Laser Derived Electron Transport Layers with Embedded p-n Heterointerfaces Enabling Planar Perovskite Solar Cells with Efficiency over 25%

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Abstract

Electron transport layers (ETLs) with pronounced electrons conducting capability are of essential for high performance planar photovoltaics, while this has been posing great challenges for the most widely used metal oxide ETLs that are featured unfortunately with intrinsic low carriers mobility. We herein demonstrate that by simply addressing the carriers loss at particle boundaries of TiO₂ ETLs through embedding in ETL p-n heterointerfaces, the electrons mobility of the ETLs can be boosted as high as three orders of magnitude. Such embedding was also found encouragingly favorable for inhibiting the formation of rutile phase TiO₂ in ETL, as well as initiating the growth of the on top high-quality and large-grain perovskite films with less defect states. By virtue of these merits, we achieve in present study formamidinium lead iodide PSCs with a champion efficiency of 25.05%, setting a new benchmark for plannar PSCs employing TiO₂ ETLs. Unencapsulated PSCs employing such ETLs deliver also much-improved environmental stability, i.e., more than 80% of their initial efficiency after 9000 h of air storage under RH of 40%, and over 90% of their initial efficiency at maximum power point under continuous illumination for 500 h. Further work exploring other laser generated p-type nanocrystals for embedding warrants the proposed strategy a universal alternative for addressing the low carriers mobility of metal oxide based ETLs from the view point of particle boundaries engineering.

Keywords: electrons transport layer, p-n heterointerfaces, electron mobility, particle boundaries, perovskite solar cells

Introduction

Planar halide perovskite solar cells (PSCs) have been attracting increased interests owing to their potentials on low-cost construction of high performance tandem and/or flexible photovoltaics.^[1-2] In addition to the important photoactive perovskite layer that absorbs incident photons to generate sufficient photogenerated carriers, the planar PSCs necessitate the electron transport layer (ETL) that plays crucial roles not only on efficiently extracting and conducting the photogenerated electrons to the out circuit,^[3] but also elaborately ruling the growth of the on-top high quality halide perovskite.^[4] Compact films of metal oxides such as TiO₂, ZnO, SnO₂, and La-BaSnO₃ have been widely explored as efficient ETLs due to their ease solution-processing that is compatible with the low-cost and low-temperature processing of the perovskite layer.^[5] As one of the most important ETLs that is traditionally used in dye-sensitized solar cells,^[6] TiO₂ is screened as the most widely adopted ETL in regular planar PSCs owing to its high optical transmittance and excellent thermal stability, low cost as well as suitable energy level.^[7] With the intensive efforts devoted on the engineering of TiO₂ ETL in planar PSCs, the power conversion efficiency (PCE) could reach a benchmark value as high as 24.8%, which unfortunately still lags behind the efficiency of PSCs employing SnO₂ as ETL.^[8] The reason behind such efficiency gap is largely due to the insufficient electrons conducting capability of the TiO₂ ETL, which principally accounts for the electron accumulation at the buried interface and the unbalanced carrier transport within PSCs, thus leading to the unsatisfied I-V hysteresis, less-than-ideal PCE and lightinduced instability.^[9-10] Encouragingly, the intrinsic low carriers mobility of the TiO₂ nanoparticles (usually at the level of $10^{-5} \sim 10^{-4} \text{ cm}^2 \text{V}^{-1} \text{s}^{-1}$) in the ETLs has recently been intensively addressed by several state-of-art engineering strategies, exemplified by a high-quality single-crystalline TiO₂ nanoparticles based ETL with the electron mobility as high as 6.04 cm²V⁻¹s⁻¹, as well as the nitrogen-doped TiO₂ ETL with an increased electrons mobility up to $0.5 \text{ cm}^2 \text{V}^{-1} \text{s}^{-1}$ that leads to the PSC with a fill factor >86%.^[11-12] While noting that TiO₂ ETL is essentially a layer composed of stacked TiO₂ nanoparticles, the inhibition of carriers loss at the particle boundaries (PBs) in between adjacent nanoparticles of ETL should be considered as equally important as improving the intrinsic carriers mobility of the nanoparticles, as described in Schemes 1a and 1b.^[13] Inspired by the grain boundaries engineering of the polycrystalline perovskite films that could alleviate the defect trap states and modulate the carriers dynamics, it would thus be highly promising to improve the electron mobility of TiO₂ based ETLs by addressing the carriers loss at PBs for highly efficient and stable planar PSCs.

Herein, we demonstrate an effective strategy of PBs engineering via laser embedding p-n heterointerfaces at ETLs, which could enable boosted electrons mobility of TiO_2 ETL by maximally three

orders of magnitude, thus results in a planar PSC with a new benchmark PCE over 25%. By virtue of the unique advantage of pulsed laser irradiation in liquid on producing well defined nanocrystals in desired solvents, we are able to embed laser generated p-type nanocrystals at the PBs of the n-type TiO_2 ETLs. The embedded p-n heterointerfaces could efficiently inhibit the carrier loss at PBs owing to not only the elimination of the electrons trapping at TiO_2 nanoparticles but also boosted electrons transfer between adjacent TiO_2 nanoparticles through localized built-in electric filed (Scheme 1c). By elaborately modulating the amount of embedded p-type CdTe nanocrystals, the electrons mobility of ETLs could be enhanced from 2.67×10^{-5} up to 1.89×10⁻² cm²V⁻¹s⁻¹. The embedding of CdTe nanocrystals were found also helpful to modulate the crystallization kinetics of TiO_2 matrix, which was verified favorable for the inhibition of rutile TiO_2 that is detrimental for the electrons conducting and light induced stability of perovskite. Owing to the boosted electrons conducting at ETL and subsequent suppressing of charge accumulation and recombination at the interface between ETL and perovskite, we achieve in present study highly efficient formamidinium lead iodide PSCs with a top record efficiency of 25.05%, which is as far as we know the highest value among TiO₂-based planar PSCs. The novel TiO₂@CdTe ETL was found helpful to initiate high-quality and large-grain perovskite films with reduced interfacial defects, leading to the construction of the mixed-cation PSCs with stable shelf storage over 9000 h (moisture RH of 40%), and FAPbI₃ PSCs that could maintain over 90% of their initial efficiency at maximum power point under continuous illumination for 500 h. We believe this study provides an efficient alternative of improving the carriers mobility of charge transport layer for high performance planar PSCs from the view point of PBs engineering.

Results

Embedding laser-generated p-type CdTe nanocrystals in TiO₂ ETL.

In order to locate the p-type nanocrystals at the PBs of TiO_2 ETL, we first use the technology of pulsed laser irradiation in liquid to generate size controlled CdTe nanocrystals (Figure 1a) in the precursor solvent used for chemical bath deposition of TiO_2 ETL. By optimizing the processing parameters such as laser fluence, irradiation time and cooling temperature (Figure S1, Supporting Information), a stable and well-dispersed ptype CdTe nano-colloids showing Mie-scattering and with an average size of 3 nm were consequently obtained, as illustrated in Figure 1b. High resolution transmission electron microscopy (HRTEM) was employed to determine the crystal structure of the as-prepared CdTe nanocrystals. The lattice fringes with a spacing of 0.37 nm shown in Figure 1c are affiliated to the characteristic plane (111) of CdTe, which is further evidenced by the corresponding fast Fourier transform (FFT) of selected area that shows (111) and (311) plane with 29.5° included angle. Similar results towards larger nanoparticles were also acquired and further confirmed by Raman spectroscopy (Figures S2-3, Supporting Information), indicating that there is no significant crystalline change in CdTe during laser irradiation. Furthermore, the elements mapping determined by TEM-energy dispersive spectroscopy (TEM-EDS) suggests homogeneous distributions of all elements throughout both nanoparticles and nanocrystals without elements segregation during laser process (Figure S4, Supporting Information), consistent with the analyses of the X-ray photoelectron spectroscopy (XPS, Figure S5, Supporting Information).

Subsequent to the acquiring of CdTe nanocrystals in the precursor solvent for depositing TiO₂, a facile onestep chemical bath co-deposition approach is adopted for *in situ* embedding laser generated sub-10 nm CdTe nanocrystals (Figure 1a and Figure S6, Supporting Information). In brief, different amounts (3%, 6% and 9% volume ratio to TiCl₄ precursor abbreviated as 3%-CdTe-TiO₂, 6%-CdTe-TiO₂, 9%-CdTe-TiO₂) of CdTe colloids solution with 0.1 mg/ml concentration were added to the TiCl₄ precursor for construction of ETL. The surface morphologies of the TiO₂ ETL with or without embedding of CdTe nanocrystals were evaluated by scanning electron microscopy (SEM) and atomic force microscopy (AFM). Compared with the pristine TiO₂, CdTe@TiO₂ exhibits more smoother surface with reduced roughness from 31.0 nm to 12.1 nm (Figures 1f-g, Figures S7-8, Supporting Information). As it was reported that the embedding of laser generated nanocrystals at the grain boundaries of perovskite could in principle alter the crystalline kinetics of the matrix, we collected the predicated TiO₂ powder after the co-deposition for further XRD analysis to investigate if the embedding of CdTe nanocrystals could influence the crystallization kinetics of TiO₂. As shown in Figure 1d, the obtained TiO₂ without embedding of CdTe exhibits a polymorphic phase including rutile and anatase in the range of temperature from 30 to 150 °C. By contrast, the rutile phase is remarkably restrained after embedding CdTe nanocrystals, as shown in Figure 1e.

Generally, anatase phase is easily formed with fast nucleation rate during the hydrolysis of TiCl₄, whereas slow nucleation rate facilitates directional arrangement of aggregates and formation of more stable rutile phase.^[14] We thus deduce that the incorporation of CdTe nanocrystals could initiate the crystal nucleation and speed up the crystallization of TiO₂, since more nanocrystals in the starting precursor result in the formation of more nuclei at the beginning of hydrolysis process and thus smaller TiO₂ nanocrystals, which favors dispersion of TiO₂ nanocrystals in solvent and formation of compact anatase TiO₂ film.^[9] This result is in agreement with the case of using nanocrystals in anti-solvent to modulate the crystalline kinetics of perovskite films. We believe such rapid crystallization induced inhibition of the rutile phase would be much favorable for

the construction of an ideal TiO_2 ETL in PSCs, not only because rutile phase TiO_2 has lower carrier mobility than that of anatase phase, but also due to the elimination of the mixed anatase-rutile TiO_2 with unfavorable photocatalysis property that are detrimental for the stability of PSCs upon AM 1.5 illumination.^[15]

In order to confirm the existence of CdTe nanocrystals within the bulk TiO₂, cross-sectional TEM with EDS mapping is performed to probe spatial distribution of CdTe nanocrystals throughout the entire thickness. As shown in Figure 2a, the EDS mapping shows homogeneous distribution of CdTe nanocrystals within overall matrix due to well-dispersed CdTe nanocrystals in the TiCl₄ precursor (Figure S9, Supporting Information). The Cd and Te compositions do not change after embedding CdTe nanocrystals in TiO₂ matrix and retain atomic fraction ratio of 1:1 (Figure S10, Supporting Information). The cross-sectional HRTEM (Figure 2c from magnified red region of Figure 2b) convincingly demonstrates successful embedding by identifying typical lattice space of 0.35 nm and 0.37 nm that successively corresponds to (101) plane of TiO₂ and (111) plane of CdTe. Such embedding is further determined by corresponding fast Fourier transform (FFT) of selected area in Figure 2c, demonstrating (101) and (200) planes of TiO₂ with 21.7° plane included angle (Figure 2d), as well as (111) and (311) planes of CdTe with 29.5° angle (Figure 2e).

The UV-vis absorption spectra and corresponding Tauc plot were conducted to determine optical bandgap of ETLs upon CdTe embedding. It is found that embedding of CdTe nanocrystals causes slight change on the bandgap of TiO₂ ETL ranging from 3.20 to 3.23 eV, but negligible influence on the optical transmittance (Figures S11-12, Supporting Information). The space charge-limited current (SCLC) was adopted to evaluate the defect density (N_t) and charge mobility of different ETLs. The results exhibit an obvious reduction of N_t from pristine 3.95×10¹⁶ to 1.32×10¹⁶, 0.97×10¹⁶ and 1.59×10¹⁶ cm⁻³ for 3%, 6%, 9%-CdTe-TiO₂ ETLs, respectively (Figure S13 and Table S1, Supporting Information), which could be attributed to the enhanced crystallinity from the embedded CdTe nanocrystals. Also, the electron mobilities (µ) were calculated to be 2.67×10⁻⁵, 8.10×10⁻³, 1.89×10⁻² and 6.07×10⁻⁴ cm² V⁻¹ s⁻¹ for the pristine and 3%, 6%, 9%-CdTe-TiO₂ ETLs, respectively (Figure 2f, Note S1, Supporting Information), indicating a maximum increment of three orders of magnitude after the embedding of CdTe nanocrystals. Simultaneously, the conductivity (σ) of target ETL increases dramatically compared to that of pristine ETL, agreeing well with the mobility result (Figure S14, Supporting Information). The increased mobility and conductivity are probably ascribed to the formation of pn heterojunction between TiO₂ and CdTe, in specific, p-type semiconductor feature of CdTe is confirmed by Ultraviolet photoelectron spectroscopy (UPS, Figure 2g) based on the optical bandgap (2.26 eV) (inset in Figure 2g), Fermi energy level (-3.92 eV), the calculated valence band value of -4.92 eV, and conduction band

value of -2.66 eV, respectively, reflected by an inverted "V-shape" Mott-Schottky plots shown in Figure 2h.^[16]

Influence of embedding of CdTe in TiO₂ ETL on the top perovskite film.

We find that the embedding CdTe in TiO₂ ETL also plays a significant role on the surface morphology and grain sizes of the top perovskite films. Considering the improved surface smoothness of TiO₂ ETL upon CdTe embedding (Figures S8, Supporting Information), the contact angle measurement was further conducted to investigate the wettability of perovskite precursor on different TiO₂ substrates (Figure S15, Supporting Information). The result shows that the contact angles sharply decrease from 15° to 2° as the content of CdTe nanocrystals increases, favoring the reduction of Gibbs free energy for improved crystallization kinetics during the perovskite deposition.^[10] As depicted in Figure 3a-b, the mixed-cation perovskite films with the composition of Cs_{0.05}(FA_{0.85}MA_{0.15})_{0.95}PbI_{2.55}Br_{0.45} (CsFAMA) deposited on CdTe@TiO₂ ETL show smoother surface and larger grain size compared to those deposited on pristine TiO₂ ETL, further evidenced by the decreased RMS roughness from AFM characterization and improved crystallinity from XRD analyses (Figures S16-17, Supporting Information).

The steady-state photoluminescence (PL) and time-resolved PL (TRPL) spectra were employed to investigate the carrier dynamics between perovskite and ETL. As shown in Figure 3c, the PL intensity of perovskite films deposited on optimal CdTe@TiO₂ ETLs significantly decreases by about four times, manifesting more effective electron transport and extraction from perovskite to ETL. Meanwhile, the carrier lifetime of perovskite films upon all ETLs extracted from time-resolved photoluminescence (TRPL) spectra shows prominent decrease from 247.04 ns to 92.55 ns and 82.14 ns for the pristine TiO₂, 3%-CdTe-TiO₂ and 6%-CdTe-TiO₂, respectively (Figure 3d and Table S2, Supporting Information). The result reveals suppressive charge accumulation and enhanced electron extraction and transport at the buried interface.^[17] We proceeded to conduct electrical impedance spectroscopy (EIS) to probe interface carrier transport behavior within the device. The result shows the decrease of the contact resistance *R*co values from 27159 Ω to 2422 Ω , and the increment of the recombination resistance *R*rec values from 1.47 × 10⁵ to 3.11 × 10⁵ after embedding nanocrystals, indicating improved charge transport and effectively suppressed charge recombination at the buried interface (Figures 3i and S18, Table S3, Supporting Information).^[18]

Subsequently, it was found that the embedding of CdTe also favors for the effectively restrained leakage current from carrier recombination by dark current evaluation, as well as higher built-in potential (V_{bi}) enabling improved open-circuit voltage (Voc) output by capacitance-voltage (C-V) measurement of targe devices (Supplementary Figure S19, Figure S22c, Supporting Information). This is attributed to more efficient charge

extraction for restrained recombination caused by serious electron accumulation, and better energy level alignment between CdTe embedded TiO_2 and perovskite due to the upward shifts of TiO_2 work function by the embedded p-n heterointerfaces for reduced the interface barrier, as proved by UPS result (Figure 3j, Note S2 and Figures S20-21, Supporting Information).^[7]

We further found that embedding of CdTe nanocrystals has important effects on tuning the defects in the perovskite absorber, which was confirmed by the temperature dependent admittance spectroscopy (TAS) combined with Mott-Schottky analysis (Note S3, Supporting Information). The capacitance-frequency (C-f) curves of different devices were measured at various temperature (210-330K) under dark to calculate their defect activation energies (Ea) to be 0.252 and 0.219 eV, respectively (Figure 3e, Figure S22a-b, Supporting Information). The trap density-of-states distribution and their energy levels are presented in Figure 3f, indicating that such embedding effectively decreases the energy level of trap states from 0.15 to 0.12 eV, as well as their density of states from 2.22×10^{16} to 1.24×10^{16} cm⁻³ eV⁻¹. In addition to the improved crystallinity of perovskite film on CdTe@TiO₂ ETL, we also ascribe such considerable decrease of N_t for perovskite films to suppression of interfacial iodine vacancy arising from the interaction between CdTe and perovskite, i.e. the binding of Cd²⁺ and iodine ions as well as under-coordinated Pb²⁺ defects by Cd²⁺(O).^[19] This was experimentally verified by the XPS analysis of the product after mixing laser generated CdTe and PbI₂, as shown in Figure 3g, where the Cd 3d_{3/2} was assigned to four peaks at 410.7, 412.1, 413.6 and 415.5 eV corresponding to Cd²⁺(O), Cd-Te, Cd-I and Cd²⁺ (complex), respectively.^[19] The peaks of I 3d_{3/2} can be fitted to two peaks at 630.7 and 632.0 eV belong to the bands of I-Pb and I-Cd that results from the interaction between Cd²⁺ and iodine for preventing iodine ions from migrating, reflected from disappearance of metallic Pb (Figure 3h and Figures S23-24, Supporting Information).^[20]

Photovoltaic performance of planar PSCs employing CdTe@TiO2 ETLs.

The regular planar PSCs with the configuration of FTO substrate/CdTe@TiO₂/perovskite/Spiro-OMeTAD/Au were fabricated to confirm the effect of CdTe nanocrystals embedding on photovoltaic performance (Figure S25, Supporting Information). The detailed fabrication process of device is shown in Figure S26, Supporting Information. Figure 4a shows the current density-voltage (*J-V*) curves of different CsFAMA-based champion devices measured under illumination of 100 mW/cm² (AM 1.5G) and the photovoltaic parameters were summarized in Table 1. The device upon 6%-CdTe-TiO₂ layer achieved the champion PCE up to 22.00% with Voc of 1.191 V, short-circuit current density (*J*_{SC}) of 23.76 mA/cm² and fill factor (FF) of 77.74%, which far exceeds those of pristine TiO₂ (20.05%), 3%-CdTe-TiO₂ (21.68%) and 9%-

CdTe-TiO₂ (20.98%) based devices. Moreover, champion device shows negligible hysteresis (hysteresis factor: 0.9%) and a stabilized PCE of 21.94% close to champion efficiency, in contrast with that of control (12.8%) (Figure 4b and Table S4, Supporting Information). To further confirm the roles of laser embedding of CdTe nanocrystals for their enhanced photovoltaic performance, we constructed pure a-phase FAPbI3 absorber based PSCs. It was found that target FAPbI₃ PSCs exhibit optimized morphology with larger grain throughout the entire thickness of film, perfect crystallinity orientation along the direction perpendicular to the substrate as well as improved interfacial carrier dynamics (Figures S27-31, and Table S5, Supporting Information). Encouragingly, we acquire a champion device delivering the efficiency of 25.05% (stabilized power output: 24.91%) with Voc of 1.167 V, J_{SC} of 25.61 mA/cm² and fill factor (FF) of 83.81%, a maximum 10% PCE enhancement relative to control (22.72%, a band gap of 1.53 eV), as shown in Figure 4c and Figure S32, Supporting Information. As shown in Figure 4d, the enhancement of EQE response in the entire range is attributed to the compact and large grains of perovskite film grown on CdTe nanocrystals embedded ETLs,^[21] resulting in improved Jsc from 24.06 and 25.11 mA/cm² for control and target devices, respectively, which agree well with the Jsc values from J-V results. The highly stabilized PCE demonstrated that devices embedded by CdTe nanocrystals enable suppression of ion migration, as illustrated by better alignment of the logarithmic I-V plots under contrast sweep direction (Figure 4e).^[22] Meanwhile, a histogram of 50 individual devices based different perovskite suggests that average PCE are improved from 18.95±0.73% to 21.23±0.49% and from 21.36±1.29% to 23.81±0.85%, respectively, as well as the enhancement of Voc, Jsc and FF (Figure 4f, Table 1 and Figures S33-34, Supporting Information).

Owing to the improved crystallinity and decreased defect states of perovskite films on CdTe@TiO₂ ETLs, we further evaluate the effect of CdTe emdedding in ETL on environmental stability of PSCs. As shown in Figure 4g, the long-term humidity durability of all unencapsulated CsFAMA-based devices was examined under 40% relative humidity (RH) in the dark at room temperature. The result shows that the target devices maintain almost 81% of initial PCE over 9000 h, whereas control devices degrade continuously to 42% after 1000 h. Likewise, the target devices exhibit excellent thermal and operational stability (Figure S35, Supporting Information). In view of the fact that FAPbI₃ perovskite has typical spontaneous phase-transition from α - to δ -FAPbI₃ under ambient atmosphere, we further examined the stability of FAPbI₃ based PSCs employing CdTe@TiO₂ ETLs. As shown in Figure S36, the humidity stability of unencapsulated target FAPbI₃ devices is significantly enhanced, retaining over 85% of initial efficiency for 1200 h under the RH of 30~40%, compared with that of control devices (less than 32%). The operational stability of FAPbI₃ devices was conducted by

maximum power point (MPP) tracking at 60 °C under 1 sun illumination in inner atmosphere (Figure 4h). After illumination of 200 h, the pristine FAPbI₃ device degrades to 68% of its initial PCE, whereas the nanocrystals embedded FAPbI₃ device maintains over 90% of its initial PCE for 500h. Such improved light stability could also be due to the prohibited formation of mixed anatase-rutile TiO₂ phases in ETLs owing to the inhibited rutile TiO₂ from CdTe induced fast crystallization of ETL (Figure 1e).

Discussion

It is worth noting that the embedding of CdTe nanocrystals has actually multifunctions for the improving of PSCs performance: i) embedded p-n heterointerfaces enable a remarkable increment of three orders of magnitude in the electron mobility of ETL (Table S6, Supporting Information), due to the synchronously eliminating of the carriers loss at the PBs of TiO₂ ETL and the electrons trapping at TiO₂ nanoparticles surface through localized built-in electric filed; ii) the introducing of CdTe nanocrystlas in the precursor solution of TiO_2 ETL could initiate the fast crystallization of TiO_2 film and thus inhibited rutile TiO_2 phase with slow nucleation rate, which accounts partially for the improved electrons mobility of ETL as well as light stability of PSCs employing CdTe@TiO2 ETLs; iii) the increased smoothness and the improved wettability of CdTe@TiO₂ ETLs, as well as the synchronous passivation of interfacial under-coordinated Pb²⁺ by Cd²⁺(O) and the reinforced anchoring effect of Cd²⁺ and iodine ions, favor for the construction of the top high-quality and large-grain perovskite with less defect states.^[10,19] Benefiting from these desirable merits, a FAPbI₃ based PSC with the top record of PCE up to 25.05%, as well as a mixed-cation PSC with pronounced moisture stability (RH of 40%) over 9000 h were achieved. It could be deduced from significant improvement of all photovoltaic parameters that the increment in Voc, Jsc and FF are due to better energy level alignment between perovskite and ETL and suppressed interfacial nonradiative recombination,^[23] suitable energy gaps and efficient charge extraction,^[24] as well as reduced contact resistance at the perovskite/ETL interfaces,^[25-26] respectively. In particular, a negligible hysteresis shown in 6%-CdTe-TiO₂ based device is due to the removal of carrier transportation imbalance at both electrodes of the devices upon enhancement of electrical property of ETL. Furthermore, 0.1 mg/ml as the optimal concentration of laser embedding is proved from J-V results employing different concentration of CdTe nanocrystals (Figures S37, Table S7, Supporting Information)

In virtue of universal and unique features of pulsed laser irradiation on generating any nanocrystals in desired solvents, we further explore the embedding other laser generated p-type semiconductor nanocrystals into TiO₂ film for enhancing photovoltaic property of PSCs. Exemplified by the p-type GaAs, we prepared a transparent gallium arsenide (GaAs) colloidal solution with well-dispersed sub-10 nm nanocrystals (average

size: 3 nm) by laser irradiation. The Mie-scattering and the HRTEM image exhibits the labeled lattice fringes (0.33 nm) that correspond to the lattice spacing of GaAs (111) plane (Figure S38a, Supporting Information). The element mapping presents the stable existence for all elements whether nanoparticles or nanocrystals after laser irradiation, whereas nanocrystals present heterogeneous element distributions due to Ga oxide shell preferentially formed based on the lower oxidation Gibbs free energy of Ga than that of As (Figures S38b, Supporting Information).^[27] The existence of GaAs nanocrystals in the TiO₂ film has been verified by SEM-EDS as well as corresponding elements mapping, (Figures S39a, Supporting Information). Likewise, the embedding of such nanocrystals contributes to more denser and smoother surface of TiO₂ and perovskite films than control, and observably improved μ from 2.67 $\times 10^{-5}$ to 1.08×10^{-2} cm² V⁻¹ s⁻¹, as well as construction of p-n junction enabling the upward shift of TiO₂ work function with better energy level alignment with perovskite (Figures S39b-40, Supporting Information). Consequently, the initial attempts on construction of PSCs employing GaAs@TiO₂ ETLs yield CsFAMA-based device with improved PCE up to 21.54% with enhanced photovoltaic parameters, as well as FAPbI3-based PSC with improved efficiency up to 23.49%, (Figures S41-42, Table S8, Supporting Information). Though the efficiency needs a substantial improvement comparing to the embedding of CdTe, we believe such attempt proves the universal feature of the strategy of embedding p-type nanocrystals at PBs of TiO₂ ETLs for addressing the unwanted carriers loss in metal oxides based ETLs, which would thus be inspiring for the design of efficient ETLs for high performance of PSCs with the planar configuration.

Conclusion

In summary, we have demonstrated an efficient strategy of boosting the electrons conducting capability of metal oxide based ETLs, which is of particularly importance for developing high performance planar PSCs. By generating well defined p-type CdTe nanocrystals in the precursor solvent of TiO₂ ETL, the nano-colloids were found helpful for regulating the crystallization kinetics of TiO₂ by inhibiting the rutile TiO₂ that is detrimental for the electrons conducting and light stability of the on top perovskite. The embedded p-n heterointerfaces are considered favorable not only for eliminating the electron trapping at TiO₂ nanoparticles surface and carrier loss at PBs, but also the upward shifts of TiO₂ ETLs, also exert significant influence on the growth of high-quality perovskite films with larger grain size, lower defect density and improved crystallinity. We consequently achieved a top record efficiency of 25.05% for TiO₂ ETL based planar FAPbI₃

PSCs as well as highly stable CsFAMA PSCs with pronounced moisture stability (RH of 40%) over 9000 h. Owing to universal and unique property of pulsed laser irradiation, embedding other laser generated nanocrystals has been demonstrated to be general for effectively addressing the carriers loss in TiO_2 ETL. This study may provide a novel pathway of developing efficient ETLs for plannar PSCs based on laser-matter interactions.

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Scheme 1. (a) Schematic illustration of conventional work for the intrinsic carriers mobility of TiO_2 nanoparticles (NPs) based the electrons trapping at TiO_2 NPs. (b) Schematic illustration of present work for the particle boundaries (PBs) engineering of TiO_2 ETL. (c) Schematic illustration of the functions of embedded p-n heterointerfaces.



Figure 1. (a) Schematic illustration of the laser process of CdTe nanocrystals and its embedding in TiO₂ ETL matrix. (b) TEM image of CdTe nanocrystals (inset: nanocrystals size distribution diagram and Mie-scattering image of colloids loading CdTe nanocrystals via laser process). (c) HRTEM and corresponding FFT of CdTe nanocrystals. In-situ xrd patterns of TiO₂ with (e) and without (d) CdTe nanocrystals (*: rutile phase). SEM (f) and AFM (g) images of pristine TiO₂ and 6%-CdTe-TiO₂ based ETLs. Scale bar: (b) 20 nm, (c) 2 nm, (f) 1 μ m, (g) 1 μ m.



Figure 2. (a) Cross-sectional HAADF and corresponding TEM-EDS image of TiO_2 film with CdTe nanocrystals embedding. (b) Cross-sectional TEM and (c) HRTEM (magnified red region) conducted on the glass/FTO/CdTe-TiO2/Pt probing the existence of CdTe nanocrystals in the TiO₂ film. FFT-transformation images for CdTe nanocrystals (d: labeled as red area) and TiO₂ nanocrystals (e: labeled as yellow area). (f) Electron mobilities of different ETLs using the SCLC model, and the inset shows the device structure of ITO/Al/ETLs/Al. (g) UPS results of CdTe nanocrystals. (h) Mott-Schottky plots of CdTe nanocrystals embedded TiO₂ matrix. Scale bar: (a) 20 nm, (b) 50 nm, (c) 5 nm.



Figure 3. SEM images of perovskite films based pristine TiO_2 (a) and 6%-CdTe-TiO_2 films (b). (c) Steadystate and (d) time-resolved PL spectra of CsFAMA perovskite films spin-coated on different TiO_2 layers. (e) Arrhenius plots of the characteristic transition frequencies extracted from the derivative of the admittance spectra. (f) Trap state density (N_T) of the perovskite solar devices was measured at 300 K. High-resolution XPS spectra of Cd $3d_{3/2}$ (g) and I $3d_{3/2}$ (h) for CdTe nanocrystals-PbI₂ powder. (i) Nyquist curves of the devices upon different TiO₂ films measured in the dark at a bias of 0.8 V. Inset: the equivalent circuit model. (j) Energy level diagram for each component of devices upon different TiO₂ layers. The energy levels of Spiro-OMeTAD and Au refer to the literatures.^[28] Scale bar: (a) 500 nm; (b) 500 nm.



Figure 4. (a) *J-V* curves of the devices based on different TiO_2 layers with schematic illustration of device architecture. (b) *J-V* plots of the champion devices upon the pristine TiO_2 and 6%-CdTe- TiO_2 layers measured both in reverse scan and forward scan. Inset: the stabilized photocurrent and power output of champion devices at MPP. (c) *J-V* curves of FAPbI₃ devices based on different TiO_2 layers with stabilized photocurrent and power output. (d) The EQE spectra and the integrated current density for different champion devices. (e) The logarithmic *I-V* plots of different devices based the pristine TiO_2 and 6%-CdTe- TiO_2 layers, respectively. (f) PCE distribution of 50 individual CsFAMA and FAPbI₃ devices. The long-term humidity of CsFAMA-based devices (g), and operational (h) stability of FAPbI₃ devices with and without CdTe nanocrystals embedding. The error bars represent the standard deviation for 20 devices.

Samples	Content of CdTe	Voc (V)	Jsc (mA/cm²)	FF (%)	PCE (۶ Average	%) Best
Pristine TiO ₂	0%	1.107±0.024	22.72±0.58	75.80±2.01	18.95±0.73	20.05
3%-CdTe-TiO ₂	3%	1.155±0.020	23.38±0.33	76.60±1.48	20.66±0.63	21.68
6%-CdTe-TiO ₂	6%	1.162±0.018	23.54±0.30	78.10±1.15	21.23±0.49	22.00
9%-CdTe-TiO ₂	9%	1.141±0.022	22.98±0.47	75.97±1.64	20.15±0.70	20.98
Pristine FAPbI ₃	0%	1.081±0.032	24.81±0.79	78.01±2.28	21.36±1.29	22.72
Target FAPbI ₃	6%	1.131±0.026	25.41±0.60	80.21±1.78	23.81±0.85	25.05

Table 1. Photovoltaic parameters of CsFAMA and FAPbI $_3$ type PSCs upon different TiO $_2$ ETLs.

Entry for the Table of Contents



Present work demonstrates an efficient strategy of the particle boundaries (PBs) embedding of multifunctional p-type semiconducting CdTe nanocrystals for inhibited carrier losses at PBs, which could serve as efficient PBs mediator for boosting the electrons mobility of TiO_2 ETL by maximally three orders of magnitude and consequently result in a new benchmark PCE over 25% in planar PSCs.

Supporting Information

Laser Derived Electron Transport Layers with Embedded p-n Heterointerfaces Enabling Planar Perovskite Solar Cells with Efficiency over 25%

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Experimental Section

Materials

Unless stated otherwise, all materials were purchased from Xi'an Polymer Light Technology Corp without further purification. Fuorine-doped tin oxide (FTO) coated glass substrates with partial etching (~ 20Ω) were purchased from Opvtech. The CdTe and GaAs bulk materials and some solvents including dimethylformamide (DMF), dimethylsulfoxide (DMSO), chlorobenzene (CB) and acetonitrile (ACN) were purchased from Sigma-Aldrich.

Preparation of ligand-free nanocrystals in desired solvent

A certain amount of the bulk materials are transferred into a chamber with 10 ml deionized water, subsequently suffering from a non-focusing nanosecond pulsed laser irradiation (Quantel, repetition rate: 10 Hz, pulse width: 8 ns, beam diameter: 8 mm) with a wavelength of 1064 nm. During the irradiation, continuous ultrasonic treatment is performed to ensure the preparation of homogeneous and well-dispersed CdTe colloids. The tailored laser fluence is applied ranging from 0-600 mJ/pulse cm².

Device Fabrication

Preparation of compact ETLs

The etched FTO substrates were cleaned sequentially in detergent, deionized water, acetone, isopropyl alcohol and ethanol under continuous sonication, and then dried with N_2 flow using a compressed nitrogen gun. After removing organic residues on the FTO substrates and improving wettability using an ozone plasma for 10 min, approximately 50 nm thick TiO₂ compact layer prepared by chemical bath deposition is deposited on clean FTO substrate immersed in a TiCl₄ aqueous solution (2.25:100 volume ratio of TiCl₄ : H₂O) at 70 °C for 1 h, and then annealed at 150 °C for 60 min.

For the target ETLs configuration, TiCl₄ precursor incorporating different content of CTNs with 0.1mg/ml concentration (3%, 6% and 9% volume ratio to TiCl₄ precursor abbreviated as 3%-CdTe-TiO₂, 6%-CdTe-TiO₂, 9%-CdTe-TiO₂) and then together with FTO substrates are subject to similar thermal processing.

Fabrication of perovskite films

CsFAMA-based perovskite: The spin-coating procedure was carried out in a nitrogen glove box. For the fabrication of $Cs_{0.05}FA_{0.81}MA_{0.14}PbI_{2.57}Br_{0.43}$ perovskite, the precursor solution (1.25 M) containing CsI (0.0625M), FAI (1.0125 M), PbI₂ (1.075 M), MABr (0.175M) and PbBr₂ (0.175M) is prepared with a successive stir in a mixture of DMF: DMSO (4:1 v/v) at 55 °C for 2 h. The as-prepared precursor (30 µL) was dropped onto the TiO₂/FTO substrate followed by a consecutive two-step spin-coating process at 2,000 and 4,000 rpm for 10 and 30 s, respectively. During the second spin-coating step, 200 µL anhydrous chlorobenzene (CB) was immediately poured on the substrate 10 s prior to the end of the program. Afterwards, the intermediate phase film is heat-treated on a hotplate at 100 °C for 60 min.

FAPbI₃-based perovskite: The perovskite precursor (1.80M) containing FAI (1.80M), PbI₂ (1.80M), FAHCOO (0.12M) and MACl (0.52M) is prepared with a successive stir in a mixture of DMF: DMSO (4:1 v/v) at 55 °C for 2 h. The as-prepared precursor (30 μ L) was dropped onto the TiO₂/FTO substrate followed by a one-step spin-coating at 6000 rpm for 50 s. During spin coating, 200 μ L chlorobenzene (CB) is immediately poured on the substrate 15 s prior to the end of the program. Subsequently, the intermediate phase film is heat-treated on a hotplate at 150 °C for 10 min.

In addition, for FAPbI₃ type perovskite, 3mg/ml phenethylammonium iodide (PEAI) was spin-coated on the perovskite/TiO₂/FTO substrate at 6000 rpm for 30 s, followed by deposition of Spiro-OMeTAD.

The deposition of HTM and Au electrode

The hole-transport layer was deposited onto the perovskite/TiO₂/FTO substrate via spin-coating at 6000 rpm for 30 s, prepared by dissolving 72.3 mg Spiro-OMeTAD, 29 mL tBP and 18 μ L Li-TFSI (a stock solution of 520 mg/mL in acetonitrile) into 1 mL chlorobenzene. Finally, the Au electrode (80 nm) was deposited on the top of devices by thermal evaporation using a shadow mask. Each electrode of devices exhibits the active area of 0.05 cm², and during testing, that of each electrode is accurately defined by a 0.045 cm² non-reflective mask to eliminate the influence of any edge effects.

Characterization

The SEM images and relevant EDS results were obtained by using a field emission SEM with an EDS detector (FEI Nova). Atomic force microscope (AFM) characterization was performed using a Bruker Dimension Icon. High-resolution transmission electron microscopy (HRTEM) with EDS were characterized employing an FEI Tecnai F30 transmission electron microscope operated at 300 kV, equipped with an Oxford Instruments EDS

detector and a high angle annular dark field (HAADF) STEM detector. The cross sections of TiO2 films are prepared by focused ion beam (FIB). The Raman spectra was acquired using a Raman microscope at an excitation laser wavelength of 532 nm (Renishaw). The X-ray diffraction (XRD) patterns were characterized on a X'pert PRO (PANalytical) using a Cu K α ($\lambda = 0.15406$ nm) as the X-ray source. The absorption plots were acquired using the ultraviolet-visible (UV-vis) spectrophotometer (Perkin-Elmer Lambda 35 UV-vis-NIR). The steady-state photoluminescence and time-resolved photoluminescence (TRPL) spectra were measured using pulse laser excitation source at the wavelength of 470 nm (Horiba FluorologFL-3). The twoprobe conductivity measurements were carried out on an electrochemical workstation (CHI660E) using linear sweep voltammetry at room temperature in the dark. The EIS was measured applying a bias of 0.8V in the dark in a frequency range from 1 MHz to 0.1 Hz (CHI660E). For Mott–Schottky analysis, capacitance–voltage measurements were performed at a frequency of 1 kHz with devices configuration of FTO/ETLs/Perovskite/Au without a hole-transport layer (CHI660E). X-ray photoelectron spectroscopy (XPS) measurements were conducted on an Axis Supra (Kratos). Ultraviolet photoelectron spectroscopy (UPS) was characterized by a VG Scienta R4000 analyzer and the HeI (21.22eV) emission line employed for excitation at a bias of -5V. The contact angles measurements were performed on a Data physics OCA-20 contact-angle system at ambient air. The J–V curves and maximum power point tracking were measured via a Keithley 2400 source meter under simulated AM 1.5G illumination (100 mW cm⁻²) produced by a xenon-lamp-based solar simulator (Oriel 67005, 150 W Solar Simulator) that was calibrated with a monocrystalline silicon reference cell (Hamamatsu S1133). The devices were measured both in reverse scan (+1.2 V to -0.1 V) and forward scan (-0.1 V to +1.2 V) with a scanning rate of 0.2V/s. The EQE was obtained by using a Enlitech EQE measurement system (QE-R3011). The dark I-V characterization of the electron-only devices were measured by a Keithley 2400 source to calculate the trap state density using SCLC model.

Calculations

Calculations for included angle of CdTe crystal planes (cubic):

$$\cos \theta = \frac{h_1 h_2 + k_1 k_2 + l_1 l_2}{\sqrt{(h_1^2 + k_1^2 + l_1^2)(h_2^2 + k_2^2 + l_2^2)}}$$
(1)

Calculations for included angle of TiO₂ crystal planes (tetragonal):

$$\cos \theta = \frac{\frac{h_1 h_2 + k_1 k_2}{a^2} + \frac{l_1 l_2}{c^2}}{\sqrt{\left(\frac{h_1^2 + k_1^2}{a^2} + \frac{l_1^2}{c^2}\right)\left(\frac{h_2^2 + k_2^2}{a^2} + \frac{l_2^2}{c^2}\right)}}$$
(2)

where θ is the included angle of crystal planes; a, c are crystal lattice \Box onstants; h₁, k₁, l₁, h₂, k₂ and l₂ are specific crystal indices.

Calculation for the electron mobility (μ) of different ETLs using SCLC model:

$$\mu = \frac{8JL^3}{9\varepsilon\varepsilon_0(V - V_r - V_{bi})^2} \tag{3}$$

where *J* is the current density; *L* is the thickness of different TiO₂ films; ε_0 and ε are the vacuum permittivity and the dielectric permittivity, respectively. *V* is the applied voltage, *V*_r is the voltage loss resulting from radiative recombination, and *V*_{bi} is the built-in potential.^[1]



Figure S1. CdTe nanoparti □es with different size distribution via various laser fluen □e of 300, 400 and 500 mJ pulse⁻¹ □m⁻² produ □ing nanoparti □es with average size of approximately 73 nm, 25 nm and 17 nm, respe □tively. S □ale bars: (a) 100 nm, (b) 50 nm, (c) 100 nm.



Figure S2. HRTEM (**a**) and Corresponding FFT-transformation (**b**) images for laser generated CdTe nanoparticles, exhibiting the (111) and (220) plane with 35.3° plane included angle. S cale bars: (**a**) 5 nm, (**b**) 1/5 nm.



Figure S3. Raman spectroscopy of CdTe with and without laser irradiation process.



Figure S4. HAADF and corresponding TEM-EDS images of CdTe nanoparticles produced by pulsed laser process with different laser fluence. Scale bar: (a) 20 nm, (b) 10 nm, (c) 10 nm.



Figure S5. XPS spectrum (a) and high-resolution XPS spectra (b) of CdTe nanocrystals. The appearances of Cd 3d_{3/2} peak at 412.1 eV, Cd 3d_{5/2} peak at 405.1 eV, Te 3d_{5/2} peak at 572.8 eV, and Te 3d_{3/2} peak at 583.2 eV confirm the existence of cadmium and tellurium species in the CdTe nanocrystals.^[2]



Figure S6. S hemati illustration of embedding of CdTe nano rystals in the TiO₂ matrix.



Figure S7. SEM top-view images of , 3%-CdTe-TiO₂ and 9%-CdTe-TiO₂ films.



Figure S8. AFM height images of pristine TiO₂, 3%-CdTe-TiO₂, 6%-CdTe-TiO₂ and 9%-CdTe-TiO₂ films.



Figure S9. TEM-EDS mapping of Tross-se Tional TEM for Ti, O, Cd, Te, Cd/Sn, Sn, and Pt elements. S ale bars: 20 nm.



Figure S10. Cross-sectional TEM-EDS results of 9%-CdTe-TiO₂ films with all elements content analyses (**b**) from selected area 1 of HADDF image (**a**) and corresponding schematic diagram (**c**).


Figure S11. Absorption spectra (**a**) of different TiO₂ films and corresponding Tauc plots (**b**) determining the optical bandgaps of TiO₂ films.



Figure S12. Optical transmittance spectra of FTO and different TiO₂ films based on FTO.



Figure S13. *I-V* curves with V_{TFL} related to the N_t of different ETLs and corresponding device structure of FTO/ETL/Au.

Table S1. The dark *I-V* plots of different devices show V_{TFL} (the onset voltage of the trap-filled limit region) kink point behavior, it could be extracted to determine the N_t of various ETLs. The N_t of different ETLs was evaluated using $N_t = \frac{2\epsilon_0 \epsilon V_{TFL}}{qL^2}$ formula, where ϵ_0 is the vacuum permittivity, ϵ is the relative dielectric constant, V_{TFL} is the onset voltage of the trap-filled limit region, q is the elemental charge, and L is the thickness of the ETLs.^[3]

ETLs	$V_{\mathrm{TFL}}\left(\mathrm{V} ight)$	$N_{\rm t}$ (cm ⁻³)
Control	1.14	3.95×10^{16}
3%-CdTe-TiO ₂	0.38	1.32×10^{16}
6%-CdTe-TiO ₂	0.28	9.70×10^{15}
9%-CdTe-TiO ₂	0.46	1.59×10^{16}

Note S1: preparation of the TiO₂ films for the measurement of the electron mobilities

We estimate the electron mobilities of the TiO₂ films as reported in the literature^[4] using the electron-only devices with a structure of ITO/Al/ETLs/Al. Detailly, Al (80 nm) is deposited on ITO substrates, and then TiO₂ powder prepared by drying TiO₂ aqueous solution with and without CdTe-NCs after suffering thermal treatment of 70 °C are dispersed in anhydrous ethanol (5 mg/ml) with the addition of titanium(diisopropoxide) bis(2,4-pentanedionate) (TIPD: 15 ul/ml),^[5] then spin-coated on ITO substrates at 5000 rpm for 40s for twice, followed by annealing at 150°C for 60 min. Finallly, Al (80 nm) is deposited on different samples by thermal evaporation.



Figure S14. Conduction properties of pristine TiO₂ and 6%-CdTe-TiO₂ films measured by a device structure of FTO/ETL/Ag. The conductivity of different ETLs was evaluated using $\sigma = \frac{Id}{VA}$ formula, where A is the area of devices, d is the thickness of ETLs. The results indicate that CdTe embedded TiO₂ yields larger conductivity based larger slope.^[6]



Figure S15. Contact angle of pristine TiO₂ and 6%-CdTe-TiO₂ films dropped by perovskite precursor.



Figure S16. XRD pattern of perovskite films upon different TiO₂ layers.



Figure S17. AFM images of perovskite films based pristine TiO₂ and 6%-CdTe-TiO₂ films.

Table S2. Summaries of parameters extracted from the fitted plots of the corresponding TRPL spectra following bi-exponential rate law: $f(t) = A_1 exp(-t/\tau_1) + A_2 exp(-t/\tau_2) + y_0$, where A_1 and A_2 represent the relative amplitudes, τ_1 represents trap-assisted recombination, and τ_2 represents free carrier recombination. The τ_{ave} is

ETLs	A_1	$\tau_1(ns)$	A_2	τ_2 (ns)	$\tau_{ave} (ns)$
Control	0.67	11.45	0.33	267.51	247.04
3 Vol%	0.58	9.27	0.42	102.91	92.55
6Vol%	0.78	6.44	0.22	99.51	82.14
9 Vol%	0.64	9.90	0.36	122.87	108.72

calculated using $\tau_{ave} = \frac{\sum A_i \tau i^2}{\sum A_i \tau i}$ formula.^[7]



Figure S18. (a) Nyquist Turves of the devices upon different TiO₂ films measured in the dark at a bias of 0.8 V. Inset: the equivalent Tir Tuit model. (b) The high-frequenty regions of Nyquist plots of devices based on different ETLs. The Nyquist plots of devices upon different TiO₂ ETLs measured in the dark under a bias of 0.8 V exhibits two diatritical theratteristicants in high-frequenty and low-frequenty assigned to the Teresistante (R to the recombination resistante (R to finance the equivalent tir Tuit model and relevant fitting parameters are detailly listed in Table S3.

	Rs	Rco	CPE1-T	CPE1-P	Rrec	CPE2-T	CPE2-P
ETLs							
	(Ω)	(Ω)	(F)	(F)	(Ω)	(F)	(F)
control	20.11	27159	6.20E-09	0.94	1.47E05	1.82E-06	0.76
3Vol%	9.86	11740	5.60E-09	0.97	2.21E05	1.11E-06	0.80
6Vol%	4.08	2422	6.30E-09	0.96	3.11E05	1.03E-06	0.87
9Vol%	11.39	20910	4.18E-09	0.99	1.85E05	1.14E-06	0.84

Table S3. Summaries of parameters extracted from the fitted plots of the impedance spectra of devices upon different ETLs using a the equivalent circuit model.



Figure S19. Dark current density-voltage plot of devices based the pristine TiO₂ and 6%-CdTe-TiO₂ films.

The Fermi level (E_F) is calculated using $E_F = E_{cut-off} - 21.22 \text{ eV}$ equation, where $E_{cut-off}$ represents cut-off binding energy, 21.22 eV is the photon energy of excitation light. The $E_{cut-off}$ values are 16.59, 17.54, 17.58, and 17.62 eV for different TiO₂ layers. The calculated E_F values show -4.63, -3.68, -3.64 and -3.60 eV for the pristine TiO₂, 3%-CdTe-TiO₂, 6%-CdTe-TiO₂ and 9%-CdTe-TiO₂ layers, respectively. And then the valence band (E_{VB}) is calculated using $E_{VB} = E_F - E_F$, edge (Fermi edge) equation, obtaining -7.40, -7.36, -7.33 and -7.24 eV for different TiO₂ layers. The band gap (E_g) values of TiO₂ are extracted from the absorption spectrum and corresponding Tauc plot, revealing negligible change on E_g (3.20, 3.22, 3.23 and 3.24 eV for different TiO₂ layers). Then the calculated conduction bands (E_{CB}) reflected from E_g and E_{VB} are -4.20, -4.14, -4.10 and -4.00 eV for different TiO₂ layers, respectively.



Figure S20. Different TiO₂ films with UPS Fermi edge (left) and the cut-off energy (right).



Figure S21. UPS Fermi edge (a) and the \Box ut-off energy (b) of CsFAMA perovskite film.

Note S3: AS and Mott-Schottky analyses

By virtue of the limitations of the space charge-limited current (SCLC) for defect analysis, temperature dependent admittance spectroscopy (AS) along with Mott-Schottky analysis is widely employed to quantitatively characterize both shallow and deep defects in PSCs. As the previous literature reported,^[8] the defect activation energies (E_a) is derived via the equation:

$$\omega_0 = \beta T^2 exp(\frac{-E_a}{K_b T})$$

where ω_0 is the characteristic transition angular frequency, β is a temperature dependent parameter, T is the temperature and K_b is the Boltzmann's constant, respectively. The derivative of the capacitance spectrum can be used to determine ω_0 from the peak value of the $[-\omega \times dC/d\omega]$ curve. According to this relation, the value of E_a can be obtained from the slope of the graph by fitting the Arrhenius formula^[8]:

$$\ln\frac{\omega_0}{T^2} = \ln\beta - \frac{E_a}{K_b T}$$

the trap density (N_t) can be derived pursuant to the equation^[8]:

$$N_t(E_\omega) = -\frac{V_{bi}}{qW} \frac{dC}{d\omega} \frac{\omega}{k_b T}$$
, $E_\omega = K_b T \ln \frac{\beta T^2}{\omega}$

where V_{bi} is the built-in potential, *C* is the capacitance, ω is the applied angular frequency, *q* is the elementary charge, and *W* is the depletion width, respectively. The values of V_{bi} and *W* were obtained from the Mott-Schottky analysis^[8]:

$$\frac{A^2}{C^2} = \frac{2(V_{bi} - V)}{q\varepsilon\varepsilon_0 N}$$

where *A* is the active area, ε is the static permittivity of perovskite, ε_0 is the permittivity of free space, *N* is the apparent doping profile in the depleted layer, and *V* is the applied bias. The Mott-Schottky plot describes a straight line where the intersection on the bias axis determines V_{bi} and the slope gives the impurity doping density *N*. Then, the depletion width $W = \sqrt{\frac{2\varepsilon\varepsilon_0 V_{bi}}{qN}}$ corresponding to the zero bias can be calculated, the distribution of trap state density can be calculated. **Figure S22a-b** show the temperature-dependent AS of the devices based pristine TiO₂ and 6%-CdTe-TiO₂ ETLs measured at temperatures ranging from 210 to 330 K under the dark. The corresponding Arrhenius plots of these two devices describe the relationship between the characteristic transition frequencies and temperatures (**Figure 3e**). The calculated defect activation energies (*E*_a) of the control and target samples are 0.252 and 0.219 eV, respectively. The built-in potential (*V*_{bi}) and the depletion width (*W*) are determined from the Mott-Schottky analysis of the data measured at a frequency of 1 kHz with bias potential from 0 to 1.2 V (**Figure S22c**). After calculation, the *V*_{bi} of the control and target devices are 0.92 and 1.02 V, and the corresponding depletion widths are 185 and 290 nm, respectively. The density distribution and energy levels of trap states in these two devices are then presented in **Figure 3f**.



Figure S22. Admittance spectra of devices based on pristine TiO_2 (a) and 6%-CdTe- TiO_2 (b) ETLs, which were measured at gradient temperatures between 210 and 330 K with a step of 15 K. (c) Mott-Schottky analysis at 1 kHz for extraction of the built-in potential (V_{bi}).



Figure S23. XPS spectrum of PbI_2 powder (a), high-resolution XPS spectra of Pb 4f (b) and I 3d (c).



Figure S24. XPS spectrum of CdTe nanocrystals-PbI₂ (a), high-resolution XPS spectra of Pb 4f (b). In order to confirm the interaction between CdTe and PbI₂, an excess of CdTe nanocrystals powder and PbI₂ (molar ratio: 3:1) were mixed together in DMF with continuous stirring at 55°C for 2h. The products were washed with toluene for several times and were finally dried under a vacuum for further characterizations.



Figure S25. Cross-sectional SEM image of device upon the best ETLs.



Figure S26. Schematic illustration for fabri ation process of devices based on CdTe nano rystals embedded

ETLs.

Scanning direction	Voc (V)	Jsc (mA/cm²)	FF (%)	PCE (%)
Reverse 6%-CdTe-TiO ₂	1.191	23.76	77.74	22.00
Forward 6%-CdTe-TiO ₂	1.172	23.77	78.27	21.80
Reverse pristine TiO ₂	1.149	22.81	76.52	20.05
Forward pristine TiO ₂	1.141	22.87	67.02	17.49

Table S4. Performance parameters of champion devices upon different TiO_2 layers obtained from the J-V curves measured both in reverse scan and forward scan.



Figure S27. Cross-sectional SEM images of FAPbI₃ based devices with the pristine TiO_2 and 6%-CdTe- TiO_2 layer. Scale bars: 1 μ m.



Figure S28. SEM images of FAPbI₃ perovskite based on the pristine TiO_2 (a) and 6%-CdTe-NCs- TiO_2 (b) layers. S ale bars: 2 μ m.



Figure S29. AFM images of FAPbI₃ perovskite based on the pristine TiO_2 a) and 6%-CdTe- TiO_2 b) layers. Scale bars: 1 μ m.



Figure S30. XRD pattern of $FAPbI_3$ upon different TiO_2 layers.



Figure S31. PL (a) and TRPL (b) spectra of FAPbI₃ perovskite based on the pristine TiO₂ and 6%-CdTe-TiO₂ layers.

ETLs	A ₁	7 1 (ns)	A_2	7 ₂ (ns)	T _{ave} (ns)
Pristine TiO ₂	0.20	175.43	0.80	1212.01	1175.81
6%-CdTe-TiO ₂	0.25	123.40	0.75	781.76	748.85

Table S5. Summaries of fitting parameters for TRPL spectra with a structure of FAPbI₃/TiO2 (with andwithout CdTe nanocrystals)/FTO substrates.



Figure S32. The Tauc plot of $FAPbI_3$ film for obtaining band gap.



Figure S33. Statistical distributions of the photovoltaic parameters including Voc, Jsc and FF for CsFAMA devices upon different TiO_2 layers. The statistics of PCEs for 50 devices is shown in Figure 4f, which exhibit an excellent reproducibility of the performance improvement with the construction of CsFAMA/CdTe-TiO₂. The average PCE rises from 18.95% to 21.23%, with Voc from 1.108 to 1.160 V, the average Jsc from 22.70

to 23.54 mA/cm², and the FF from 75.71 to 78.21%. It is obvious that all the parameters, have been improved possibly owing to the better film quality and excellent electronic structure in CsFAMA/CdTe-TiO₂

devices.



Figure S34. Statistical distributions of the photovoltaic parameters including Voc, Jsc and FF for FAPbI₃ devices upon different TiO₂ layers. The statistics of PCEs for 50 devices is shown in Figure 4f, which exhibit an excellent reproducibility of the performance improvement with the construction of FAPbI₃/CdTe-TiO₂. The average PCE increases from 21.36% to 23.81%, with Voc from 1.081 to 1.131 V, the average Jsc from 24.81 to 25.41 mA/cm², and the FF from 78.01 to 80.53%.



Figure S35. The long-term thermal (a) and operational (b) stability of CsFAMA-based devices with and without CdTe nanocrystals embedding. The target devices exhibit excellent thermal stability with PCE degradation of less than 20% in contrast with that of control devices (over 80%) after continuous thermal treatment of 85 °C for 500 h, as well as operational stability by retaining 90% of initial PCE over 200 h in contrast with that of control devices (approximately 40% over 150 h) at MPP tracking under full-sun illumination (AM 1.5G) in ambient air with RH of 70%.



Figure S36. The long-term humidity stability of FAPbI₃-based devices with and without CdTe nanocrystals embedding.

Electron Mobility/ Best Device electrical PCE Type Ref conductivity structure (%) $(cm^2 V^{-1} s^{-1}/S \cdot m^{-1})$ ITO/TiO2-Urea NH₂/Cs_{0.05}(FA_{0.87}MA_{0.13})_{0.95}PbI_{2.55}Br_{0.45}/Spir 21.33 7.56×10^{-3} [9] o-OMeTAD/Au FTO/D-TiO₂/Cs_{0.05}FA_{0.81}MA_{0.14}PbI_{2.55}Br_{0.45} Dopamine 20.93 $10^{-4} \sim 10^{-3}$ [10] Spiro-OMeTAD/Au FTO/LD-TiO₂/FAMAPbI₃/Spiro-LD-TiO₂ 24.81 4.51×10^{-5} [11] OMeTAD/Au FTO/cp-TiO₂/TiO₂ or B-TiO₂/Spiro- 1.69×10^{-4} Boron 20.51 [6] OMeTAD/Au Au FTO/TiO₂/C-PCBA/CsFAPb(IBr)₃/Spiro-C-PCBA 5.76×10^{-3} 24.80 [12] OMeTAD/Au 3.83×10^{-4} [BMIM]BF₄ FTO/TiO₂/[BMIM]BF₄/MAPbI₃/PTAA/Au 19.62 [13] 9.29×10^{-3} CO₂-plasma 15.39 [14] FTO/ TiO₂/CO₂-plasma/MAPbI₃/Spiro-

Table S6. Reported values for the conductivity of different planar TiO_2 ETLs.

OMeTAD/Au

Cl-SnO ₂ QDs	FTO/SnO ₂ -QDs/TiO ₂ /Perovskite/ Spiro- OMeTAD/Au	14.90		[15]
CuI	FTO/TiO2@CuI/CsFAMA/Spiro- OMeTAD/Au	19.00	3.98×10 ⁻⁴	[16]
Phosphoren e	FTO/TiO2/phosphorene/Perovskite/Spiro- OMeTAD/Au	17.85		[17]
PDI ₂	ITO/TiO2/PDI2/MAPbI3/Spiro- OMeTAD/Ag	19.84		[18]
Heparin sodium	FTO/TiO2/heparin sodium/MAPbI3/Spiro- OMeTAD/Au	20.1		[19]
GQD	FTO/c-TiO ₂ @ GQD/MAPbI ₃ /Spiro- OMeTAD/Au	19.11		[20]
GQD	FTO/c-TiO2/GQD /MAPbI3/GQD/Spiro-OMeTAD/Au	19.89		[21]
CdTe nanocrystals	<u>FTO/TiO2@CdTe/Cs0.05(FA0.85MA0.15)0.95PbI2.55B</u> <u>r0.45</u> /Spiro-OMeTAD/Au	22.00	1.89 × 10 ⁻²	Our work

CdTe				Our
nanocrystals	FTO/TiO2@CdTe/FAPbI3/Spiro-OMeTAD/Au	25.05	1.89×10^{-2}	work


Figure S37. *J-V* curves of the CsFAMA devices based on TiO₂ layers with different concentration of CdTe nanocrystals.

ETLs	Concentration of CdTe nanocrystals	Voc (V)	Jsc (mA/cm²)	FF (%)	PCE (%)
Pristine TiO ₂	0 mg/ml	1.130	22.85	75.80	19.57
0.05mg/ml CdTe- TiO ₂	0.05 mg/ml	1.141	23.41	78.28	20.91
0.1mg/ml CdTe- TiO ₂	0.1 mg/ml	1.161	23.61	78.43	21.51
0.2mg/ml CdTe- TiO ₂	0.2 mg/ml	1.158	22.97	77.50	20.61
0.3mg/ml CdTe- TiO ₂	0.3 mg/ml	1.125	22.41	75.25	18.97

Table S7. Performance parameters of champion CSFAMA devices based on CdTe@TiO2 layers.



Figure S38. (a) TEM image of GaAs nanocrystals (inset: Mie-scattering image of colloids, HRTEM of GaAs nanocrystals and nanocrystals size distribution diagram,). (b) HAADF and corresponding TEM-EDS images of GaAs nanoparticles. Scale bar: (a) 50 nm, inset: 5 nm, (b) top: 10 nm, middle: 5 nm, bottom: 5 nm.



Figure S39. (a) EDS mapping images for 9%-GaAs-TiO₂ film. (b) SEM top-view images of different TiO₂ and CsFAMA perovskite films deposited on the corresponding TiO₂ films. AFM height images of different TiO₂ and CsFAMA perovskite films based the pristine TiO₂ and 6%-GaAs-NCs-TiO₂ films.



Figure S40. (a) The optical bandgaps of GaAs nanocrystals estimated by UV-visible diffuse-reflectance spectrum (UV-Vis DRS); (b) UPS spectrum of GaAs nanocrystals with UPS Fermi edge (c) and the cut-off energy (d); (e) Mott-Schottky plots of GaAs nanocrystals embedded TiO₂ matrix; (f) UPS spectra of different TiO₂ films with and without GaAs nanoTystals embedding; (g) Electron mobilities of different ETL using the SCLC model, and the inset shows the device structure of ITO/Al/6%-GaAs-TiO₂/Al.



Figure S41. (a) *J-V* curves of CsFAMA-based devices with and without GaAs embedding. (b) *J-V* curves of FAPbI₃-based devices with and without GaAs embedding.

ETLs	Voc (V)	Jsc (mA/cm ²)	FF (%)	PCE (%)
pristine TiO2/CsFAMA	1.126	22.50	77.83	19.73
6%-GaAs- TiO2/CsFAMA	1.172	23.60	77.89	21.54
pristine TiO ₂ /FAPbI ₃	1.089	25.00	80.29	21.86
6%-GaAs- TiO2/FAPbI3	1.130	25.65	81.05	23.49

Table S8. Performance parameters of champion devices upon the pristine and 6%-GaAs-NCs-TiO₂ layers.



Figure S42. The distribution of photovoltaic parameters for CsFAMA-based devices with and without GaAs embedding.

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