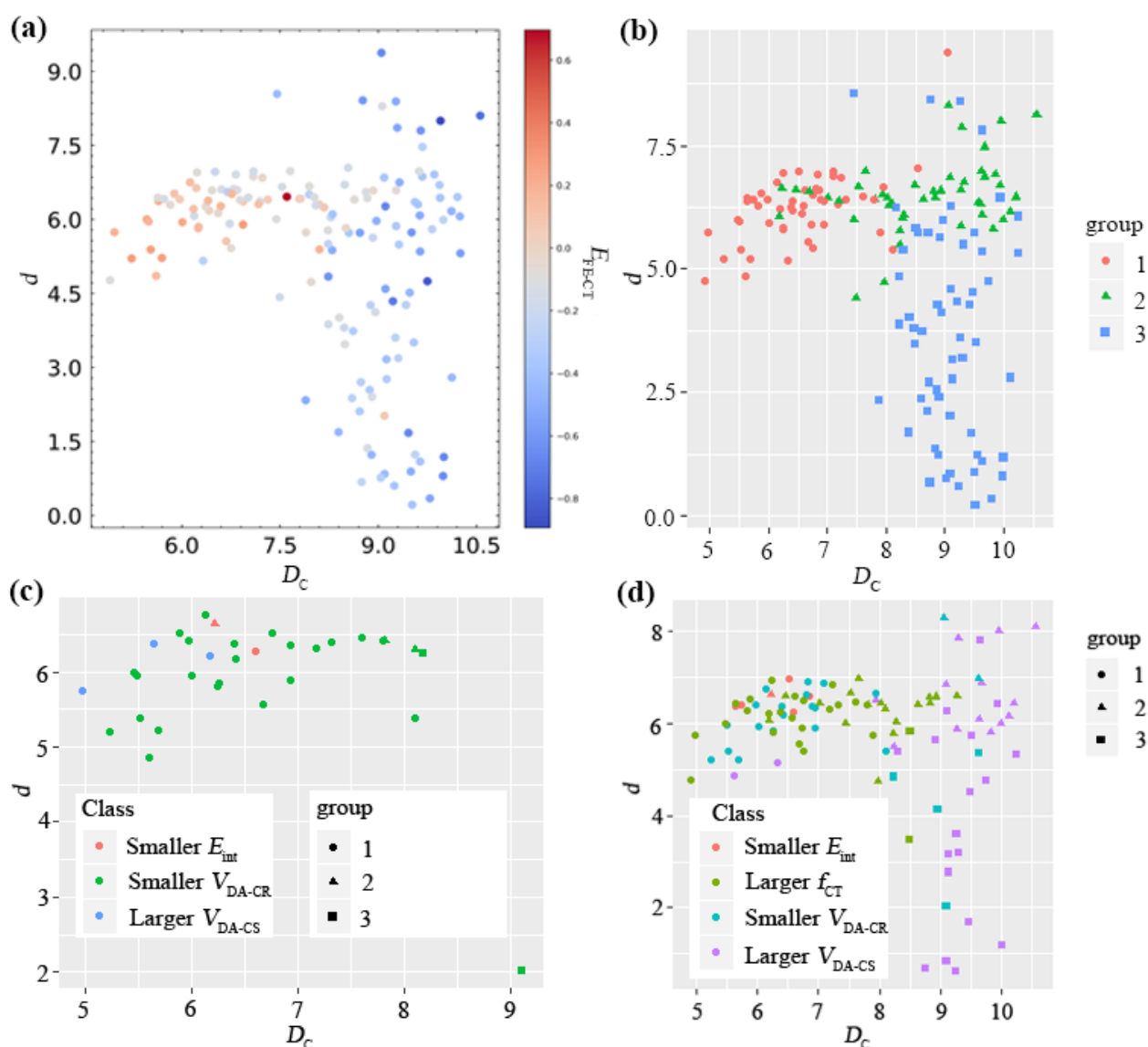
coefficients between two properties larger than 0.5 are presented in the SI.)

Figure 5a plots the variation of  $E_{\text{FE-CT}}$  along with characteristic values, trying to show the correlation between the energy and the structural characteristics. Combined with Figure 5b, we can see that different group tends to have distinguished region of  $E_{\text{FE-CT}}$  values. A preliminary test shown in the SI was conducted to explore the effect of the excess energy  $E_{\text{FE-CT}}$  on the charge-transfer mechanism by the method adopted in our previous work.<sup>30</sup> We can find that when the  $E_{\text{FE-CT}}$  is less than zero, as the  $E_{\text{FE-CT}}$  increases, the maximum time-averaged outgoing charge  $\bar{P}_{\text{out}}$  increases slowly at first and then rises substantially. When  $E_{\text{FE-CT}}$  values are approaching zero, more and more charges in bound-CT states are separated into outgoing charges. When  $E_{\text{FE-CT}}$  is greater than zero,  $\bar{P}_{\text{out}}$  decreases drastically as  $E_{\text{FE-CT}}$  increases. This is because the systems adopt hot mechanism, and the generated charges in bound-CT state are small. Considering that  $\bar{P}_{\text{out}}$  can reach its peak when  $E_{\text{FE-CT}}$  is close to zero, we can imagine that the appearance of Frenkel exciton state and CT state as a hybrid state will facilitate charge separation. It is in accordance with our another work<sup>64</sup> that the hybrid state is favorable to improve charge dissociation.

In order to provide a specific distinction between charge-transfer mechanisms for the three groups, the position distributions along with the corresponding charge-transfer mechanisms, namely hot mechanism and direct mechanism (charge dissociation can take place directly into charge), are detailed in Figure 5c-5d. It can be seen from Figure 5b that both  $D_C$  and  $d$  for group 1 are mostly located in the range of (5.0 Å, 7.5 Å), while for group 2,  $D_C$  are relatively larger with the same range of  $d$  with group 1, and for group 3,  $D_C$  are also large with  $d$  varying widely (0~9 Å). Overall, when  $D_C$  is smaller than 8 Å, most configurations belong to group 1, and more stacking patterns of group 1 are assigned to the hot mechanism with reference to Figure 5c, which only considers lower  $E_{\text{int}}$ , larger  $V_{\text{DA-CS}}$  and smaller  $V_{\text{DA-CR}}$  (The relative lower or larger value are selected on the basis of the distribution of data. According to the distance between two adjacent points, the data is arranged in order from small to large. If the distance between two adjacent points is relatively large, the middle of these two points is the split point.) for hot mechanism. We know that effective charge dissociation needs larger  $V_{\text{DA-CS}}$  and smaller  $V_{\text{DA-CR}}$ . The corresponding promising region ( $D_C = (5.5, 7.8)$  and  $d = (6.2, 6.6)$ ) for hot mechanism in group 1 are presented in the SI. While for group 2 and group 3, hot mechanism is not considered because the majority of negative  $E_{\text{FE-CT}}$  values are in both groups. In comparison, the stacking patterns



assigned to direct mechanism show a larger proportion according to the dense distribution in Figure 5d, whereas, the position distributions with effective charge dissociations are still concentrated at smaller intermolecular distance. The promising regions for direct mechanism of group 1 ( $D_C = (5.5, 7.4)$  and  $d = (6.0, 6.7)$ ) and group 2 ( $D_C = (7.0, 9.0)$  and  $d = (6.0, 7.0)$ ) are shown in the SI. Therefore, we get the conclusion that the hot mechanism mainly appears in group 1, while direct mechanism mainly appears in group 1 and group 2, and more stacking patterns in group 3 behave ineffective charge dissociations considering their smaller  $V_{DA-CS}$  and nearly zero  $f_{CT}$ .



**Figure 5.** (a) Distribution of  $E_{FE-CT}$  along with  $D_C$  and  $d$  for the dimers; (b) Distributions of  $D_C$  and  $d$  for all of 155 dimers; (c) Distributions of  $D_C$  and  $d$  for the dimers with hot mechanism (only consider  $E_{FE-CT} > 0$ ). Region  $D_C = (5.5, 8.0)$  (Å) and  $d = (5.5, 6.5)$  (Å) is promising; (d) Distributions of  $D_C$  and  $d$  for the dimers with direct mechanism. Region  $D_C = (5.0, 8.0)$  (Å) and  $d = (5.0, 7.0)$  (Å) is promising.

In conclusion, we performed a statistical analysis on the key parameters related to the charge dissociation in an active layer model trying to search possible charge-transfer mechanisms. We find that different charge dissociation mechanisms become more likely depending on the packing of the D/A pair. A hot exciton mechanism is the most likely to occur when the acceptor and donor are the closest with acceptor lying on the central of donor (in about one third of configurations). While the packings for acceptor faced on the donor molecule could have a strong propensity to have direct excitation of the charge-transfer state. As for those acceptor edged on the donor are hardly satisfied with either hot or direct mechanism. In addition, the quantum dynamics simulation suggests that when the excess energies come close to zero, the separation charge yield is approaching the peak. This work demonstrates the importance of considering large sample of intermolecular geometries to draw useful conclusions on the actual charge dissociation mechanism.

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